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**Modeling unsteady state leachate flow in a landfill using finite
difference and boundary element methods**

Ahmed, Shabbir, Ph.D.

City University of New York, 1992

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A

**MODELING UNSTEADY STATE LEACHATE FLOW
IN A LANDFILL USING FINITE DIFFERENCE
AND BOUNDARY ELEMENT METHODS**

by

SHABBIR AHMED

A dissertation submitted to the Graduate Faculty in Engineering in partial fulfillment of the requirements for the degree of Doctor of Philosophy, The City University of New York.

1992

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SHABBIR AHMED

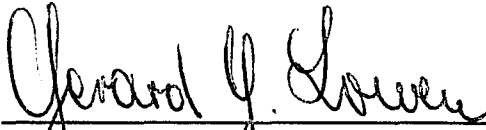
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This manuscript has been read and accepted for the Graduate Faculty in Engineering in satisfaction of the dissertation requirement for the degree of Doctor of Philosophy.

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ABSTRACT

MODELING UNSTEADY STATE LEACHATE FLOW IN A LANDFILL USING FINITE DIFFERENCE AND BOUNDARY ELEMENT METHODS

by

Shabbir Ahmed

Advisor: Professor Reza M. Khanbilvardi

The physical processes involving leachate flow in a solid waste landfill are described by the unsaturated flow through the refuse to the saturated leachate mound at the bottom of a landfill. The moisture-flow in the unsaturated zone helps build up the saturated leachate mound at the bottom of a landfill. The moisture content in the unsaturated zone is obtained by solving the two-dimensional unsaturated moisture-flow equation using numerical techniques. A two-dimensional unsteady state Flow Investigation for Landfill Leachate (FILL) model, based on the implicit finite-difference technique, has been developed to describe the leachate flow process in a landfill.

To obtain accuracy and efficiency in numerical modeling, it is important to investigate the numerical solution techniques suitable to solve the governing equations. Accuracy and efficiency of the boundary integral method over the finite-difference method has been investigated. Two approaches, direct Green's function

and perturbation Green's function formulations, have been developed to solve the unsaturated flow problem. Direct Green's function and perturbation Green's function boundary integral solutions are found to be more accurate than both the Gauss-Seidel iteration and Gauss-Jordan elimination method of finite-difference solution.

The efficiency of the boundary integral formulation for the computation of the moisture-flux is an advantage that is useful to estimate leachate accretion in a landfill. A close agreement of the internal fluxes with the exact solution shows the ability of the boundary integral methods to compute accurate recharge from the unsaturated zone to the saturated leachate mound.

To gain confidence in the finite-difference scheme for the saturated leachate flow equation, two numerical methods, Gauss-Seidel iteration and Gauss-Jordan elimination, are used to solve the finite-difference expressions. Satisfactory performance of the two methods are observed in computing leachate mound-head for leachate accretion from the unsaturated zone. The finite-difference FILL model representing the unsaturated-saturated flow has been applied to an example landfill section to demonstrate the effectiveness of the model in providing information such as evapotranspiration, surface runoff, recharge, and leachate flow rates in both the lateral and vertical directions.

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CHAPTER 1

INTRODUCTION AND BACKGROUND

1.1 INTRODUCTION

Leachate is formed from that portion of precipitation that infiltrates through a solid waste landfill. The water extracts dissolved and/or suspended solids from the landfilled wastes. The resulting contaminated water is called leachate; it has the potential to pollute both groundwater and surface water resources. Because of its importance when investigating the risks due to groundwater and surface water contamination, an estimate of the magnitude and timing of the discharged leachate is desirable.

Mathematical models are used for evaluating a variety of scenarios involved in the leachate flow processes. The complexities involved in the simulation of leachate flow process by mathematical models are a major focus of research at the present time. Of particular interest and difficulty are the simulation of leachate transport that occur in both the unsaturated and saturated zones in a landfill. To obtain accuracy and efficiency in the simulation for moisture transport, considerable research is invested for the development of high-quality numerical models.

To estimate the net infiltration into the landfill, the amount of surface runoff and

evapotranspiration needs to be predicted accurately. In a landfill, side slope has an important bearing on the amount of surface runoff. However, models currently used in the field of landfills do not take into consideration the effect of slope and surface roughness (Fenn et al., 1975; Schroeder et al., 1984; Demetracopoulos et al., 1986).

Evapotranspiration depends on the climatologic and soil data of the landfill area. The modified Penman method developed by Ritchie (1972) and used in Chemical Runoff and Erosion from Agricultural Management Systems (CREAMS) model (Knisel, 1980) and Hydrologic Evaluation of Landfill Performance (HELP) model (Schroeder et al., 1984), takes into consideration the effects of climatologic (temperature and solar radiation) and soil data (soil moisture, porosity, field capacity, and wilting point) to compute evapotranspiration. This method can be used for a realistic estimate of evapotranspiration from a landfill surface.

The total precipitation minus the runoff and evapotranspiration equals the amount of water that contributes to the leachate already present within the landfill. Some of the infiltrated water, however, is retained by the soil matrix in the unsaturated zone. The moisture variation in the unsaturated zone generates leachate accretion. The saturated leachate mound is developed by the leachate accretion from the unsaturated zone. The governing partial differential equations describing the leachate flow phenomena in both the unsaturated and saturated zones of a landfill are obtained from the continuity equation and Darcy's Law. The time-history of

leachate accretion is obtained by numerical solution of the partial differential equation describing the moisture-flow in the unsaturated zone in a landfill.

The numerical solution of the unsaturated moisture-flow equation can be done either in steady state or unsteady state flow condition. Since the process of the development of leachate mound-head is strictly unsteady due to changes in the hydraulic conductivity and diffusivity of the medium, it would be realistic to develop numerical models for leachate accretion rate based on the unsteady state solution of the unsaturated moisture-flow equation. The unsteady leachate flow models can be based on the solution of either one-dimensional or two-dimensional moisture-flow equation in the unsaturated zone in a landfill. The unsteady one-dimensional leachate flow model considers leachate accretion only in the vertical direction. However, the leachate accretion rate in the field changes with respect to distance and time. So, a realistic estimate of the leachate flow can be obtained by solving the two-dimensional moisture-flow equation in the vertical section of a landfill.

High-quality numerical models depend on the efficient and accurate numerical solution procedure. Finite-difference method was used extensively in porous media flow problems since 1960s (Tysen and Weber, 1963; Richtmeyer and Morton, 1967; Cooley, 1971; Prickett and Lonquist, 1971; Remson et al., 1971, Thomas, 1973; Abnish, 1975; Abbott, 1982; Rulon et al., 1985; Demetracopoulos et al., 1986; McKeon and Chu, 1987; and Demetracopoulos, 1988). The finite-difference method requires discretization of the whole domain along with the adjustment of

the geometry of the boundary to fit the finite-difference grid spacing. The method requires to establish grid size such that the variable is satisfactorily approximated everywhere in the domain.

Finite element method was used since 1970s (Neuman, 1973; Gupta and Tanji, 1976; Kuppusamy et al., 1987). This method permits to consider small elements in regions of rapidly varying flow and along the boundary to represent the physical geometry of the domain more realistically. In the finite element method, it is required to find an element size in order to approximate the partial differential equation everywhere in the domain. In such method, the internal fluxes can not be calculated directly in the solution region. The internal fluxes are important parameter to describe leachate flow problem specifically to estimate leachate accretion from the unsaturated zone in a landfill. In both the finite-difference and finite element methods, the internal fluxes can not be calculated directly in the solution domain.

Recently, researchers have begun to employ boundary integral method (also known as boundary element method) to solve problems in saturated porous media flow (Aral and Tang, 1988; Fogden et al., 1988; Lafe and Cheng, 1987). The basic idea is that through some mathematical identity the governing differential equation on the domain is transformed to an integral equation on the boundary. There is always a singularity associated with the transformation which must be calculated. In the process of transformation of the differential equation to the boundary integral

equation, a three-dimensional problem is reduced to a surface integral (two-dimensional problem) and a two-dimensional problem is reduced to a line integral (one-dimensional problem). The new integral equation and boundary conditions are then solved through discretization of the boundary which leads to a system of algebraic equations. Solution of the system of equations results in information about parameters on the boundary which permits the solution of the differential equation at any point in the field.

One important feature of the boundary integral method is that one can determine the dependent variable directly from the mathematical formulation at any point in the field, whereas, in finite element or finite difference method, one can only calculate the dependent variable at node points. The velocity at the internal points is an important parameter for solving groundwater flow problems. An accurate evaluation of the velocity of the hydraulic flow is needed to solve contaminant transport problems. The knowledge of the velocities of leachate flow in the unsaturated zone in a landfill is useful to determine leachate accretion. Unlike finite-difference and finite element methods, the velocities at any point in the domain can be calculated directly using the boundary integral formulation. The ability of the boundary integral method to determine velocities at any point in the domain help estimate recharge to the saturated leachate mound in a landfill.

In this dissertation research, an unsteady two-dimensional Flow Investigation for Landfill Leachate (FILL) model is developed using the finite difference method. In

the FILL model, the unsaturated and the saturated leachate flow equations are solved in a two-dimensional vertical domain using the finite difference method. A finite-difference scheme is also used for the kinematic wave equation to compute surface runoff on the upper boundary. Infiltration is a function of moisture content, hydraulic conductivity, and depth of water on the soil surface. The infiltration equation developed by Philip (1969), which considers the effect of the aforementioned parameters, is used in the present model.

The boundary integral equation method is used to investigate the leachate flow in the unsaturated zone of a landfill. Two boundary integral approaches are developed to solve the two-dimensional unsaturated leachate flow equation. In one approach, the time-dependent Green's function is used to compute the moisture content and the moisture flux in the domain. In the second approach, a perturbation Green's function boundary integral solution process is developed to compute both moisture content and moisture-flux in a landfill.

The accuracy, efficiency and representation of the physical system using finite difference and boundary element methods are compared by solving the two-dimensional leachate flow equation in the unsaturated zone of a landfill. The two methods are also evaluated with respect to the ability to determine the function and flux in the domain of a landfill.

1.2 PRESENT STATUS OF THE RESEARCH

Water balance method is one of the simple procedures to compute leachate flow from a landfill. This method was used in the Environmental Protection Agency (EPA) model developed by Fenn et al. (1975). In the water balance method, the amount of water percolating through the solid waste landfill is obtained by subtracting surface runoff, change in soil moisture of the cover soil, and evapotranspiration from the total precipitation. In this method, moisture movement through the refuse is not considered. The EPA model was used in a number of studies to estimate leachate generation rates for landfills in Cincinnati, Ohio; Orlando, Florida; and Los Angeles, California. Later models developed by Dass et al. (1977), Perrier and Gibson (1980), and Gee (1981) were also based on the water balance techniques. None of the aforementioned models include leachate flow in the unsaturated zone and computation of leachate mound-head generated due to leachate accumulation at the bottom of a landfill.

Recent mathematical models consider moisture-flow through the landfill in both steady and unsteady state flow conditions, and estimate leachate flow vertically and laterally in a landfill. Schroeder et al. (1984) developed the Hydrologic Evaluation of Landfill Performance (HELP) model, which is based on a quasi two-dimensional and quasi steady state approach for estimating leachate flow from waste disposal sites. This model is quasi two-dimensional, because it does not compute the vertical and lateral components of flow in each layer of the landfill

profile simultaneously. The layers in which the hydraulic conductivity is large enough so that the flow occurs only downward and is not significantly restricted are called vertical percolation layers. According to HELP, in these layers, lateral flow components are assumed to be zero. Layers in which hydraulic conductivity is high, but the flow is restricted in the downward direction by an impervious layer, are considered lateral drainage layers. The model is called quasi steady state because the vertical flow component is computed using an unsteady moisture routing equation and the lateral flow component is obtained from the steady state solution of the Boussinesq equation.

In the HELP model, the surface runoff at the upper boundary of the landfill is computed using a soil conservation service (SCS) curve-number method, which does not consider the effect of surface roughness and slope along the side of the landfill. The modified Penman method developed by Ritchie (1972) and used in the Chemical Runoff and Erosion from Agricultural Management Systems (CREAMS) model (Knisel, 1980) is used to compute the potential evapotranspiration in the HELP model. Despite the quasi two-dimensional and quasi steady state treatment of moisture movement, the HELP model is a deterministic computer-based model that uses climatologic and soil data to compute the leachate flow rates in a landfill.

Application of flow models to simulate and estimate unsteady, unsaturated flows in solid waste landfills is still at an early stage. The governing equation used to characterize the leachate accretion on the leachate mound in the unsaturated zone

is a second order partial differential equation. There are few analytical solutions for the unsaturated moisture-flow equation due to the nonlinearity and heterogeneities involved in the properties (for example hydraulic conductivity and diffusivity) of the landfill wastes. Exact analytical solutions can be obtained by making simplifying assumptions regarding soil moisture characteristics and the flow domain. Analytical solutions are described by Philip (1968, 1971), Wooding (1968), Lomen and Warrick (1974), Warrick (1975), and Batu (1982).

Since the irregularities in the domain and the complexities in the moisture-flow equation limit the representation of the real flow field by analytical models, many researchers have attempted to develop numerical models. Abbott (1982) described the numerical solution of one-dimensional unsaturated moisture-flow equation by direct explicit schemes for which a stability criteria has to be satisfied to compute the moisture content. A time-varying, one-dimensional leachate flow model was developed by Korfiatis (1984) using the finite-difference numerical technique. Demetracopoulos et al. (1986) described the implicit finite-difference scheme for the one-dimensional unsteady state moisture-flow equation which is solved to obtain the leachate accretion on the leachate mound. The implicit finite-difference expression for the moisture-flow equation was solved by the successive application of the Gaussian elimination method. Korfiatis et al. (1984) performed an experimental investigation with a small solid waste column, which was used to calibrate and verify the unsteady one-dimensional leachate flow model. The measured and predicted results were reasonably close to each other (Korfiatis et

al., 1984).

McKeon and Chu (1987) described a steady state numerical model to compute suction head in unsaturated soils by the multigrid method. This method adopts an iterative solution technique that approximates the differential equation at various grid levels. The implicit finite-difference scheme for one-dimensional unsaturated flow were also described by Wright and Miller (1982). The model developed by Wright and Miller (1982) predicts suction head and leakage through landfill clay cover.

A Finite Element Model of Waste (FEMWASTE) transport was developed by Yeh and Ward (1981) that predicts the fate and transport of contaminants through saturated-unsaturated porous media. The finite element technique was also used by Yeh (1987) to solve the saturated-unsaturated flow problems through porous media. The three-dimensional Finite Element Model of Water (FEMWATER) flow through saturated-unsaturated porous media, developed by Yeh (1987) does not represent the specific flow problems related to the leachate flow in a landfill. The model does not predict surface runoff and evapotranspiration that have considerable effects on the leachate flow in a landfill.

To obtain accuracy and efficiency in numerical modeling, considerable research is being done for the development of high-quality numerical simulation techniques. In this study, the boundary integral method is employed to simulate leachate

transport in the unsaturated zone of a landfill. The application of the boundary integral method to flow in porous media was first carried out by Liggett (1977), Liu and Liggett (1979), and Lennon et. al. (1980). These problems are governed by the Laplace equation for the velocity potential with a kinematic boundary condition on the free surface. The boundary integral equation for a moving interface between two fluids in porous media was solved by Liu et. al. (1981).

Taigbenu et al. (1984) used the boundary integral method to solve problems of seawater intrusion into a freshwater aquifer. Taigbenu and Liggett (1986) developed boundary integral model for contaminant transport in groundwater system. Taigbenu and Liggett (1986) also developed boundary integral method to solve Boussinesq equation for flow in an unconfined aquifer. Liggett and Liu (1979) developed two boundary integral formulations, Laplace transform solution and direct Green's function solution, to model unsteady flow in confined aquifers. Taigbenu and Liggett (1985) used the time-dependent Green's function to model groundwater flow problems considering a dimensionless diffusion equation.

Heterogeneity in the groundwater flow problems were investigated by Lape et. al. (1981) by dividing the aquifer into a three-dimensional grid of zones. Zoning technique was employed to subdivide the domain into piecewise homogeneous regions. Cheng (1984) also developed a boundary integral equation method for Darcy's law with variable permeability. The formulation was done using the Green's function that varies with the permeability. A perturbation boundary element code

was developed by Lafe and Cheng (1987) for the solution of the steady state groundwater equation in a heterogeneous media. Aral and Tang (1988) developed a numerical time integration procedure for the boundary element solution of the time-dependent anisotropic aquifer problems.

Cheng and Liggett (1984) developed boundary integral formulations for the solution of problems related to fracture propagation and consolidation of fluid-infiltrated porous media. The boundary element model for porous-elasticity was used to simulate the consolidation of partially saturated soil. The model results revealed that a small reduction in the degree of saturation would drastically alter the consolidation and pore pressure distribution of partially saturated soils. The model would be useful for simulating the problems incorporating the effects of fracture and consolidation in porous media flow.

The aforementioned applications of the boundary integral method were mostly used to solve saturated flow problems in the groundwater system. So far, the boundary integral equation method for unsteady flow of moisture transport in the unsaturated zone of a landfill has not been reported in the literature.

1.3 OBJECTIVE

The objectives of this research are:

1. To develop a two-dimensional unsteady state leachate flow model capable of simulating leachate flow rates in space and time. The model is developed with the capability to predict the unsteady variation of leachate mound-head inside the landfill;
2. To develop a runoff submodel that is incorporated into the main two-dimensional leachate flow model by considering the effect of side slope on the total runoff that occurs from a landfill;
3. To investigate the time-varying net infiltration rate that causes the build up of leachate mound at the bottom of the landfill;
4. To investigate the effect of finite-difference and boundary integral method in solving the moisture-flow equation in the unsaturated zone of a landfill. Results obtained by these techniques are compared with the simplified analytical solutions; and
5. To apply the two-dimensional unsteady state finite-difference FILL model to an example landfill section to demonstrate its effectiveness in predicting evapotranspiration, runoff, net infiltration, recharge, lateral and vertical flow rates in a landfill. Sensitivity analysis is also carried out in order to evaluate the influence of input parameters upon the model performance.

CHAPTER 2

LEACHATE FLOW PROCESS

This chapter presents the governing equations defining the leachate flow process causing the development of leachate mound in a landfill. Formulation of the FILL model along with the boundary conditions are described. The input to the upper boundary of the solid waste landfill is the precipitation less runoff and evapotranspiration. Runoff occurs at the top of the landfill and depends on the slope, roughness, hydraulic conductivity, and moisture content. The portion of the precipitation that enters into the landfill surface is the infiltration which is dependent on the depth of flow on the surface and also on the hydraulic conductivity and moisture content of the top soil. The moisture flow in the landfill basically occurs in three phases - flow through the cover, flow through the wastes, and flow through the bottom liner or virgin ground beneath the wastes as shown in Figure 2-1. The fraction of infiltration that enters into the unsaturated zone in a landfill changes the moisture content and thus leads to a build-up of leachate to form a saturated leachate mound. Then the lateral and vertical flows can be estimated when the hydraulic gradient of the leachate mound-head in the saturated zone is known.

The movement of moisture inside a landfill can be estimated by using a simple one-dimensional model. In this case, it is assumed that the flow in one vertical

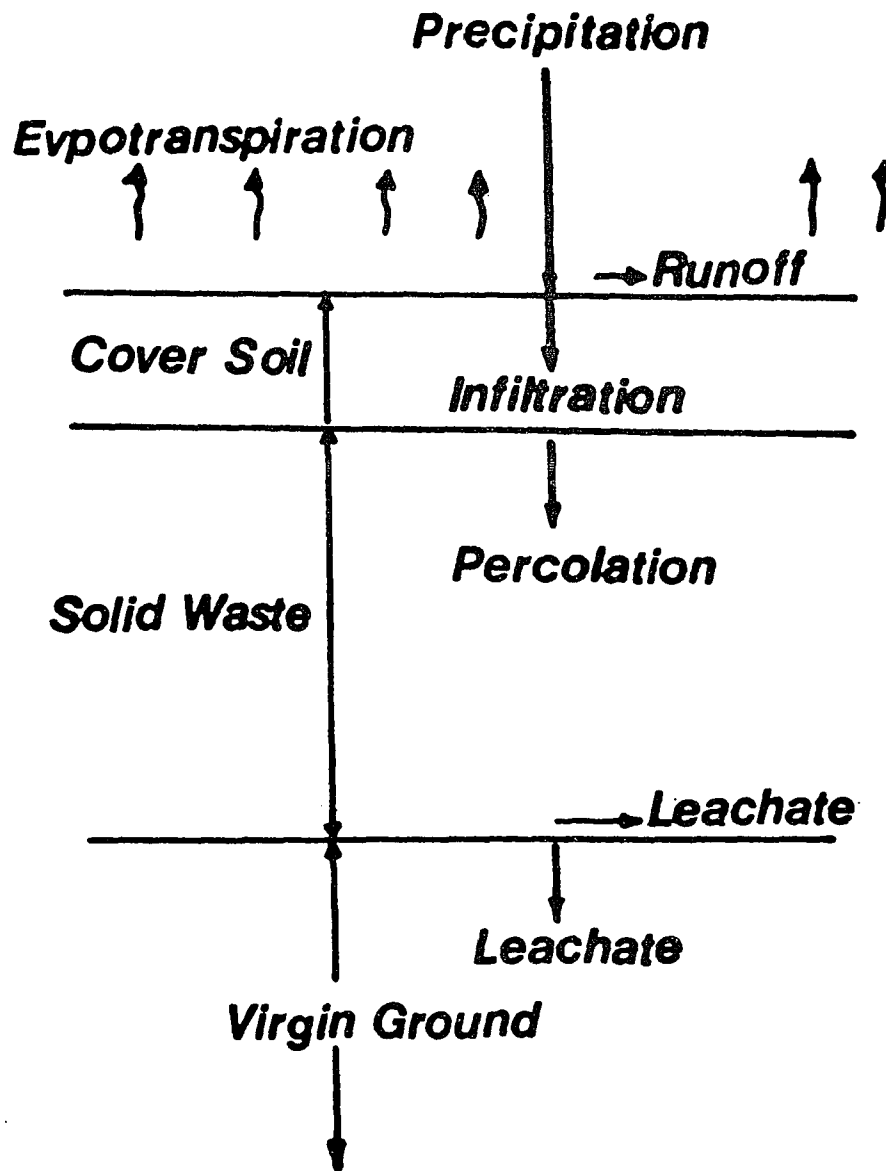


Figure 2-1. Configuration of Leachate Flow Process in a Landfill.

column of a landfill is not dependent on the changes in the moisture content in other parts of the landfill. However, moisture flow in one column depends on changes in the moisture content in the other columns of a landfill. These effects can be taken into account by considering a two-dimensional flow condition. Since the leachate flow varies in both space and time, an unsteady state two-dimensional flow with time-varying boundary conditions can be considered to predict the time-history of leachate flow inside a landfill for the real field simulation.

2.1 GOVERNING EQUATION IN THE UNSATURATED ZONE

The unsteady moisture-flow in the unsaturated zone can be formulated by using Darcy's Law and the equation of continuity. Darcy's Law was formulated for flow through a saturated porous media where hydraulic conductivity is a function of the pore properties of the medium and the properties of the fluid. Buckingham (1907) established that Darcy's Law might apply to the liquid flow in the unsaturated media in a modified form in which the hydraulic conductivity K , instead of being a constant, is regarded as a function of the volumetric moisture content. The extension of Darcy's Law was verified later by several investigators (Gardner, 1958; Childs and Collis-George, 1959; Brooks and Corey, 1964; Brutsaert, 1967; Klute, 1972; Mualem, 1976; and Clapp and Hornberger, 1978) and it is now a generally accepted theory for unsaturated liquid flow. According to Darcy's Law, the specific discharge in the unsaturated zone is:

$$q = -K(\theta) \nabla \phi \quad (2-1)$$

where,

q = volume flow rate per unit cross-sectional area normal to the direction of flow (L/T),

ϕ = total potential (L),

$K(\theta)$ = unsaturated hydraulic conductivity (L/T),

$\theta = V_w/V$ = volumetric moisture content (vol/vol),

V_w = volume of water in the soil mass (L³), and

V = total volume of the soil mass (L³).

The total potential head ϕ is illustrated in Figure 2-2 which shows the negative suction head developed in the unsaturated zone. Figure 2-2 shows that the total potential head is:

$$\phi = -z - \Delta \quad (2-2)$$

For unsaturated soil Δ is the suction head and considering $\Delta = \psi$, the total potential head ϕ is:

$$\phi = -\Psi - z \quad (2-3)$$

The continuity equation can be written as (Willis and Yeh, 1987):

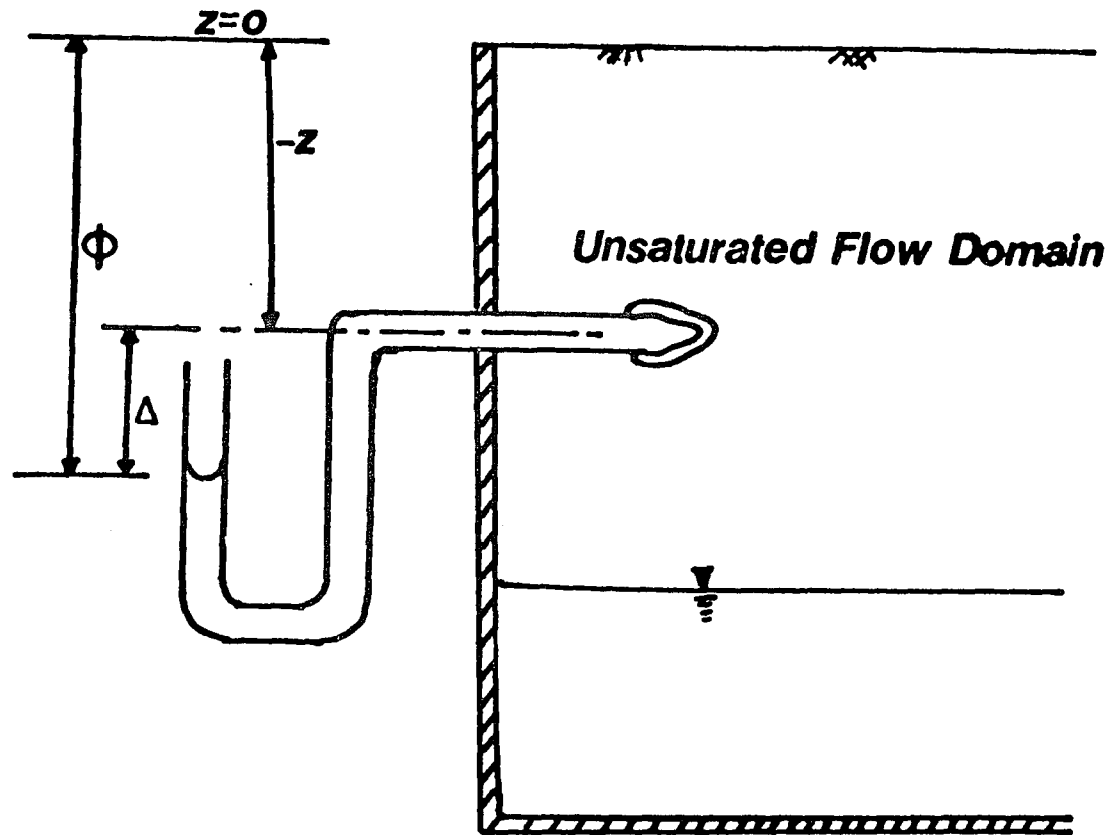


Figure 2-2. Definition Diagram for Total Potential (Bear, 1979).

$$-\nabla \cdot (\rho \mathbf{q}) = \frac{\partial (\rho \theta)}{\partial t} \quad (2-4)$$

where,

ρ = density of water.

Assuming incompressible flow:

$$-\nabla \cdot \mathbf{q} = \frac{\partial \theta}{\partial t} \quad (2-5)$$

From equations (2-1), (2-3), and (2-5):

$$-\nabla \cdot K(\theta) \nabla \psi - \frac{\partial K(\theta)}{\partial z} = \frac{\partial \theta}{\partial t} \quad (2-6)$$

Equation (2-6) can be written as:

$$\frac{\partial \theta}{\partial t} = \frac{\partial}{\partial x} \left[-K(\theta) \frac{\partial \psi}{\partial x} \right] + \frac{\partial}{\partial y} \left[-K(\theta) \frac{\partial \psi}{\partial y} \right] + \frac{\partial}{\partial z} \left[-K(\theta) \frac{\partial \psi}{\partial z} \right] - \frac{\partial K(\theta)}{\partial z} \quad (2-7)$$

Considering a two-dimensional flow domain, equation (2-7) can be written as:

$$\frac{\partial \theta}{\partial t} = \frac{\partial}{\partial x} \left[-K(\theta) \frac{\partial \psi}{\partial x} \right] + \frac{\partial}{\partial z} \left[-K(\theta) \frac{\partial \psi}{\partial z} \right] - \frac{\partial K(\theta)}{\partial z} \quad (2-8)$$

Assuming ψ as the single-valued function of θ , equation (2-8) can be written as:

$$\frac{\partial \theta}{\partial t} = \frac{\partial}{\partial x} \left[-K(\theta) \frac{\partial \psi}{\partial \theta} \frac{\partial \theta}{\partial x} \right] + \frac{\partial}{\partial z} \left[-K(\theta) \frac{\partial \psi}{\partial \theta} \frac{\partial \theta}{\partial z} \right] - \frac{\partial K(\theta)}{\partial z} \quad (2-9)$$

The soil moisture diffusivity $D(\theta)$ can be written as:

$$D(\theta) = -K(\theta) \frac{\partial \psi}{\partial \theta} \quad (2-10)$$

Therefore, the two-dimensional unsaturated moisture flow equation can be written as:

$$\frac{\partial \theta}{\partial t} = \frac{\partial}{\partial x} \left[D(\theta) \frac{\partial \theta}{\partial x} \right] + \frac{\partial}{\partial z} \left[D(\theta) \frac{\partial \theta}{\partial z} \right] - \frac{\partial K(\theta)}{\partial z} \quad (2-11)$$

This equation is solved to obtain the moisture content in the unsaturated zone of a landfill. The solution of moisture content in the unsaturated zone generates the time-history of leachate input rate on the saturated leachate mound.

2.1.1 Hydraulic Conductivity and Diffusivity Functions

The values of the unsaturated hydraulic conductivity vary with the moisture content in the soil. In the unsaturated media, the hydraulic conductivity is expressed in a modified form as a function of the volumetric moisture content. The unsaturated

hydraulic conductivity as a function of moisture content θ can be expressed as (Campbell, 1974):

$$K(\theta) = K_s \left(\frac{\theta}{\theta_s} \right)^{2b+3} \quad (2-12)$$

where,

$K(\theta)$ = unsaturated hydraulic conductivity as a function of moisture content (L/T),

θ_s = saturation moisture content (vol/vol),

K_s = saturated hydraulic conductivity (L/T),

b = slope of moisture characteristic curve plotted on log-log paper and is known as Campbell's parameter.

From equation (2-10), it is clear that to determine the diffusivity coefficient $D(\theta)$, the suction head has to be differentiated with respect to moisture content θ . The suction head ψ as a function of moisture content θ can be expressed as (Bristow and Williams, 1987):

$$\psi(\theta) = \psi_e \left(\frac{\theta}{\theta_s} \right)^{-b} \quad (2-13)$$

where,

$\psi(\theta)$ = suction head as a function of moisture content θ ,

ψ_e = air entry suction head, and

b is the same as defined before.

Reasonable estimates for ψ_e can be obtained from the moisture characteristic curves. In practice, moisture content at air entry $\theta_e = 0.9\theta_s$ and in general, θ_e varies between $0.8\theta_s$ and θ_s (Rogowski, 1971). Therefore, air entry suction head ψ_e corresponds to the suction head value for 90 percent moisture content in the soil matrix.

From equations (2-10), (2-12), and (2-13), the diffusivity relationship in the unsaturated zone can be expressed as:

$$D(\theta) = bK(\theta) \left(\frac{\psi_e}{\theta_s} \right) \left(\frac{\theta_s}{\theta} \right)^{b+1} \quad (2-14)$$

The above relationships for unsaturated diffusivity and hydraulic conductivity are used in the unsaturated moisture-flow equation to compute the moisture content and hence the moisture-flow rate in the unsaturated zone.

2.2 BOUNDARY AND INITIAL CONDITIONS

Boundary conditions define the dependent variable at the boundaries of the domain. In the FILL model, two types of upper boundary conditions are considered, the unsaturated condition and the saturated condition. For the unsaturated flow condition, the moisture-flux is composed of precipitation less evapotranspiration and runoff. This condition at the upper boundary occurs when

net precipitation intensity is less than the saturated hydraulic conductivity (Demetracopoulos et al., 1986). The precipitation intensity that produces runoff is known as the net precipitation intensity. Therefore, the input at the upper boundary for the unsaturated flow condition can be written as:

$$\text{INPUT} = P - \text{ET} - R \quad (2-15)$$

where,

P = precipitation (L),

ET = evapotranspiration (L), and

R = runoff (L).

The saturated boundary condition exists for the upper surface of the landfill when the net precipitation intensity is greater than or equal to the saturated hydraulic conductivity (Demetracopoulos et al., 1986). In this situation, the upper boundary condition is expressed as:

$$\theta = \theta_s \quad (2-16)$$

where,

θ_s = moisture content at saturation.

At the lower boundary, a free-drainage condition is considered, which is expressed as (Korfiatis et al., 1984):

$$\frac{\partial \theta}{\partial x} = 0 \quad (2-17a)$$

$$\frac{\partial \theta}{\partial z} = 0 \quad (2-17b)$$

The initial condition is defined as the known values of moisture content in the landfill section. Therefore, at time $t=0$, the initial water contents are:

$$\theta(x, z, 0) = g(x, z) \quad (2-18)$$

where $g(x,z)$ is the known volumetric water content at all points in the x - z plane.

2.3 RUNOFF

Conservation of mass requires that the change in water level (H) with time, equals the sum of the negative of the change of flow rate with distance and difference of rainfall and infiltration. This can be expressed as [(Brass, 1990; Chow et al., 1988; and Hillel, 1980), Figure 2-3]:

$$\frac{\partial Q}{\partial x} + \frac{\partial H}{\partial t} = \bar{q} \quad (2-19)$$

where,

Q = overland flow rate (L^3/T),

H = depth of runoff (L),

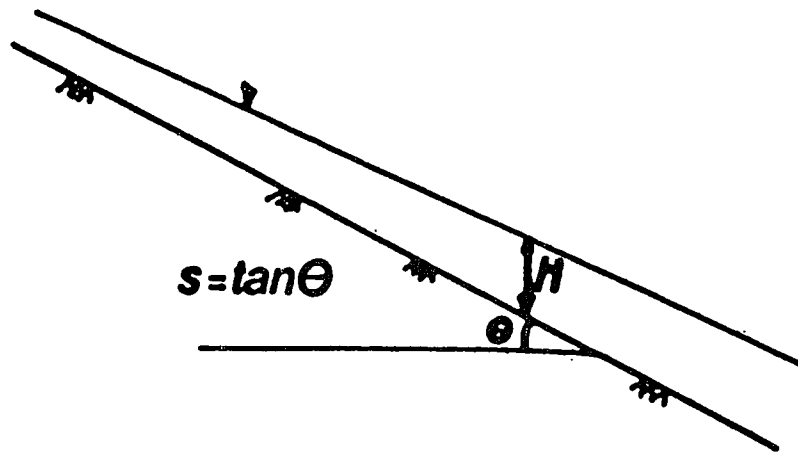


Figure 2-3. A Schematic Representation of Runoff along an Inclined Surface.

$$\bar{q} = P - f$$

P = rainfall rate (L/T), and

f = infiltration rate (L/T).

For overland flow Q is expressed as (Chow et al., 1988):

$$Q = \frac{1.49}{n} H^{\frac{5}{3}} S^{\frac{1}{2}} \quad (2-20)$$

where,

n = Manning's roughness coefficient, and

S = Slope of the surface.

From equation (2-20), the depth of overland flow can be written as:

$$H = \alpha Q^{\beta} \quad (2-21)$$

where,

$$\alpha = \left(\frac{n}{1.49\sqrt{S}} \right)^{\beta}$$

and

$$\beta = 0.6.$$

From equations (2-19) and (2-21), the kinematic wave equation is expressed as

(Chow et al., 1988):

$$\frac{\partial Q}{\partial x} + \alpha \beta Q^{\beta-1} \frac{\partial Q}{\partial t} = \bar{q} \quad (2-22)$$

When capillary forces are dominant, the cumulative infiltration equation can be expressed as (Haverkamp et al., 1987):

$$I(t) = S_o(t)^{1/2} \quad (2-23)$$

From which the infiltration capacity f_c can be expressed as:

$$f_c = \frac{dI}{dt} = \frac{S_o}{2}(t)^{-1/2} \quad (2-24)$$

where,

$I(t)$ = cumulative infiltration per unit surface,

S_o = sorptivity, reflecting the capacity of soil to absorb water by simple capillary forces, and

t = time.

Adjustment in the value of S_o for changing boundary condition has been made by the expression developed by Philip (1969):

$$S_o = [2K_s(\theta_s - \theta_i)(H - h_f)]^{1/2} \quad (2-25)$$

where,

K_s = hydraulic conductivity at saturation (L/T),

θ_s = moisture content at saturation (vol/vol),

θ_i = initial moisture content (vol/vol),

H = depth of water at the soil surface (L), and

h_f = the pressure value defining wetting front, which is constant in space and time (L).

Yu-Si (1987) defined h_f as the capillary head $-h_c$ at the wetting front in the expression for sorptivity S_o . Thus, equation (2-25) can be written as:

$$S_o = [2K_s(\theta_s - \theta_i)(H + h_c)]^{1/2} \quad (2-26)$$

This equation for sorptivity is used to compute the rate of infiltration along the upper boundary of the landfill. The value of the sorptivity from equation (2-26) is used in equation (2-24) at different time steps to compute the infiltration capacity. If rainfall intensity is higher than the infiltration capacity, the infiltration rate is considered the same as the infiltration capacity otherwise the total amount of rainfall enters into the soil surface. The infiltration rate is subtracted from the rainfall rate (as shown in equation 2-19), which is called the rainfall excess and this is the source term for the kinematic wave equation (2-22). In the kinematic wave equation (2-22), the right hand side, i.e. rainfall excess is known and the value of the depth of runoff is computed at different positions on the landfill surface considering the slope and surface roughness.

2.4 EVAPOTRANSPIRATION

The modified Penman method is used to compute potential evapotranspiration as developed by Ritchie (1972) and used in CREAMS (Knisel, 1980). The expression for the potential evapotranspiration is (Schroeder et al., 1984):

$$E_{oi} = \frac{1.28A_iR_i}{(A_i+G)25.4} \quad (2-27)$$

where,

E_{oi} = potential evapotranspiration on day i (inches),

A_i = slope of saturation vapor pressure curve on day i,

R_i = net solar radiation on day i (langleys), and

G = psychrometric constant which is assumed to remain constant at 0.68.

Evapotranspiration of available surface water occurs first. After available surface water is depleted, water from the soil column through evaporation and plant transpiration occurs. Evaporation from surface moisture is given by (Schroeder et al., 1984):

$$ESS_i = \begin{cases} E_{oi} & \text{for } E_{oi} \leq SNO_{i-1} + P_i \\ SNO_{i-1} + P_i & \text{for } E_{oi} > SNO_{i-1} \end{cases} \quad (2-28)$$

where,

ESS_i = surface water evaporation on day i (inches),

SNO_{i-1} = amount of snow at the end of day i-1, and

P_i = actual precipitation on day i (inches).

In the presence of actively transpiring plants, the potential soil evaporation is computed as follows (Schroeder et al., 1984):

$$ES_{oi} = E_{oi} e^{-0.4LAI} \quad (2-29)$$

where,

ES_{oi} = potential soil evaporation on day i (inches),

LAI = leaf area index on day i, and E_{oi} as defined previously.

The leaf area index is a factor that indicates the vegetative cover for a landfill area and is defined as the ratio of the total leaf area to the soil surface area. During the nongrowing or dormant period, the dead or dormant vegetative cover reduces the potential soil evaporation. Then the potential soil evaporation for winter cover is expressed as (Schroeder et al., 1984):

$$ES_{oi} = E_{oi} e^{-0.4WCF} \quad (2-30)$$

where,

WCF = winter cover factor.

Soil evaporation occurs in two stages. In stage-one, soil evaporation is equal to

the potential soil evaporation ES_{0i} . In stage-two, evaporation is controlled by the rate at which water can be transmitted through the soil.

An upper limit is set to control the stage-one soil evaporation. The upper limit to the stage-one evaporation is expressed as (Knisel, 1980):

$$U = \frac{[9 (a_s - 3)^{0.42}]}{25.4} \quad (2-31)$$

where,

U = upper limit for stage-one evaporation (inches) and

a_s = soil transmissivity parameter for evaporation (mm/ $\sqrt{\text{day}}$).

The total of soil evaporation less infiltration is:

$$ES1T_i = \sum_{k=m}^i (ES_k - IN_k) \quad (2-32)$$

where,

ES_k = soil evaporation on day k (inches),

m = last day when $ES1T_i$ equalled zero, and

IN_k = precipitation - surface runoff - surface water evaporation (inches).

When $ES1T_i$ is greater than U , stage-one evaporation stops and stage-two evaporation starts. Stage-two evaporation is controlled by soil transmissivity.

Stage-two evaporation from the soil is estimated as (Knisel, 1980):

$$ES2_i = \frac{1}{25.4} [\sqrt{t_i} - \sqrt{(t_i - 1)}] \quad (2-33)$$

where,

$ES2_i$ = stage-two soil evaporation for day i (inches), and

t_i = days since stage-one evaporation ended.

The potential plant transpiration is computed as (Schroeder et al., 1984):

$$EP_{oi} = \frac{E_{oi} LAI_i}{3.0} \quad (2-34)$$

where,

EP_{oi} = potential plant transpiration on day i (inches).

The actual plant transpiration demand is expressed as (Schroeder et al., 1984):

$$EPD_i = \begin{cases} EP_{oi} & \text{for } EP_{oi} + ESS_i + ES_i \leq E_{oi} \\ E_{oi} - ESS_i - ES_i & \text{for } EP_{oi} + ESS_i + ES_i > E_{oi} \end{cases} \quad (2-35)$$

where,

ES_i = actual soil evaporation on day i (inches).

Using the relationship developed by Shanholtz and Lillard (1970), the actual plant transpiration is computed as:

$$EP_1 = EPD_1 \left[1.2 - 4EPD_1 + \frac{SM_1 - WP}{FC - WP} \right] \quad (2-36)$$

where,

SM_i = soil moisture on day i (inches),

WP = wilting point (VOL/VOL), and

FC = field capacity (VOL/VOL).

Field capacity is defined as the amount of soil moisture content remaining in a unit volume of soil after a prolonged period of gravity drainage. Wilting point is defined as the lowest soil moisture content that can be achieved by plant transpiration.

The total evapotranspiration (ET) is equal to the sum of the evaporation from the soil (ES) and the plant transpiration (EP):

$$ET_1 = \begin{cases} ES_1 + EP_1 & \text{when } (ES_1 + EP_1) \leq (E_{o1} - ESS_1) \\ E_{o1} - ESS_1 & \text{when } (ES_1 + EP_1) > (E_{o1} - ESS_1) \end{cases} \quad (2-37)$$

2.5 SNOWMELT

If the mean daily temperature is below 32°F, the precipitation is stored on the surface as snow. When the temperature exceeds 32°F, snow continues to melt and contributes to the infiltration and/or runoff. The amount of snowmelt is added

to the rainfall for use in equation (2-22). The amount of snowmelt without rain is computed as (Viessman et al., 1977):

$$M = K'(1-F) (0.004I_1) (1-A) + E(0.0084V) \{0.22 (T_a-32) + (0.78 (T_d-32))\} + 0.029F (T-32) \quad (2-38)$$

where,

M = melt, inches/day,

K' = melting coefficient, $0.9 < K' < 1.1$

F = forest cover factor,

I_1 = solar radiation, langley/day,

A = albedo (reflection coefficient),

E = exposure coefficient $0.3 < E < 1.0$,

V = wind velocity (mph),

T_a = air temperature °F, and

T_d = dew point temperature °F.

If there is rain on snow, the equation for snowmelt is (Wanielista, 1990):

$$M = 0.09 + (0.029 + 0.084EV + 0.007P) (T_a - 32) \quad (2-39)$$

where,

V = wind velocity (mph),

P = rainfall (inches per day), and

T_a = air temperature, °F.

2.6 GOVERNING EQUATION IN THE SATURATED ZONE

The leachate head at the bottom of a landfill develops due to the leachate input from the unsaturated zone to the mound. The hydraulics of the leachate mound at the bottom of a landfill has been studied by Korfiatis et al. (1984), Schroeder et al. (1984) and Demetracopoulos (1988). The submodel to compute the leachate head can be formulated by applying Darcy's Law and continuity equation in the saturated zone as shown in Figure 2-4. According to Darcy's Law, the lateral drainage in the saturated zone for an inclined bottom can be written as (Bear, 1972):

$$q_L = -K_s h \frac{\partial z}{\partial x} \quad (2-40)$$

where,

K_s = saturated hydraulic conductivity (L/T),

h = saturated thickness of the landfill section (L), and

z = elevation of water table from a reference datum (L).

Figure 2-4 shows that:

$$z = h + d$$

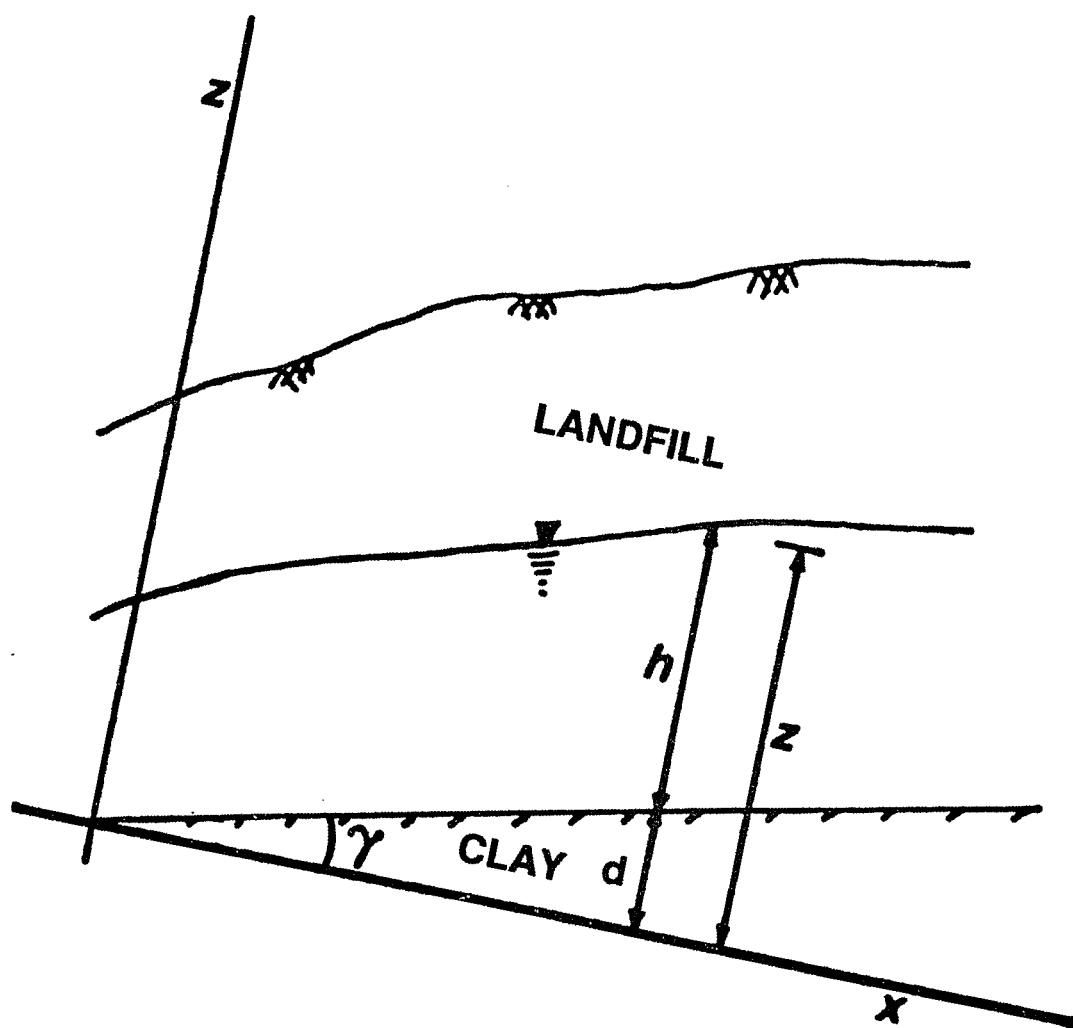


Figure 2-4. Saturated Flow at the Bottom of a Landfill.

$$\text{or, } z = h + x \tan \gamma$$

From which:

$$\frac{\partial z}{\partial x} = \frac{\partial h}{\partial x} + \tan \gamma$$

Therefore, the lateral drainage can be written as:

$$q_L = -K_s h \left[\frac{\partial h}{\partial x} + \tan \gamma \right] \quad (2-41)$$

Assuming a fully saturated bottom liner, the leakage rate can be expressed as (Korfiatis and Demetracopoulos, 1986):

$$q_v = K_c \left(\frac{h + d}{d} \right) \quad (2-42)$$

where,

K_c = saturated hydraulic conductivity of the clay liner (L/T), and

d = thickness of clay at the bottom of the landfill (L).

The mass conservation principle states that the sum of all flow rates into the control volume must equal the rate of volume change in the control volume. This yields (Korfiatis and Demetracopoulos, 1986):

$$-\frac{\partial q_L}{\partial x} + N_l - K_c \left(\frac{h+d}{d} \right) = n_e \frac{\partial h}{\partial t} \quad (2-43)$$

where,

N_l = leachate accretion onto the drainage layer (L/T),

n_e = effective porosity of the solid waste material,

Substituting equation (2-41) into equation (2-43):

$$\frac{\partial}{\partial x} \left[K_s h \left(\frac{\partial h}{\partial x} + \tan \gamma \right) \right] + N_l - K_c \left(\frac{h+d}{d} \right) = n_e \frac{\partial h}{\partial t} \quad (2-44)$$

Equation (2-44) is solved to obtain the leachate head at the bottom of the landfill which is used to compute lateral and vertical flow of leachate from the landfill to the perimeter or collection system.

After solving equation (2-11) for soil moisture content, the leachate accretion or drainage to the saturated leachate mound is computed. The governing equation for the leachate accretion is expressed as (Korfiatis et al., 1984):

$$N_l = K(\theta) - D(\theta) \frac{\partial \theta}{\partial z} \quad (2-45)$$

The major uncertainties in making more accurate estimate of leachate accretion to the saturated leachate mound are in the material properties such as the

coefficients related to the soil moisture diffusivity $D(\theta)$ and hydraulic conductivity $K(\theta)$. For this reason, after computing the soil moisture content, the water balance is checked for change in the moisture contents by considering the net infiltration (at the boundary), inflow and outflow due to the diffusivity and the unsaturated hydraulic conductivity in the domain. The drainage rate to the saturated leachate mound computed by the governing equation (2-45) is compared with the drainage rates obtained by the water balance in the domain and by using empirical formulas described by Bras (1990). An appropriate value is estimated by assuming that the recharge to the saturated leachate mound can not exceed the total net infiltration over a period of time.

A common empirical form of the percolation rate is expressed as (Bras, 1990):

$$N_d = a_s \theta^{b_s} \quad (2-46)$$

where,

$$a_s = K_s \psi_e$$

b_s = a calibration parameter.

If water content is between field capacity and saturation, the soil moisture drainage rate is considered as (Bras, 1990):

$$N_d = f_c \left(1 - \frac{\theta_s - \theta_i}{\theta_s - \theta_c} \right)^3 \quad (2-47)$$

where,

f_c = minimum infiltration rate (L/T),

θ_i = initial soil moisture content (VOL/VOL), and

θ_c = field capacity moisture content (VOL/VOL).

CHAPTER 3

FINITE-DIFFERENCE FORMULATION

The governing partial differential equations incorporated into the two-dimensional leachate flow model, represent runoff at the upper boundary, moisture-flow in the unsaturated zone along with the boundary conditions, and movement of leachate at the bottom of a landfill. All these equations are discretized using the finite-difference technique to obtain a system of algebraic equations. In this method, the derivatives are approximated in the finite grid points to transform the partial differential equation into a set of simultaneous algebraic equations. These equations are solved simultaneously for the dependent variables using numerical solution methods. Application of the finite-difference techniques is described for each of the governing partial differential equations in this chapter.

3.1 UNSATURATED MOISTURE-FLOW EQUATION

In the finite-difference method, the unsaturated moisture-flow equation (2-11) is discretized in x - z - t space. The horizontal and vertical lines in x - and z -directions are divided into segments of lengths of Δx and Δz respectively. These constitute rectangular nodal areas in the vertical plane having a nodal point at the centroid

of the rectangles as shown in Figure 3-1. The time is divided into intervals of duration Δt . An implicit finite-difference scheme was employed by Demetracopoulos et al. (1986) to discretize the one-dimensional unsaturated moisture-flow equation. In the present model, the same concept of the finite-difference scheme is used to discretize the two-dimensional unsaturated moisture-flow equation (2-11). According to Figure 3-1, discretization of equation (2-11) for the (i,j)th grid point yields:

$$\begin{aligned} \frac{\theta_{i,j}^{k+1} - \theta_{i,j}^k}{\Delta t} = & \frac{1}{(\Delta x)^2} [D_{i+1/2,j}(\theta_{i+1,j} - \theta_{i,j}) - D_{i-1/2,j}(\theta_{i,j} - \theta_{i-1,j})] \\ & + \frac{1}{(\Delta z)^2} [D_{i,j+1/2}(\theta_{i,j+1} - \theta_{i,j}) - D_{i,j-1/2}(\theta_{i,j} - \theta_{i,j-1})] \\ & - \frac{1}{\Delta z} (K_{i,j+1/2} - K_{i,j-1/2}) \end{aligned} \quad (3-1)$$

where,

k = initial time level;

$k+1$ = forward time level;

$D_{i+1/2,j}$ = diffusivity between nodes i,j and $i+1,j$;

$D_{i-1/2,j}$ = diffusivity between nodes i,j and $i-1,j$;

$D_{i,j+1/2}$ = diffusivity between nodes i,j and $i,j+1$;

$D_{i,j-1/2}$ = diffusivity between nodes i,j and $i,j-1$;

$K_{i,j+1/2}$ = unsaturated hydraulic conductivity between nodes i,j and $i,j+1$; and

$K_{i,j-1/2}$ = unsaturated hydraulic conductivity between nodes i,j and $i,j-1$.

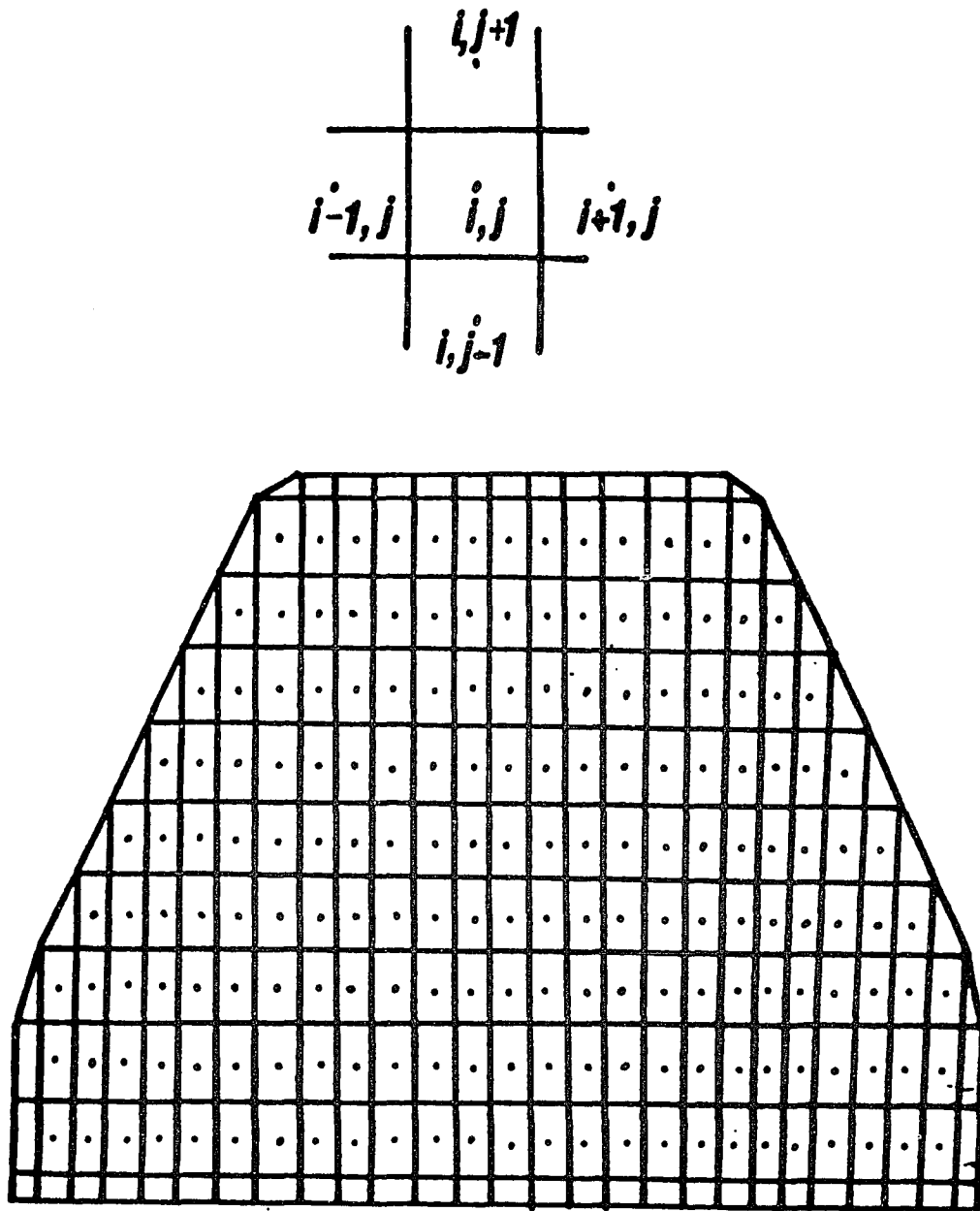


Figure 3-1. Finite-difference Grid Representation of a Landfill Section.

The dependent variable θ with the subscripts i,j ; $i+1,j$; $i-1,j$; $i,j+1$; and $i,j-1$ defines the moisture content in the finite-difference grid-points. The diffusivity D , and hydraulic conductivity K in equation (3-1) are functions of the moisture content θ . The values of $D_{i+1/2,j}$, $D_{i-1/2,j}$, $D_{i,j+1/2}$, $D_{i,j-1/2}$, $K_{i,j+1/2}$, and $K_{i,j-1/2}$ at a particular time step are determined by the extrapolation from the previous time steps as described by Rubin and Steinhardt (1963). By evaluating $\theta_{i+1,j}$, $\theta_{i-1,j}$, $\theta_{i,j+1}$, and $\theta_{i,j-1}$ at the forward time level, the implicit finite-difference equation is obtained as follows:

$$\begin{aligned} \frac{\theta_{i,j}^{k+1} - \theta_{i,j}^k}{\Delta t} &= \frac{1}{(\Delta x)^2} [D_{i+1/2,j}(\theta_{i+1,j}^{k+1} - \theta_{i,j}^{k+1}) - D_{i-1/2,j}(\theta_{i,j}^{k+1} - \theta_{i-1,j}^{k+1})] \\ &+ \frac{1}{(\Delta z)^2} [D_{i,j+1/2}(\theta_{i,j+1}^{k+1} - \theta_{i,j}^{k+1}) - D_{i,j-1/2}(\theta_{i,j}^{k+1} - \theta_{i,j-1}^{k+1})] \\ &- \frac{1}{\Delta z} (K_{i,j+1/2} - K_{i,j-1/2}) \end{aligned} \quad (3-2)$$

It is difficult to obtain a regular network of grid points having equal space intervals Δx and Δz throughout the domain of a landfill with irregular boundaries. In this case, it is necessary to develop special finite-difference formulas for irregular spacing of the grid points. Following the Taylor series expansion described by Remson et al. (1971), the implicit finite-difference expression for irregular grid points is obtained with respect to Figure 3-2 as:

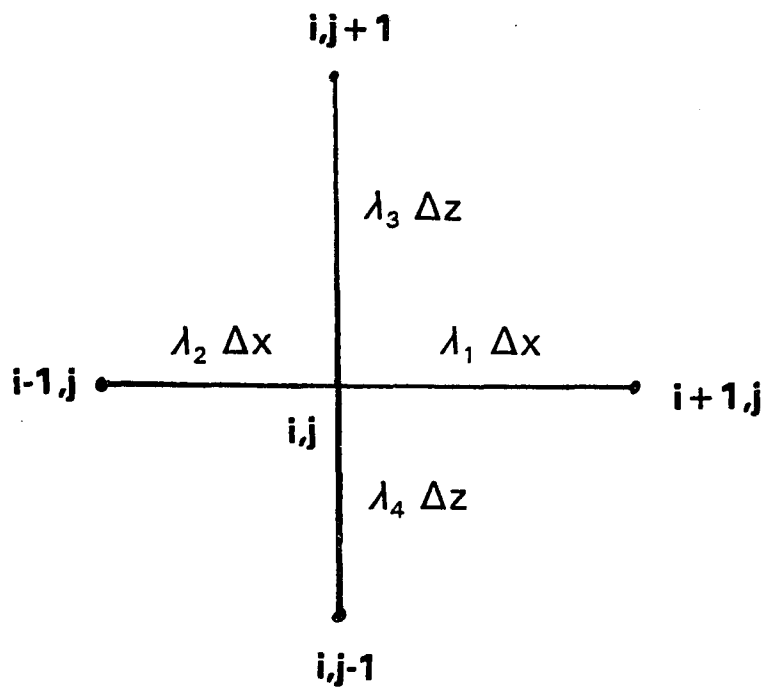


Figure 3-2. Irregular Finite-difference Grid Spacing.

$$\begin{aligned}
\frac{\theta_{i,j}^{k+1} - \theta_{i,j}^k}{\Delta t} &= \frac{2}{(\lambda_1 + \lambda_2) \Delta x} \left[D_{i+1/2,j} \frac{(\theta_{i+1,j}^{k+1} - \theta_{i,j}^{k+1})}{\lambda_1 \Delta x} - D_{i-1/2,j} \frac{(\theta_{i,j}^{k+1} - \theta_{i-1,j}^{k+1})}{\lambda_2 \Delta x} \right] \\
&+ \frac{2}{(\lambda_3 + \lambda_4) \Delta z} \left[D_{i,j+1/2} \frac{(\theta_{i,j+1}^{k+1} - \theta_{i,j}^{k+1})}{\lambda_3 \Delta z} - D_{i,j-1/2} \frac{(\theta_{i,j}^{k+1} - \theta_{i,j-1}^{k+1})}{\lambda_4 \Delta z} \right] \\
&- 2 \left(\frac{K_{i,j+1/2} - K_{i,j-1/2}}{\lambda_3 \Delta z + \lambda_4 \Delta z} \right) \tag{3-3}
\end{aligned}$$

The system of finite-difference expression developed, can be solved by the Gauss-Seidel iteration method. The iteration procedure is continued until two consecutive sets of moisture values agree within a prescribed tolerance criterion.

3.2 BOUNDARY CONDITIONS

The upper boundary of a landfill includes negligible slope at the top and steep slopes along the sides. The upper boundary condition for the boundaries at the top and at the side slope for unsaturated moisture-flow condition can be described by mathematical expression and then a finite-difference technique can be applied to get a system of algebraic equations to compute moisture values at the surface.

The upper boundary conditions at the side slope are defined by considering the flow of moisture from the boundary node (i,j) to the adjacent nodes (i+1,j) and (i,j+1) in both the horizontal and vertical directions respectively as shown in Figure 3-3.

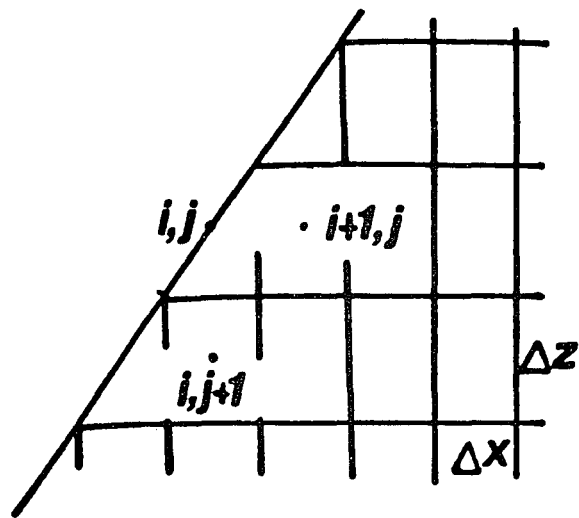


Figure 3-3. Boundary Node along the Upper Surface of a Landfill.

The unsaturated moisture flow equation (2-5) can be written as:

$$-\frac{\partial q}{\partial x} - \frac{\partial q}{\partial z} = \frac{\partial \theta}{\partial t} \quad (3-4)$$

Moisture flow in x-direction (q_x) can be written as:

$$q_x = -D(\theta) \frac{\partial \theta}{\partial x} \quad (3-5)$$

where,

$$D(\theta) = -K(\theta) \frac{\partial \psi}{\partial \theta}$$

ψ = suction head (L),

$D(\theta)$ = diffusivity (L^2/T), and

$K(\theta)$ = unsaturated hydraulic conductivity (L/T).

Equation (3-4) can be written for the boundary node (i,j) as:

$$-\frac{q_{i+1,j} - q_{i,j}}{\Delta x} - \frac{q_{i,j+1} - q_{i,j}}{\Delta z} = \frac{\theta_{i,j}^{k+1} - \theta_{i,j}^k}{\Delta t} \quad (3-6)$$

where,

$q_{i,j}$ = rainfall rate - (evapotranspiration rate + runoff rate)

= net infiltration (L/T).

From equation (3-5), the finite-difference expression $q_{i+1,j}$ can be written as:

$$q_{i+1,j} = - D_{i+1/2,j} \left(\frac{\theta_{i+1,j} - \theta_{i,j}}{\Delta x} \right) \quad (3-7)$$

From equation (2-1), the specific discharge in the z-direction in the unsaturated zone is:

$$q_z = K_z(\theta) - D(\theta) \frac{\partial \theta}{\partial z} \quad (3-8)$$

From which:

$$q_{i,j+1} = K_{i,j+1/2} - D_{i,j+1/2} \left(\frac{\theta_{i,j+1} - \theta_{i,j}}{\Delta z} \right) \quad (3-9)$$

Substituting (3-7) and (3-9) in (3-6), the implicit finite-difference expression for the unsaturated moisture-flow equation at the upper boundary is obtained as:

$$\begin{aligned} - \frac{\theta_{i,j}^{k+1} - \theta_{i,j}^k}{\Delta t} &= - \frac{1}{\Delta z} \left[K_{i,j+1/2} - D_{i,j+1/2} \left(\frac{\theta_{i,j+1}^{k+1} - \theta_{i,j}^{k+1}}{\Delta z} \right) \right] + \frac{q_{i,j}^{k+1}}{\Delta z} \\ &\quad - \frac{1}{\Delta x} \left[- D_{i+1/2,j} \left(\frac{\theta_{i+1,j}^{k+1} - \theta_{i,j}^{k+1}}{\Delta x} \right) \right] + \frac{q_{i,j}^{k+1}}{\Delta x} \end{aligned} \quad (3-10)$$

The implicit finite-difference expression (3-10) is solved to obtain the moisture content along the upper boundary which depends on precipitation less evapotranspiration and surface runoff. The flow at the top flat surface of a landfill is considered predominantly vertical. Similar finite-difference expressions are

developed for top flat surface considering moisture movement only vertically from the upper to the lower node.

At the bottom of the unsaturated zone in a landfill, where the saturated leachate mound exists, the moisture gradient along the lower boundary is postulated herein as:

$$\frac{\partial \theta}{\partial x} = 0 \quad (3-11a)$$

$$\frac{\partial \theta}{\partial z} = 0 \quad (3-11b)$$

This type of postulation has been used successfully for one-dimensional moisture flow equations (Korfiatis et al., 1984), which in fact implies free drainage from the bottom due to gravity.

3.3 RUNOFF EQUATION

The overland flow equation or the kinematic wave equation (2-22) is transformed into a system of algebraic expressions using a finite-difference scheme. The finite-difference scheme is explained for two-dimensional grid in x-t plane as shown in Figure 3-4. Different finite-difference expressions yield different numerical schemes associated with their own stability and convergence criteria. Chow et al. (1988) described the finite-difference scheme which is stable and convergent. The finite-

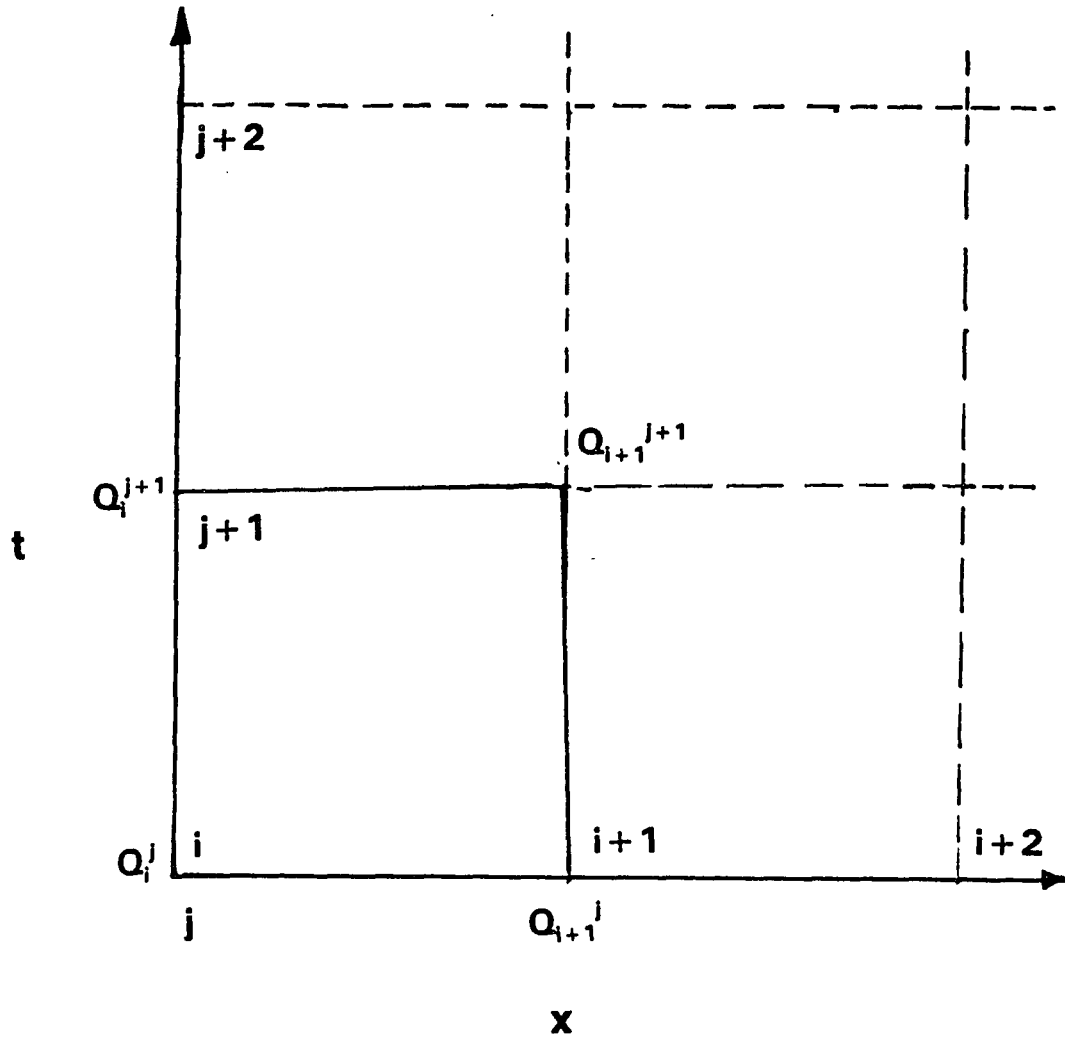


Figure 3-4. Finite-difference Scheme for the Kinematic Wave Equation.

difference expression with respect to Figure 3-4 is:

$$\frac{Q_{i+1}^{k+1} - Q_i^{k+1}}{\Delta x} + \alpha\beta \left(\frac{Q_{i+1}^k + Q_i^{k+1}}{2} \right)^{\beta-1} \left(\frac{Q_{i+1}^{k+1} - Q_i^k}{\Delta t} \right) = \bar{q} \quad (3-12)$$

The unknown variable Q_{i+1}^{k+1} is then obtained as follows:

$$Q_{i+1}^{k+1} = \frac{\left[\frac{\Delta t}{\Delta x} Q_i^{k+1} + \alpha\beta Q_{i+1}^k \left(\frac{Q_{i+1}^k + Q_i^{k+1}}{2} \right)^{\beta-1} + \Delta t \bar{q} \right]}{\left[\frac{\Delta t}{\Delta x} + \alpha\beta \left(\frac{Q_{i+1}^k + Q_i^{k+1}}{2} \right)^{\beta-1} \right]} \quad (3-13)$$

The solution procedure at each segment consists of solving for the flow rate Q at time t and then advancing the solution to forward time level $t + \Delta t$, considering the values of Q at the previous time level t as the initial condition.

3.4 SATURATED LEACHATE FLOW EQUATION

The finite-difference expression for the saturated leachate flow equation (2-44) can be obtained by discretizing it in the finite-difference grid as shown in Figure 3-5.

The transformed finite-difference equation can be written as:

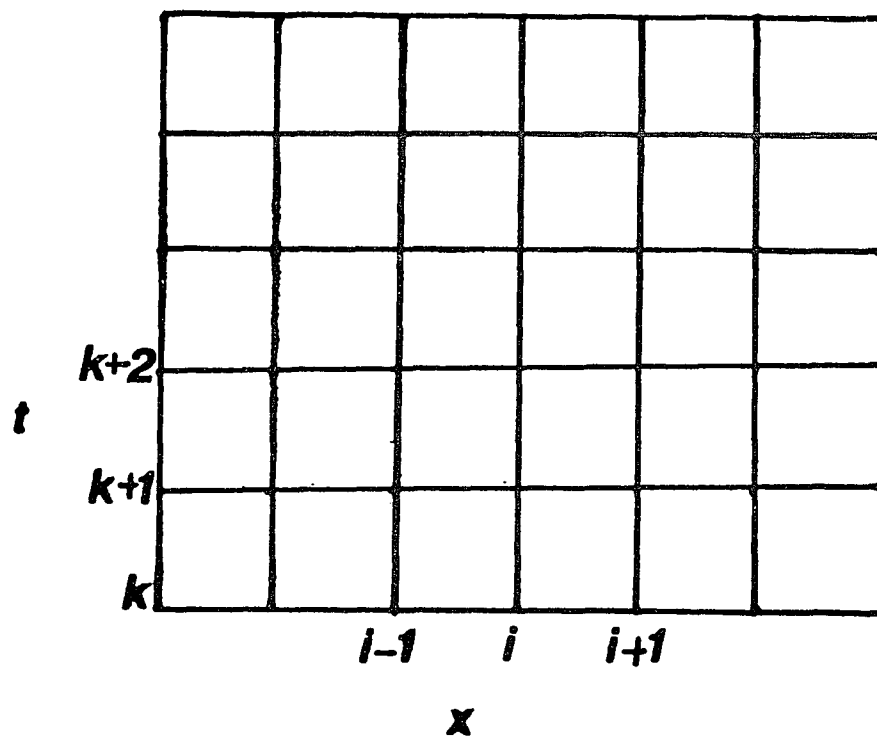


Figure 3-5. Finite-difference Scheme for the Saturated Leachate Flow Equation.

$$\begin{aligned} & \frac{1}{(\Delta x)^2} [T_{i+1/2}(h_{i+1} - h_i) - T_{i-1/2}(h_i - h_{i-1})] - K_{si} \tan \gamma \left(\frac{h_{i+1} - h_{i-1}}{2\Delta x} \right) \\ & + N_{1i} - K_{ci} \left(\frac{h_i + d_i}{d_i} \right) = n_{ei} \left(\frac{h_i^{k+1} - h_i^k}{\Delta t} \right) \end{aligned} \quad (3-14)$$

This system of finite-difference equations can be solved simultaneously for each grid-point by evaluating all the dependent variables at the forward time level, $k+1$.

Therefore, the implicit expression for the equation (3-14) can be written as:

$$\begin{aligned} & \frac{1}{(\Delta x)^2} [T_{i+1/2}(h_{i+1}^{k+1} - h_i^{k+1}) - T_{i-1/2}(h_i^{k+1} - h_{i-1}^{k+1})] - K_{si} \tan \gamma \left(\frac{h_{i+1}^{k+1} - h_{i-1}^{k+1}}{2\Delta x} \right) \\ & + N_{1i} - K_{ci} \left(\frac{h_i^{k+1} + d_i}{d_i} \right) = n_{ei} \left(\frac{h_i^{k+1} - h_i^k}{\Delta t} \right) \end{aligned} \quad (3-15)$$

The leachate accretion N_{1i} is obtained after computing moisture values which helps build up the saturated leachate mound head. The values of leachate mound head h at each grid point are the only unknowns in equation (3-15) and are solved simultaneously at each time step using Gauss-Jordan elimination method.

3.5 NUMERICAL SOLUTION FOR UNSATURATED FLOW EQUATION

The system of finite-difference equations (3-2) derived from the partial differential

equations describing moisture-flow in the unsaturated zone is solved by applying Gauss-Seidel iteration technique. The finite-difference expression (3-15) for the saturated leachate flow equation is solved using Gauss-Jordan elimination method. The iteration method is not used for the saturated leachate flow equation in the FILL model because for large values of leachate accretion from the unsaturated zone, the solution for the system of equations (3-15) does not converge rapidly. In the elimination method, the solution is obtained directly and no iteration is required.

In the unsaturated zone, iteration method is suitable to avoid storing large coefficient matrix for grid-points in the two-dimensional vertical domain. In the saturated zone, the coefficient matrix is not large and elimination method is used in the main model to avoid large central processing unit time required for iteration. However, for comparison with the analytical solution, both iteration and elimination methods are used for the solution of the finite-difference equations describing flow in both the unsaturated and saturated zones of a landfill. The two methods are used and compared with the analytical solution to establish the credibility of the model results.

To solve the system of algebraic equations for moisture content in the unsaturated zone, the relaxation scheme (Thomas, 1973) of the Gauss-Seidel iteration method is used. The values of the moisture content θ are assumed and the improved values are then calculated with the known initial and boundary conditions. A

residual flow is caused by the discrepancy between the exact solution and the approximate values of the moisture content. This is gradually diminished by the process of relaxation and more accurate values of moisture content are obtained.

From equation (3-2) the residual value can be expressed as:

$$\begin{aligned}
 \text{RES}_{i,j} = & -\frac{1}{\Delta t}(\theta_{i,j}^{k+1} - \theta_{i,j}^k) + \frac{D_{i+1/2,j}}{(\Delta x)^2}(\theta_{i+1,j}^{k+1} - \theta_{i,j}^{k+1}) \\
 & - \frac{D_{i-1/2,j}}{(\Delta x)^2}(\theta_{i,j}^{k+1} - \theta_{i-1,j}^{k+1}) + \frac{D_{i,j+1/2}}{(\Delta z)^2}(\theta_{i,j+1}^{k+1} - \theta_{i,j}^{k+1}) \\
 & - \frac{D_{i,j-1/2}}{(\Delta z)^2}(\theta_{i,j}^{k+1} - \theta_{i,j-1}^{k+1}) - \frac{1}{\Delta z}(K_{i,j+1/2} - K_{i,j-1/2}) \quad (3-16)
 \end{aligned}$$

Adding $\Delta\theta_{i,j}^{k+1}$ to $\theta_{i,j}^{k+1}$ in equation (3-16):

$$\begin{aligned}
 \frac{\theta_{i,j}^{k+1} + \Delta\theta_{i,j}^{k+1} - \theta_{i,j}^k}{\Delta t} = & \frac{D_{i+1/2,j}}{(\Delta x)^2}(\theta_{i+1,j}^{k+1} - \theta_{i,j}^{k+1} - \Delta\theta_{i,j}^{k+1}) \\
 & - \frac{D_{i-1/2,j}}{(\Delta x)^2}(\theta_{i,j}^{k+1} + \Delta\theta_{i,j}^{k+1} - \theta_{i-1,j}^{k+1}) \\
 & + \frac{D_{i,j+1/2}}{(\Delta z)^2}(\theta_{i,j+1}^{k+1} - \theta_{i,j}^{k+1} - \Delta\theta_{i,j}^{k+1}) \\
 & - \frac{D_{i,j-1/2}}{(\Delta z)^2}(\theta_{i,j}^{k+1} + \Delta\theta_{i,j}^{k+1} - \theta_{i,j-1}^{k+1}) \\
 & - \frac{1}{(\Delta z)}(K_{i,j+1/2} - K_{i,j-1/2}) \quad (3-17)
 \end{aligned}$$

From which:

$$\Delta\theta_{i,j}^{k+1} = \frac{\text{RES}_{i,j}}{\frac{1}{\Delta t} + \frac{1}{(\Delta x)^2}(D_{i+1/2,j} + D_{i-1/2,j}) + \frac{1}{(\Delta z)^2}(D_{i,j+1/2} + D_{i,j-1/2})} \quad (3-18)$$

The coefficient of residual is termed relaxation coefficient. Therefore,

$$\text{RELAX}_{i,j} = \frac{1.0}{\frac{1}{\Delta t} + \frac{1}{(\Delta x)^2}(D_{i+1/2,j} + D_{i-1/2,j}) + \frac{1}{(\Delta z)^2}(D_{i,j+1/2} + D_{i,j-1/2})} \quad (3-19)$$

Thus, the moisture content at grid-point (i,j) for the new time step $t = k + 1$, is:

$$\theta_{i,j}(\text{new}) = \theta_{i,j}(\text{old}) + \Delta\theta_{i,j}(\text{new}) \quad (3-20)$$

which is:

$$\theta_{i,j}(\text{new}) = \theta_{i,j}(\text{old}) + \text{RELAX}_{i,j} * \text{RES}_{i,j} \quad (3-21)$$

Successive application of this iteration technique is continued until two consecutive sets of values agree within a prescribed tolerance criterion. A value of 10^{-5} is considered as the tolerance criterion for the convergence of the solution at a finite-difference grid-point.

3.6 NUMERICAL SOLUTION FOR SATURATED FLOW EQUATION

Finite-difference expressions (3-15) for the saturated leachate flow equation (2-44) are solved simultaneously using the Gauss-Jordan elimination method. Equation (3-15) can be rearranged as:

$$\tilde{b}h_i^{k+1} + \tilde{c}h_{i-1}^{k+1} + \tilde{d}h_{i+1}^{k+1} = \tilde{f} \quad (3-22)$$

where,

$$\tilde{b} = -\frac{1}{(\Delta x)^2} [T_{i+1/2} + T_{i-1/2}] - \frac{K_c}{d} - \frac{K_c}{\Delta t}$$

$$\tilde{c} = \frac{T_{i-1/2}}{(\Delta x)^2} - \frac{K_s \tan \gamma}{2\Delta x}$$

$$\tilde{d} = \frac{T_{i+1/2}}{(\Delta x)^2} - \frac{K_s \tan \gamma}{2\Delta x}$$

$$\tilde{f} = -N_1 + K_c - n_e \frac{h_i^k}{\Delta t}$$

For n number of nodes, n number of equations of the type (3-22) are obtained.

They may be expressed in matrix form as:

$$[A] [h] = [F] \quad (3-23)$$

where,

$$[A] = \begin{bmatrix} a_{11} & a_{12} & a_{13} & \cdot & \cdot & \cdot & \cdot & a_{1n} \\ a_{21} & a_{22} & a_{23} & \cdot & \cdot & \cdot & \cdot & a_{2n} \\ a_{31} & a_{32} & a_{33} & \cdot & \cdot & \cdot & \cdot & a_{3n} \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ a_{n1} & a_{n2} & a_{n3} & \cdot & \cdot & \cdot & \cdot & a_{nn} \end{bmatrix}$$

$$[h] = \begin{bmatrix} h_1^{k+1} \\ h_2^{k+1} \\ h_3^{k+1} \\ \cdot \\ \cdot \\ \cdot \\ h_n^{k+1} \end{bmatrix} \quad [F] = \begin{bmatrix} f_1 \\ f_2 \\ f_3 \\ \cdot \\ \cdot \\ \cdot \\ f_n \end{bmatrix}$$

In using Gauss-Jordan elimination method, the coefficient matrix in equation (3-23) is reduced to a diagonal matrix. In this process, the i th row is used to eliminate the coefficients above and below it in the i th column. Thus at the end of the solution, the diagonal elements of $[A]$ become unity with zero values in the remaining elements. The elimination process to make zero in the i th column of $[A]$ with the corresponding operations on the right-hand side column vector $[F]$ can be given as follows:

$$a_{i,j}^{(i+1)} = a_{i,j}^{(i)} / a_{i,i}^{(i)}, \quad j=i,n$$

$$a_{k,m}^{(i+1)} = a_{k,m}^{(i)} - a_{k,i}^{(i)} * a_{i,m}^{(i)} / a_{i,i}^{(i)}$$

$$k=1,n; \quad m=i,n; \quad k \neq i.$$

$$f_i^{(i+1)} = f_i^{(i)} / a_{i,i}^{(i)}$$

$$f_k^{(i+1)} = f_k^{(i)} - a_{k,i}^{(i)} * f_i^{(i)} / a_{i,i}^{(i)}, \quad k=1,n \text{ and } k \neq i.$$

Rearrangement of rows and columns, called pivoting is performed when a diagonal element is nearly zero to ensure numerical stability (Dahlquist and Bjorck, 1974). The set of equations (3-22) are checked whether they retain large values of the diagonal elements at each step of the elimination.

CHAPTER 4

BOUNDARY INTEGRAL FORMULATION

The boundary integral formulations presented in this chapter are derived using the time-dependent fundamental solution of the governing partial differential equation. Two boundary integral approaches are presented. The first approach solves the unsaturated moisture-flow problem by directly applying the time-dependent Green's function. The second method is based on the expansion of the moisture content into perturbation series. The governing equation is decomposed into a moisture-flow equation without the gravity term on the right-hand side and a sequence of moisture flow equation with known gravity term on the right-hand side. The perturbation equations are then solved in succession by evaluating the gravity term of the unsaturated leachate flow equation from the preceding level of solution. The two boundary integral approaches, direct Green's function and perturbation Green's function, are closely related because the same time-dependent fundamental solution is used to formulate the boundary integral expressions. The direct Green's function formulation solves the boundary integral expression in one step. However, the perturbation Green's function approach requires repeated solution of the integral expression for homogeneous boundary conditions.

4.1 DIRECT GREEN'S FUNCTION (DGF) FORMULATION

The present boundary integral procedure for the time-dependent moisture-flow problem is similar to the technique employed by Taigbenu (1985) that was used to solve the unsteady aquifer flow problems. The method uses the time-dependent Green's function which satisfies the diffusion equation. Liggett and Liu (1979) used this method to predict the hydraulic head in a confined aquifer. To apply this method to unsaturated moisture-flow equation (2-11), it is necessary to linearize it to the form of the diffusion equation. Rubin and Steinhardt (1963) suggested that the linearization of the moisture flow equation can be done by computing the unsaturated moisture-flow characteristics by using the moisture content from the previous time steps. In this case, the linearization is also done by considering the diffusivity and the gravity term to be dependent on the moisture content computed from the previous time step.

The linearized form of the unsteady moisture-flow equation (2-11) can be written as:

$$\frac{\partial \theta}{\partial t} = \hat{D} \nabla^2 \theta - \frac{\partial K(\theta^k)}{\partial z} \quad (4-1)$$

where,

$$\hat{D} = D(\theta^k) \text{ and}$$

θ^k = moisture content θ at previous time step $t=k$.

The free-space Green's function for the present problem is used that satisfies the following equation:

$$\frac{\partial G}{\partial t} = \hat{D} \nabla^2 G + \delta(r-r_i, t-\tau) \quad (4-2)$$

where,

$r-r_i$ = distance between any field point (x,z) and the source point (x_i,z_i) , and

τ = a time variable.

Morse and Feshbach (1953) gives the solution for Green's function as:

$$G(r, t) = \frac{1}{4\pi\hat{D}(t-\tau)} \exp\left[-\frac{r^2}{4\hat{D}(t-\tau)}\right] \quad (4-3)$$

Following the direct boundary integral formulation, which employs the Green's identity, the expression (4-1) can be written in the integral form as (Liggett and Liu, 1979):

$$\lambda\theta = \hat{D} \int_{\Gamma} \int_{t_0}^{t_1} \left[G \frac{\partial \theta}{\partial n} - \theta \frac{\partial G}{\partial n} \right] d\tau d\Gamma + \int_{\Omega} \theta_0 G d\Omega - \int_{\Omega} \int_{t_0}^{t_1} G \frac{\partial K(\theta^k)}{\partial z} d\tau d\Omega \quad (4-4)$$

where,

Ω is the domain and Γ is the boundary,

θ_0 = initial moisture content,

$\lambda = 1$ if $(x,z) \in \Omega$,

$\lambda = 0.5$ if $(x,z) \in \Gamma$ and Γ is smooth at (x,z) , and

$n =$ outward unit normal at the boundary.

Some of the data for θ and $\partial\theta/\partial n$ are unknown on the boundary which are obtained by solving boundary integral equation (4-4). After solution of the boundary integral equation (4-4), the values of θ and $\partial\theta/\partial n$ are known everywhere on the boundary. Once these values at the boundary are known, the solution for θ in the domain can be obtained by writing the discretized boundary integral expression for an interior node i as:

$$\begin{aligned} \theta_i = & \sum_{j=1}^N \left\{ \left(\frac{\partial\theta}{\partial n} \right)_j \int_{\Gamma_j} \int_{t_0}^{t_1} G d\tau d\Gamma_j - \theta_j \int_{\Gamma_j} \int_{t_0}^{t_1} \frac{\partial G}{\partial n} d\tau d\Gamma_j \right\} \\ & + \int_{\Omega} \theta_0 G d\Omega - \int_{\Omega} \int_{t_0}^{t_1} G \frac{\partial K(\theta^k)}{\partial z} d\tau d\Omega \end{aligned} \quad (4-5)$$

where,

$N =$ total number of boundary elements.

4.2 TIME INTEGRATION

In time integration, it is assumed that the function θ and $\partial\theta/\partial n$ remain constant on time. The time integration is performed stepwise in which the function G and $\partial G/\partial n$

vary with respect to time. The expression for $\partial G/\partial n$ is obtained by differentiating equation (4-3) as follows:

$$\frac{\partial G}{\partial n} = -\frac{1}{8\pi\hat{D}^2(t_1-\tau)^2} \exp\left[-\frac{r^2}{4\hat{D}(t_1-\tau)}\right] r \frac{\partial r}{\partial n}$$

According to the local coordinate system as shown in Figure 4-1:

$$\frac{\partial r}{\partial n} = \frac{\partial r}{\partial \eta}$$

$$r = \sqrt{\eta^2 + \xi^2}$$

$$\frac{\partial r}{\partial \eta} = \frac{\eta}{r}$$

Therefore,

$$\frac{\partial r}{\partial n} = \frac{\eta}{r}$$

and

$$\frac{\partial G}{\partial n} = -\frac{\eta}{8\pi\hat{D}^2(t_1-\tau)^2} \exp\left[-\frac{r^2}{4\hat{D}(t_1-\tau)}\right] \quad (4-6)$$

Now, the integrals are evaluated in time as follows (Brebbia et al., 1984):

$$\int_{t_0}^{t_1} G d\tau = \frac{1}{4\pi\hat{D}} E(a) \quad (4-7)$$

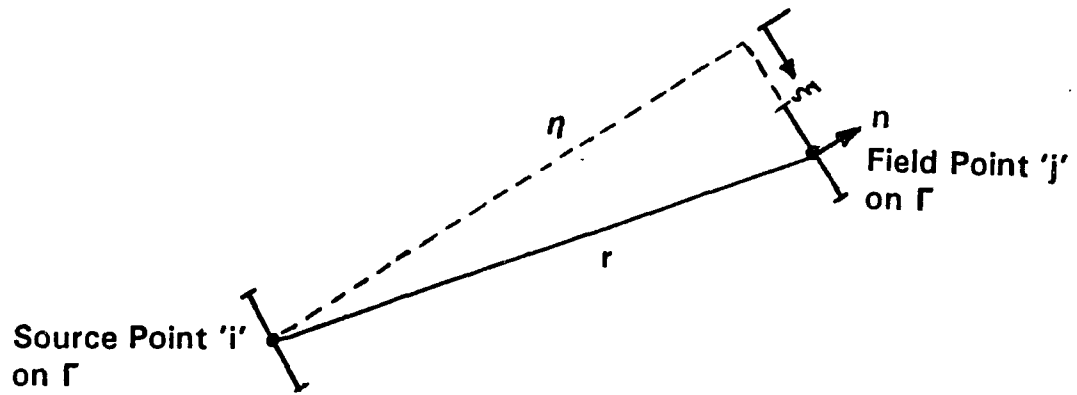


Figure 4-1. Geometrical Definition of a Local Coordinate System.

$$\int_{t_0}^{t_1} \frac{\partial G}{\partial n} = -\frac{1}{2\pi \hat{D} r^2} e^{-a} \quad (4-8)$$

where,

$$a = \frac{r^2}{4\hat{D}\Delta t}$$

$E(a)$ = exponential-integral function.

The exponential-integral function can be expressed in the form of a series as (Brebbia et al., 1984):

$$E(z) = -C - \ln z + \sum_{n=1}^{\infty} (-1)^{n-1} \frac{z^{n-1}}{n \cdot n!} \quad (4-9)$$

where,

$$C = 0.57721566$$

For convenience in numerical computation, the exponential-integral function is evaluated by polynomial and rational expansions as (Abramowitz and Stegan, 1970):

$$E(z) = -\ln z + a_0 + a_1 z + a_2 z^2 + a_3 z^3 + a_4 z^4 + a_5 z^5 + \varepsilon(z) \quad (4-10)$$

$$0 \leq 1; |\varepsilon(z)| < 2 \cdot 10^{-7}$$

After integration in time, the boundary integral equation (4-4) becomes:

$$E(z) = \frac{1}{ze^z} \left[\frac{z^4 + b_1 z^3 + b_2 z^2 + b_3 z + b_4}{z^4 + c_1 z^3 + c_2 z^2 + c_3 z + c_4} \right] \quad (4-11)$$

$$1 \leq z \leq \infty; |e(z)| \leq 2 * 10^{-8}$$

$$\begin{aligned} \lambda \theta = & \int_{\Gamma} \frac{\partial \theta}{\partial n} E(a) \frac{1}{4\pi} d\Gamma - \int_{\Gamma} \theta \left\{ -\frac{\eta e^{-a}}{2\pi r^2} \right\} d\Gamma + \int_{\Omega} \theta_o G d\Omega \\ & - \int_{\Omega} \frac{\partial K(\theta^k)}{\partial z} \frac{E(a)}{4\pi \hat{D}} d\Omega \end{aligned} \quad (4-12)$$

Equation (4-12) now contains only the boundary and domain integrals which are evaluated to solve for the unknowns θ or $\partial\theta/\partial n$ at the boundary. The values of the boundary and area integrals in equation (4-12) can be calculated using numerical quadrature rules. The singularities associated with the boundary and area integrals are evaluated using special quadrature formulae. After evaluating the integrals, the coefficients of the Dirichlet and Neumann boundary conditions in equation (4-12) are obtained. The whole set of equations for N number of nodes at the boundary can be written in matrix form as:

$$[\tilde{A}] \{\theta_i\} = [\tilde{B}] \left\{ \left(\frac{\partial \theta}{\partial n} \right)_i \right\} + [\tilde{H}] \quad (4-13)$$

$$i = 1, 2, 3, \dots, N$$

where,

\tilde{A} and \tilde{B} are the coefficient matrices, and

\tilde{H} is a known column vector that corresponds to the sum of the values of the domain integrals in equation (4-12).

The system of equations (4-13) reveals that N_1 values of θ and N_2 values of $\partial\theta/\partial n$ are known and therefore, there are a set of N ($N = N_1 + N_2$) unknowns in the expression (4-13). Rearranging the equation (4-13) in such a way that all the unknowns are on the right hand side, one can write the expression in matrix form as:

$$[\tilde{C}] [\tilde{X}] = [\tilde{F}] \quad (4-14)$$

Here, on the left hand side \tilde{C} is the known coefficient matrix, \tilde{X} is the unknown (for θ or $\partial\theta/\partial n$) column vector and on the right hand side \tilde{F} is the known column vector correspond to the sum of the values of the domain integrals and the value of the integrals with known boundary conditions. The system of equations (4-14) can be solved for the unknown boundary conditions using Gauss-Jordan elimination method.

4.3 BOUNDARY INTEGRATION

The boundary integrations associated with equation (4-12) are carried out by considering constant elements with mid-node as shown in Figure 4-2. In a typical element, the values of θ and $\partial\theta/\partial n$ are assumed constant and equal to the value at mid-node of the element. The boundary is divided into N elements of which N_1 belongs to Γ_1 (Dirichlet boundary) and N_2 belongs to Γ_2 (Neumann boundary). For a node 'i' at the boundary, equation (4-12) can be written as:

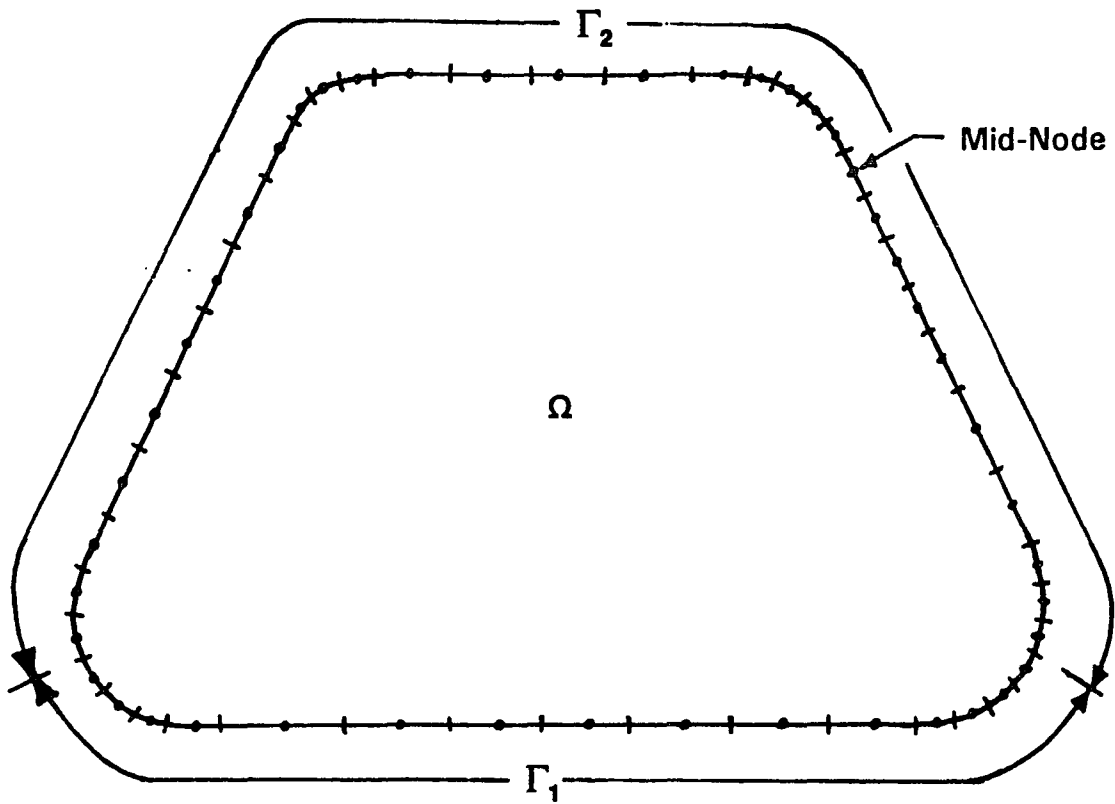


Figure 4-2. Constant Boundary Elements.

$$\lambda \theta_i = \sum_{j=1}^N \left(\frac{\partial \theta}{\partial n} \right)_j A_{ij} - \sum_{j=1}^N \theta_j B_{ij} + C_i - D_i \quad (4-15)$$

where,

$$A_{ij} = \int_{\Gamma} \frac{E(a)}{4\pi} d\Gamma \quad (4-16)$$

$$B_{ij} = \int_{\Gamma} \left\{ -\frac{\eta}{2\pi r^2} \right\} e^{-a} d\Gamma \quad (4-17)$$

$$C_i = \int_{\Omega} \theta_o G d\Omega \quad (4-18)$$

$$D_i = \int_{\Omega} \frac{\partial K(\theta^k)}{\partial z} \frac{E(a)}{4\pi \hat{D}} d\Omega \quad (4-19)$$

The values of the integrals A_{ij} and B_{ij} are calculated using a four-point Gauss quadrature rule. In Gauss quadrature rule, the integration is performed as follows (Carnahan et al., 1972):

$$\int_{-1}^1 f(\xi) d\xi = \sum_{m=1}^n \omega_m f(\xi_m) \quad (4-20)$$

where ξ_m are the base-point values and ω_m are the weight factors corresponding to the base-points. The values of the base-points and the weight factors are obtained from Brebbia (1978). The integrals A_{ij} and B_{ij} are of the type as follows:

$$A_{ij} = \sum_{m=1}^4 \frac{E(a_{im})}{4\pi} \omega_m \frac{\Delta s}{2} \quad (4-21)$$

$$B_{ij} = \sum_{m=1}^4 -\frac{\eta}{2\pi r_{im}^2} \exp(-a_{im}) \frac{\Delta s}{2} \quad (4-22)$$

where,

$$a_{im} = \frac{r_{im}^2}{4\hat{D}\Delta t}$$

r_{im} = distance between the source point and the numerical integration points.

Δs = length of field element,

ω_m = weight factor for the m th integration point.

The integral A_{ij} has a logarithmic singularity in the exponential-integral function as $j \rightarrow i$. The integral A_{ij} for singularity is evaluated as (Brebbia et al., 1982):

$$A_{ij} = \frac{\Delta s}{4\pi} \left[-C - \ln \alpha + 2 + \sum_{m=1}^{\infty} (-1)^{m-1} \frac{\alpha^m}{m(2m+1)m!} \right] \quad (4-23)$$

for $i=j$

where,

$$\alpha = \frac{(\Delta s)^2}{16\hat{D}\Delta t}$$

4.4 AREA INTEGRATION

The integrals defined by equations (4-18) and (4-19) are evaluated by dividing the domain into a series of triangular cells as shown in Figure 4-3. On each cell the integration can be performed as (Brebbia, 1978):

$$C_i = \sum_{m'=1}^{n_e} \left\{ \sum_{m=1}^n \omega_m(\theta_o G)_m \right\} A_{m'} \quad (4-24)$$

$$D_i = \sum_{m'=1}^{n_e} \left\{ \sum_{m=1}^4 \omega_m \left[\frac{\partial K(\theta^k)}{\partial z} \frac{E(a)}{4\pi\hat{D}} \right]_m \right\} A_{m'} \quad (4-25)$$

where,

n_e = number of elements,

n = number of integration points on each element,

ω_m = weighting function, and

$A_{m'}$ = area of a cell or an element.

In this case, the integration points are considered as the vertices of triangular elements. The weighting functions for the integration points are used from Tables given by Zienkiewicz and Taylor (1989).

The singularity arises for the domain integral D_i when the source point coincides with one of the vertices of the triangular element. The logarithmic singularity in D_i due to the exponential-integral function is evaluated using the special quadrature

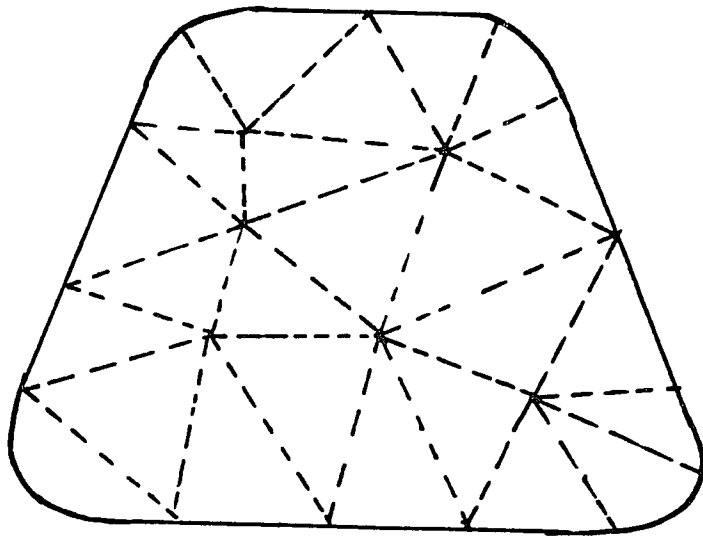


Figure 4-3. Triangular Elements for Numerical Integration in the Domain.

rule suggested by Liggett (1982). A general situation for singularity is shown at the origin (point 1) in Figure 4-4. The special quadrature formula for the integral due to logarithmic singularity (I_s) associated with the exponential-integral function in D_i is expressed as (Liggett, 1982):

$$I_s = -\tilde{D}\eta_s^2\left[(\tilde{u}-\tilde{v})\left(\ln\eta_s - \frac{3}{2}\right) + \frac{\tilde{u}}{2}\ln(1+\tilde{u}^2) - \frac{\tilde{v}}{2}\ln(1+\tilde{v}^2) + \tan^{-1}\tilde{u} - \tan^{-1}\tilde{v}\right] \quad (4-26)$$

where,

$$\tilde{D} = \frac{1}{4\pi\hat{D}} \frac{\partial K(\theta^k)}{\partial z}$$

u = slope of line 1-3 as shown in Figure 4-4,

v = slope of line 1-2 as shown in Figure 4-4, and

η_s = perpendicular distance from point 1 to line 2-3 as shown in Figure 4-4.

In the series expansion (4-9) for exponential-integral function, the log singularity associated with the integral D_i is separated to evaluate using the above expression. The rest of the series in integral D_i is evaluated using the four-point Gauss quadrature rule.

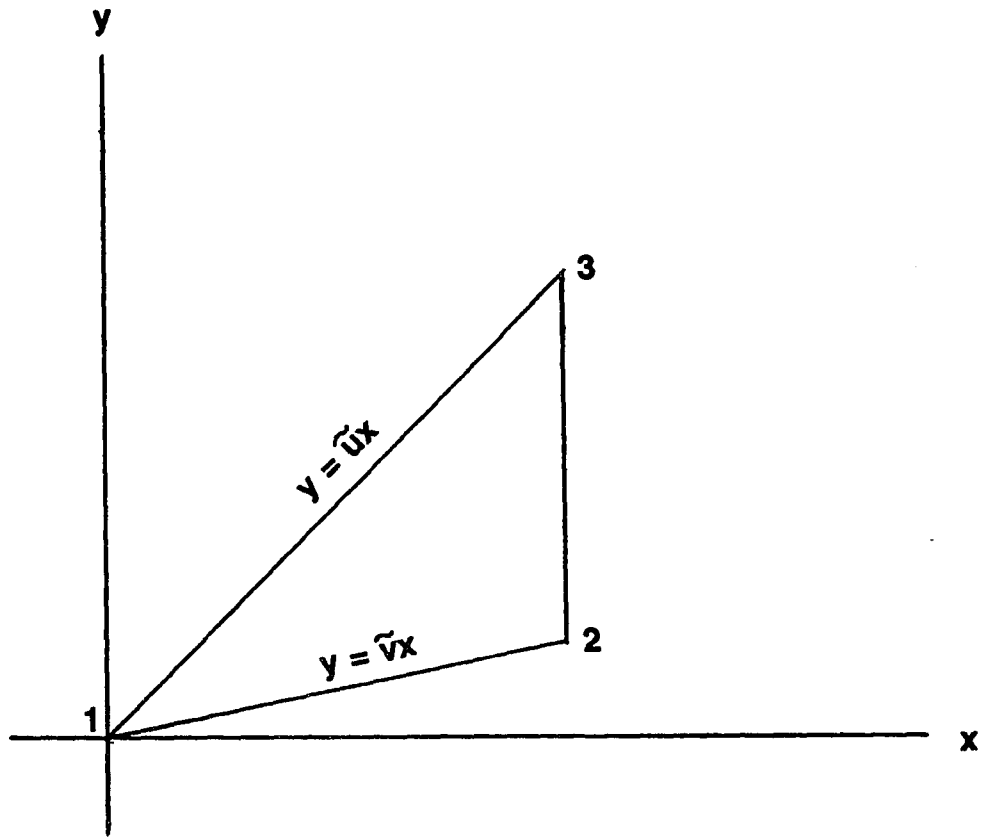


Figure 4-4. Definition Sketch of a Triangle for Singular Integral Computation (Liggett, 1982).

4.5 INTERNAL FLUXES

The velocities of the moisture flow in the unsaturated zone is of particular interest to predict leachate flow in a landfill. The moisture flux is used to compute leachate accretion onto the saturated leachate mound. The recharge to the saturated leachate mound depends on the moisture-flux in the unsaturated zone of a landfill. In the boundary integral formulation, the point value of the moisture-flux is calculated by differentiating the boundary integral formulation of the dependent variable. The moisture flow velocity in the x-direction is obtained by differentiating equation (4-5) as:

$$\begin{aligned} \frac{\partial \theta}{\partial x} = & \hat{D} \int_{\Gamma} \int_{t_0}^{t_1} \frac{\partial G}{\partial x} \frac{\partial \theta}{\partial n} d\tau d\Gamma - \hat{D} \int_{\Gamma} \int_{t_0}^{t_1} \theta \frac{\partial^2 G}{\partial x \partial n} d\tau d\Gamma + \int_{\Omega} \theta_0 \frac{\partial G}{\partial x} d\Omega \\ & - \int_{\Omega} \int_{t_0}^{t_1} \frac{\partial G}{\partial x} \frac{\partial K(\theta^k)}{\partial z} d\tau d\Omega \end{aligned} \quad (4-27)$$

Similarly in the z-direction, the moisture-flux is expressed as:

$$\begin{aligned} \frac{\partial \theta}{\partial z} = & \hat{D} \int_{\Gamma} \int_{t_0}^{t_1} \frac{\partial G}{\partial z} \frac{\partial \theta}{\partial n} d\tau d\Gamma - \hat{D} \int_{\Gamma} \int_{t_0}^{t_1} \theta \frac{\partial^2 G}{\partial z \partial n} d\tau d\Gamma + \int_{\Omega} \theta_0 \frac{\partial G}{\partial z} d\Omega \\ & - \int_{\Omega} \int_{t_0}^{t_1} \frac{\partial G}{\partial z} \frac{\partial K(\theta^k)}{\partial z} d\tau d\Omega \end{aligned} \quad (4-28)$$

After time integration, the expressions (4-27) and (4-28) contains the boundary and area integrals and are expressed as:

$$\begin{aligned} \frac{\partial \theta}{\partial x} = & \int_{\Gamma} \left\{ -\frac{e^{-a}}{2\pi r} \frac{\partial r}{\partial x} \right\} \frac{\partial \theta}{\partial n} d\Gamma - \int_{\Gamma} \left\{ \frac{\eta(1+a)e^{-a}}{\pi r^3} \frac{\partial r}{\partial x} - \frac{e^{-a}}{2\pi r^2} \frac{\partial \eta}{\partial x} \right\} \theta d\Gamma \\ & + \int_{\Omega} \theta_o \left\{ \frac{-re^{-a}}{8\pi \hat{D}^2 (\Delta t)^2} \frac{\partial r}{\partial x} \right\} d\Omega - \int_{\Omega} \frac{\partial K(\theta^k)}{\partial z} \left\{ \frac{e^{-a}}{2\pi r D} \frac{\partial r}{\partial x} \right\} d\Omega \quad (4-29) \end{aligned}$$

$$\begin{aligned} \frac{\partial \theta}{\partial z} = & \int_{\Gamma} \left\{ -\frac{e^{-a}}{2\pi r} \frac{\partial r}{\partial z} \right\} \frac{\partial \theta}{\partial n} d\Gamma - \int_{\Gamma} \left\{ \frac{\eta(1+a)e^{-a}}{\pi r^3} \frac{\partial r}{\partial z} - \frac{e^{-a}}{2\pi r^2} \frac{\partial \eta}{\partial z} \right\} \theta d\Gamma \\ & + \int_{\Omega} \theta_o \left\{ \frac{-re^{-a}}{8\pi \hat{D}^2 (\Delta t)^2} \frac{\partial r}{\partial z} \right\} d\Omega - \int_{\Omega} \frac{\partial K(\theta^k)}{\partial z} \left\{ \frac{e^{-a}}{2\pi r D} \frac{\partial r}{\partial z} \right\} d\Omega \quad (4-30) \end{aligned}$$

The boundary and area integrations are done as explained before. Once the unknown boundary values θ and $\partial\theta/\partial n$ are determined, they are substituted into equations (4-29) and (4-30) to obtain internal fluxes.

4.6 PERTURBATION-GREEN'S FUNCTION (PGF) FORMULATION

A perturbation technique similar to Lafe and Cheng (1987) is developed to solve equation (4-1) by evaluating the gravity term $\partial K(\theta)/\partial z$ at the current time level. It is assumed that the volumetric moisture content θ in the x-z-t space can be expressed as:

$$\theta = \theta_o + \theta_1 + \theta_2 + \theta_3 + \dots \quad (4-31)$$

where,

$$\theta_0 \gg \theta_1 \gg \theta_2 \gg \dots$$

It is required that the leading term in the series satisfies the moisture-flow equation with no source term. A sequence of moisture-flow equations with known right-hand sides are written. Therefore, the following set of equations can be written as:

$$\frac{\partial \theta_0}{\partial t} - \hat{D} \nabla^2 \theta_0 = 0 \quad (4-32)$$

$$\frac{\partial \theta_i}{\partial t} - \hat{D} \nabla^2 \theta_i = -\frac{\partial K(\theta_{i-1})}{\partial z} \quad (4-33)$$

$$i = 1, 2, 3, 4, 5, \dots$$

Once equation (4-32) is solved for θ_0 , then the right-hand side of equation (4-33) indicating the gravity term $\partial K(\theta_0)/\partial z$ is known. Finally, equation (4-33) is solved for $\theta_1, \theta_2, \theta_3, \dots$, using the value of θ from the preceding level of solution.

In the perturbation quantities, it is required that the leading term absorb the entire boundary conditions:

$$\theta_o = \bar{\theta} \quad \text{on } \Gamma_1$$

$$\frac{\partial \theta_o}{\partial n} = i_n \quad \text{on } \Gamma_2$$

$$\theta_i = 0 \quad \text{on } \Gamma_1$$

$$\frac{\partial \theta_i}{\partial n} = 0 \quad \text{on } \Gamma_2$$

$$i = 1, 2, 3, 4, \dots$$

Therefore, a new system of equations is obtained from the original boundary value problem. The new system consists of solving the homogeneous equation with the original boundary and initial conditions. A sequence of nonhomogeneous equations is solved with homogeneous boundary conditions and known right-hand side obtained from the preceding level of solution. Successive solutions are required to find the final result.

The boundary integral formulation is the same as for the DGF formulation explained before. In the boundary integral formulation for the homogeneous equation (4-32), the domain integral D_i defined by equation (4-19) does not exist. The boundary and domain integrations for the nonhomogeneous equations (4-33) are the same as described in direct Green's function formulation. In the domain integral D_i for the nonhomogeneous equation (4-33), the gravity term is evaluated as $\partial K(\theta_{i-1})/\partial z$, where θ_{i-1} is the value of θ at the preceding level of solution.

CHAPTER 5

COMPARISON OF NUMERICAL MODELS WITH ANALYTICAL SOLUTION

This chapter evaluates the performance of the finite-difference and boundary integral solution techniques to solve the unsaturated leachate flow equation. The finite-difference numerical solution technique is compared with the analytical solution for the saturated leachate flow equation. Comparisons of the numerical solution techniques are made with the analytical solution of the governing partial differential equations in simplified forms for idealized flow conditions with regular boundaries of the domain. Although the comparisons of the numerical models are made for idealized flow conditions, confidence is gained for the applicability of the models for complex real field conditions if the models reproduce agreeable true solution.

Analytical solution is derived for the unsaturated leachate flow equation for a rectangular domain defined by both Dirichlet and Neumann boundary conditions. The finite-difference solution of the moisture-flow equation using Gauss-Seidel iteration and Gauss-Jordan elimination methods are compared with the exact solution. The direct Green's function and perturbation Green's function boundary integral solutions are also compared with the finite-difference and exact solutions.

The analytical solution for the saturated leachate flow equation is obtained for a drainage layer with known leachate accretion onto it. The two numerical solution techniques, Gauss-Seidel iteration and Gauss-Jordan elimination are used to solve the simplified form of the saturated leachate flow equation. The numerical solution results are compared with the exact solutions.

5.1 ANALYTICAL SOLUTION

Analytical solutions are obtained for both the governing partial differential equations in the unsaturated and saturated zones of a landfill. The analytical solution for the unsaturated flow is done on a two-dimensional rectangular domain considering unsteady Dirichlet boundary condition. The steady Neumann boundary condition is assumed for both the unsaturated and saturated leachate flow equation.

5.1.1 Unsaturated Flow Equation

Analytical solution is difficult to obtain for the complex moisture-flow equation including all the flow components active in a landfill. However, analytical solutions can be obtained for the unsteady moisture-flow equation with the help of some simplifying assumptions. The assumptions made in the derivation of the two-dimensional moisture flow equation are :

- A well-defined rectangular domain is considered with Dirichlet and Neumann boundary conditions;
- The diffusivity is a function of moisture content. It is estimated using the moisture content from the previous time step; and
- The flow due to gravity is obtained from the moisture content at the previous time-step. This condition is obtained by defining the gravity term as $\partial K(\theta^k)/\partial z$, where θ^k is the moisture content at the previous time step $t=k$.

With the above simplifying assumptions, the unsaturated moisture-flow equation (2-11) becomes:

$$\frac{\partial \theta}{\partial t} = \hat{D} \left(\frac{\partial^2 \theta}{\partial x^2} + \frac{\partial^2 \theta}{\partial z^2} \right) - p \quad (5-1)$$

where,

$$p = \frac{\partial K(\theta^k)}{\partial z}$$

Analytical solution for the above equation is obtained using the method of separation of variables with respect to Figure 5-1a for the following boundary conditions:

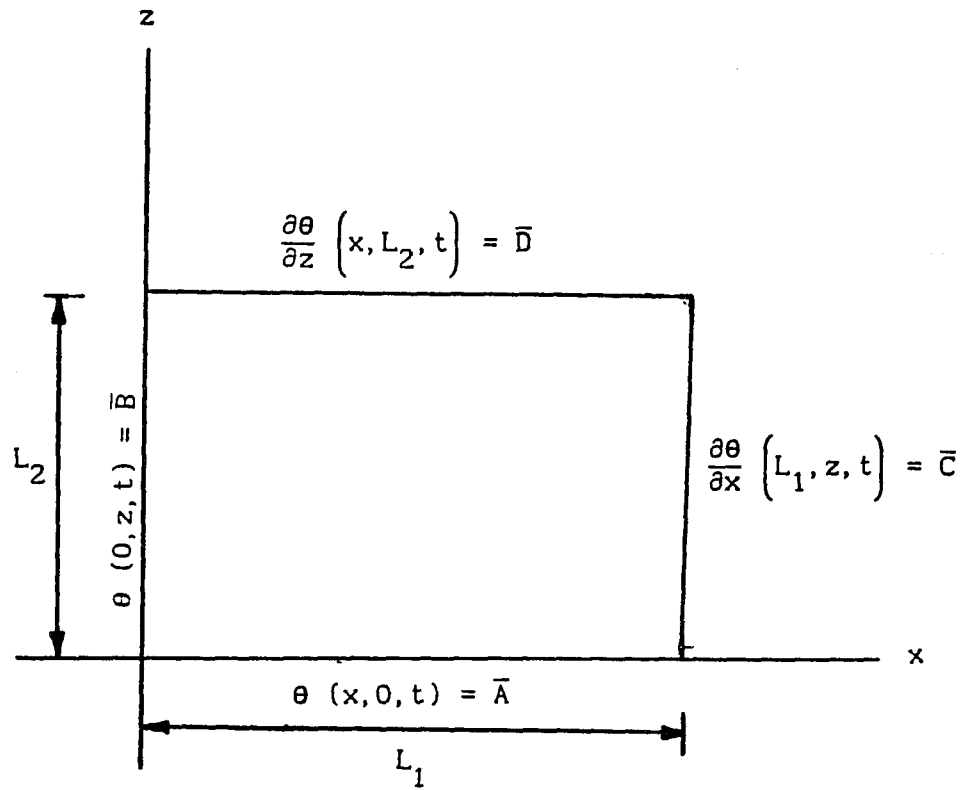


Figure 5-1a. Computational Domain for Analytical Solution.

$$\theta(x, 0, t) = \bar{A}$$

$$\theta(0, z, t) = \bar{B}$$

$$\frac{\partial \theta}{\partial x}(L_1, z, t) = \bar{C}$$

$$\frac{\partial \theta}{\partial z}(x, L_2, t) = \bar{D}$$

The initial condition is as follows:

$$\theta(x, z, 0) = \theta_0$$

The solution for the moisture content θ , obtained for the above boundary and initial conditions is expressed as:

$$\theta = \sum_{n=0}^{\infty} \left[\sum_{m=0}^{\infty} C_{mn}(t) \sin \frac{(2m+1)\pi x}{2L_1} + \beta_n + \frac{x^2}{2L_1} \hat{C} \right] \sin \frac{(2n+1)\pi z}{2L_2} + \bar{A} + \frac{z^2}{2L_2} \bar{D} \quad (5-2)$$

where,

$$C_{mn}(t) = \frac{t_{mn}}{\alpha} + \gamma e^{-\alpha t}$$

$$t_{mn} = \frac{4}{(2m+1)\pi} [t_{mn1} - t_{mn2}]$$

$$t_{mn1} = \bar{t}_n + \alpha_c - \alpha_\beta$$

$$t_{mn2} = \frac{\hat{u}}{L_1} \left\{ \frac{2L_1^2}{(2m+1)\pi} \sin \frac{(2m+1)\pi}{2} - \left[\frac{2L_1}{(2m+1)\pi} \right]^2 \right\}$$

$$\bar{t}_n = \frac{4}{(2n+1)\pi} [\alpha_D - p]$$

$$\alpha_D = \frac{\hat{D}\bar{D}}{L_2}$$

$$\alpha_c = \frac{\hat{D}\hat{C}}{L_1}$$

$$\hat{C} = \frac{4\bar{C}}{(2n+1)\pi}$$

$$\alpha_\beta = \hat{D} \left\{ \frac{(2n+1)\pi}{2L_2} \right\}^2 \beta_n$$

$$\beta_n = \frac{4}{(2n+1)\pi} (\beta_{n1} - \beta_{n2})$$

$$\beta_{n1} = \bar{B} - \bar{A}$$

$$\beta_{n2} = \frac{\bar{D}}{L_2} \left\{ \frac{2L_2^2}{(2n+1)\pi} \sin \frac{(2n+1)\pi}{2} - \left[\frac{2L_2}{(2n+1)\pi} \right]^2 \right\}$$

$$\hat{u} = \hat{D}\hat{C} \left\{ \frac{(2n+1)\pi}{2L_2} \right\}^2$$

The detail derivation of the above analytical solution is given in Appendix A.

The analytical solution for the unsaturated moisture-flow equation (5-1) can be obtained in one-dimensional form by disregarding the gravity term $\partial K(\theta^k)/\partial z$. The

initial and boundary conditions are shown in Figure 5-1b and can be expressed as follows:

$$\theta = \theta_{in} \quad z \leq 0, t = 0$$

$$\theta = \theta_{oo} \quad z = 0, t > 0$$

$$\theta \text{ finite for } z \rightarrow -\infty, t > 0$$

The solution for the one-dimensional moisture-flow problem with the given initial and boundary conditions is expressed as (Eagleson, 1970):

$$\theta(z, t) = \theta_{oo} + (\theta_{in} - \theta_{oo}) \operatorname{erf} \left[\frac{|z|}{2(\hat{D}t)^{1/2}} \right] \quad (5-3)$$

The unsaturated moisture content can be obtained using the above one-dimensional analytical solution. The boundary integral solution for moisture content in a soil column can be compared with the exact solution expressed in equation (5-3).

5.1.2 Saturated Flow Equation

Analytical solution for the saturated leachate flow equation is derived by considering the following assumptions:

- The top of the underlying bottom clay (Figure 2-4) is considered horizontal;

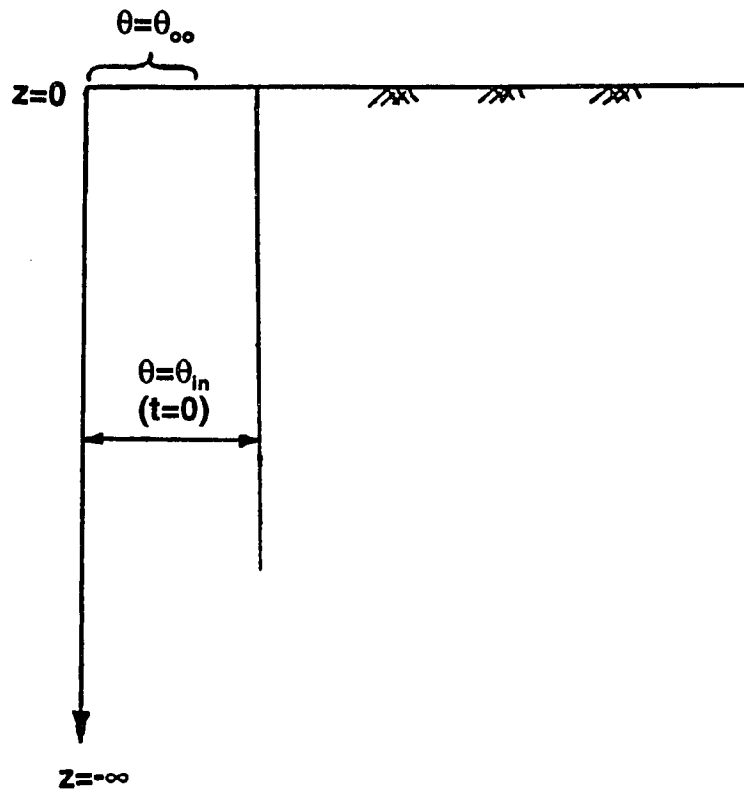


Figure 5-1b. Boundary and Initial Conditions for a One-Dimensional Moisture-Flow Problem (Willis and Yeh, 1987).

- The transmissivity is computed at a certain time level as $T=K_s h^k$, where h^k is the hydraulic head computed at previous time step $t=k$. For this consideration, transmissivity becomes independent of the dependent variable at any time level of computation; and
- The steady state Dirichlet and Neumann boundary conditions are considered.

Saturated leachate flow equation (2-44) is written in a simplified form as:

$$\frac{\partial h}{\partial t} = \alpha \frac{\partial^2 h}{\partial x^2} + \beta h - \bar{p} \quad (5-4)$$

where,

$$\alpha = \frac{T}{n_e}$$

$$\beta = - \frac{K_c}{n_e d} \text{ and}$$

$$\bar{p} = \frac{-N_l + K_c}{n_e}$$

Equation (5-3) is solved for the following initial and boundary conditions:

where,

L = length of the saturated leachate mound from the origin at $x=0$.

$$h(x, t=0) = h_i$$

$$\frac{\partial h}{\partial x}(x=0, t) = i_o$$

$$h(x=L, t) = h_L$$

The method of separation of variables is used to obtain the solution for equation (5-3) for the given boundary and initial conditions. The hydraulic head $h(x,t)$ at any distance x and at any time t is expressed as:

$$h(x, t) = \sum_{n=0}^{\infty} C_n(t) \cos\left[\frac{(2n+1)\pi x}{2L}\right] + i_o x \left(1 - \frac{x}{L}\right) + \frac{x^2 h_L}{L^2} \quad (5-5)$$

where,

$$C_n(t) = \frac{\Phi_n}{\hat{v}} + g_s e^{-\nu t}$$

$$\Phi_n = \frac{4\alpha u_o}{L^2} \left[-i_o + \frac{h_L}{L}\right] + \frac{2}{L} \left[-\bar{p}u_o + \beta i_o u_1 - \frac{\beta i_o}{L} u_2 + \frac{\beta h_L}{L^2} u_2\right]$$

$$u_o = \frac{2L}{(2n+1)\pi} (-1)^n$$

$$u_1 = \frac{2L^2}{(2n+1)\pi} \left[(-1)^n - \frac{2}{(2n+1)\pi}\right]$$

$$u_2 = \frac{2L}{(2n+1)\pi} \left[(-1)^n L^2 - \frac{8(-1)^n}{\pi^2} \left(\frac{L}{2n+1}\right)^2\right]$$

$$g_s = \frac{2}{L} \left[h_1 u_o - i_o u_1 + \frac{i_o}{L} u_2 - \frac{h_L}{L^2} u_2 \right]$$

$$\varphi = \alpha \left[\frac{(2n+1)\pi}{2L} \right]^2 - \beta$$

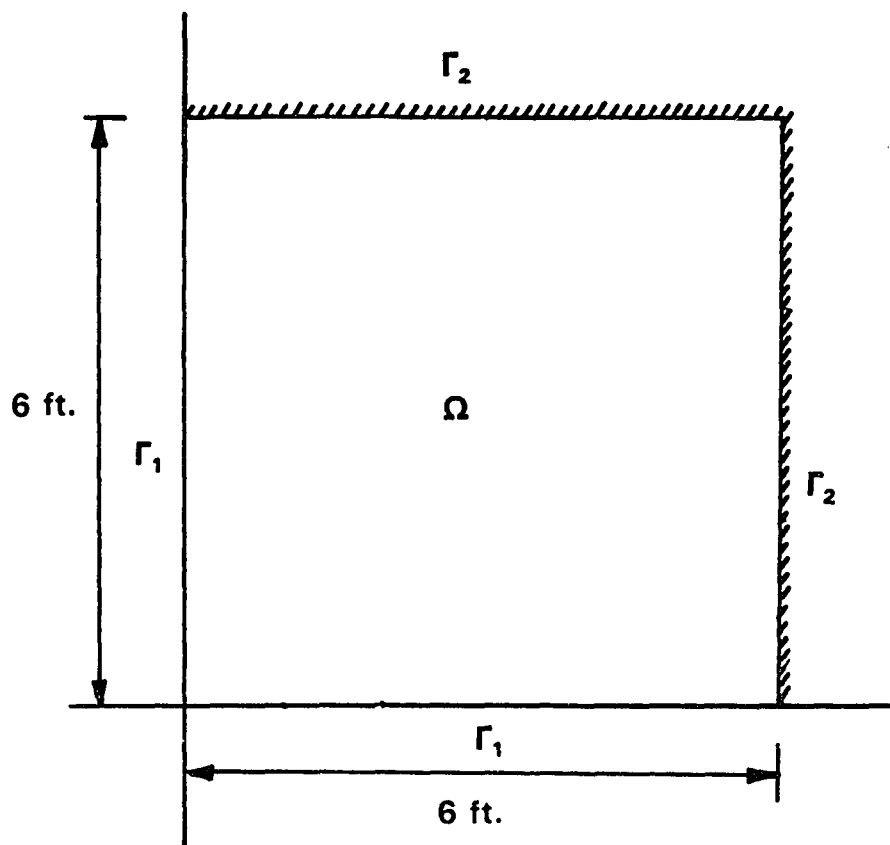
The detail derivation of the above analytical solution is given in Appendix B.

5.2 COMPARISON OF ANALYTICAL AND NUMERICAL SOLUTION

The numerical calculations for the unsaturated leachate flow problem are done on a two-dimensional square domain of size 6 ft. as shown in Figure 5-2. The analytical solution is formulated using simplifying assumptions for the problem defined by Neumann and Dirichlet boundary conditions. The numerical solutions are also obtained for the saturated flow equation. The results for both the unsaturated and saturated flow problems are compared with the exact solutions.

5.2.1 Unsaturated Flow Problem

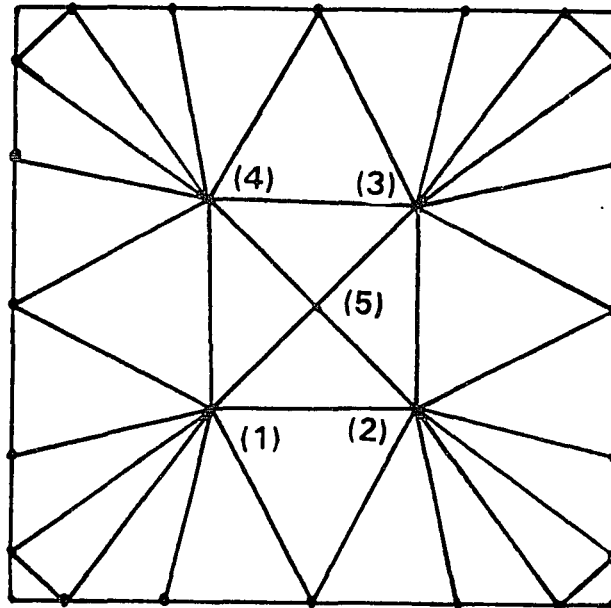
The square region is divided into a series of triangular cells for domain integration in the boundary integral formulation as shown in Figure 5-3. The computational region for the boundary integral solution consists of 20 boundary elements, 28 triangular elements, and 5 internal nodes.



Γ_1 : Dirichlet Boundary Condition

Γ_2 : Neumann Boundary Condition

Figure 5-2. Computational Domain for Comparison of Analytical and Numerical Solution.



(1) : Internal Node No. 1.

Figure 5-3. Computational Domain for Boundary Integral Solution.

For the finite-difference solution, the area is discretized into a network of squares as shown in Figure 5-4. The Neumann boundary condition is considered steady over the computational period. The Dirichlet boundary condition is considered unsteady and the variation is shown Figure 5-5. An initial moisture value of 0.373 is assumed. The other necessary information for the model run are given in Table 5-1.

The direct Green's function and perturbation Green's function boundary integral solutions are compared with the exact solution for node numbers 1 and 3 as shown in Figures 5-6a and 5-6b respectively. Figure 5-6a shows the comparison at $x = 2.0$ ft. and $z = 2.0$ ft. and Figure 5-6b shows the comparison at $x = 4.0$ ft. and $z = 4.0$ ft. From these Figures, it is observed that the boundary integral solutions obtained by direct Green's function and perturbation Green's function formulations are in close agreement with the exact solution. The two boundary integral techniques give similar results.

The assumption for computing the gravity term $\partial K(\theta)/\partial z$ using the moisture content from the previous time step is made in the analytical solution. In the direct green's function solution the same assumption is made to compute the gravity term. However, the gravity term in the perturbation boundary integral method is obtained from the previous higher order solution in the perturbation quantities. The assumption of considering constant gravity term at a certain time step is not possible to use in the perturbation technique. But agreeable results with the direct

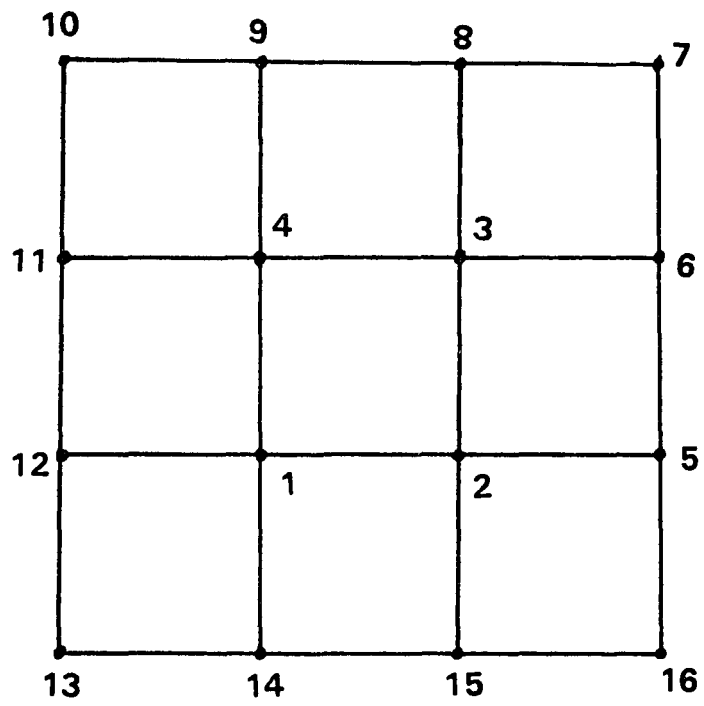


Figure 5-4. Computational Domain for Finite-Difference Solution.

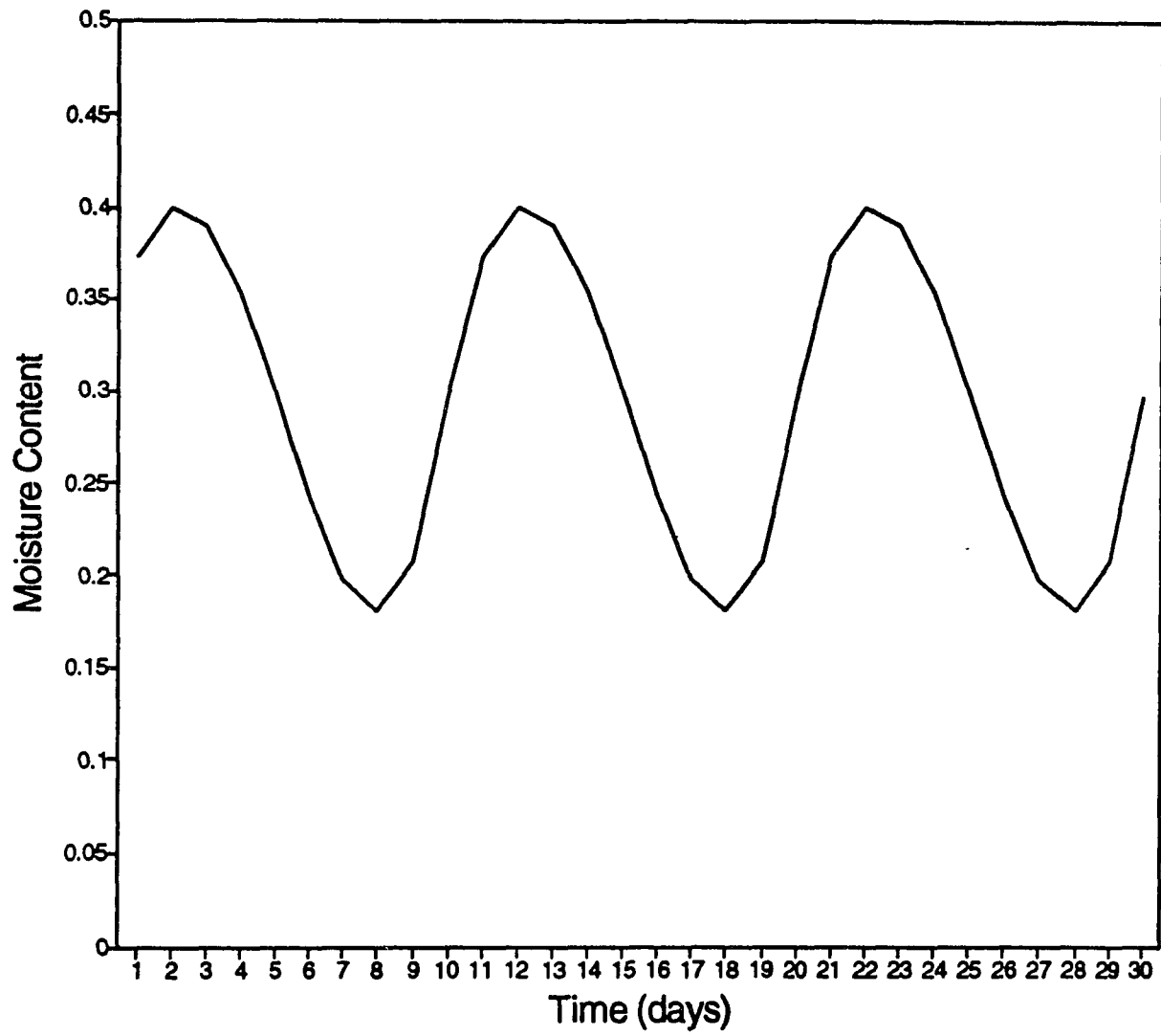


Figure 5-5. Unsteady Dirichlet Boundary Condition.

TABLE 5-1

DATA USED TO SOLVE THE UNSATURATED FLOW EQUATION

Hydraulic conductivity	28 ft./day
Porosity	0.417
Air entry suction head	-2.57 ft.
Campbell's parameter (b)	4.0

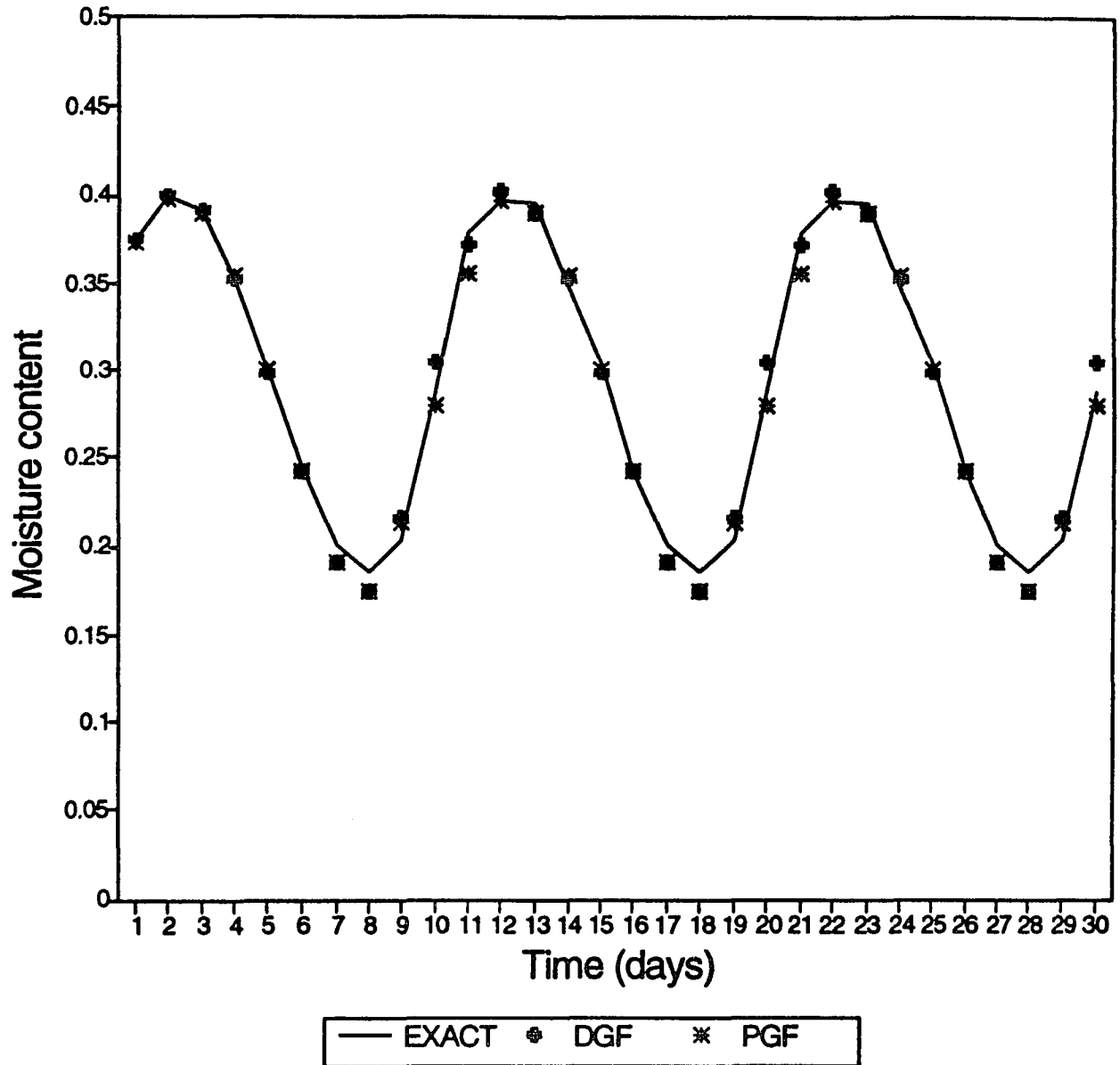


Figure 5-6a. Comparison of Analytical and Boundary Integral Solution at Node No. 1.

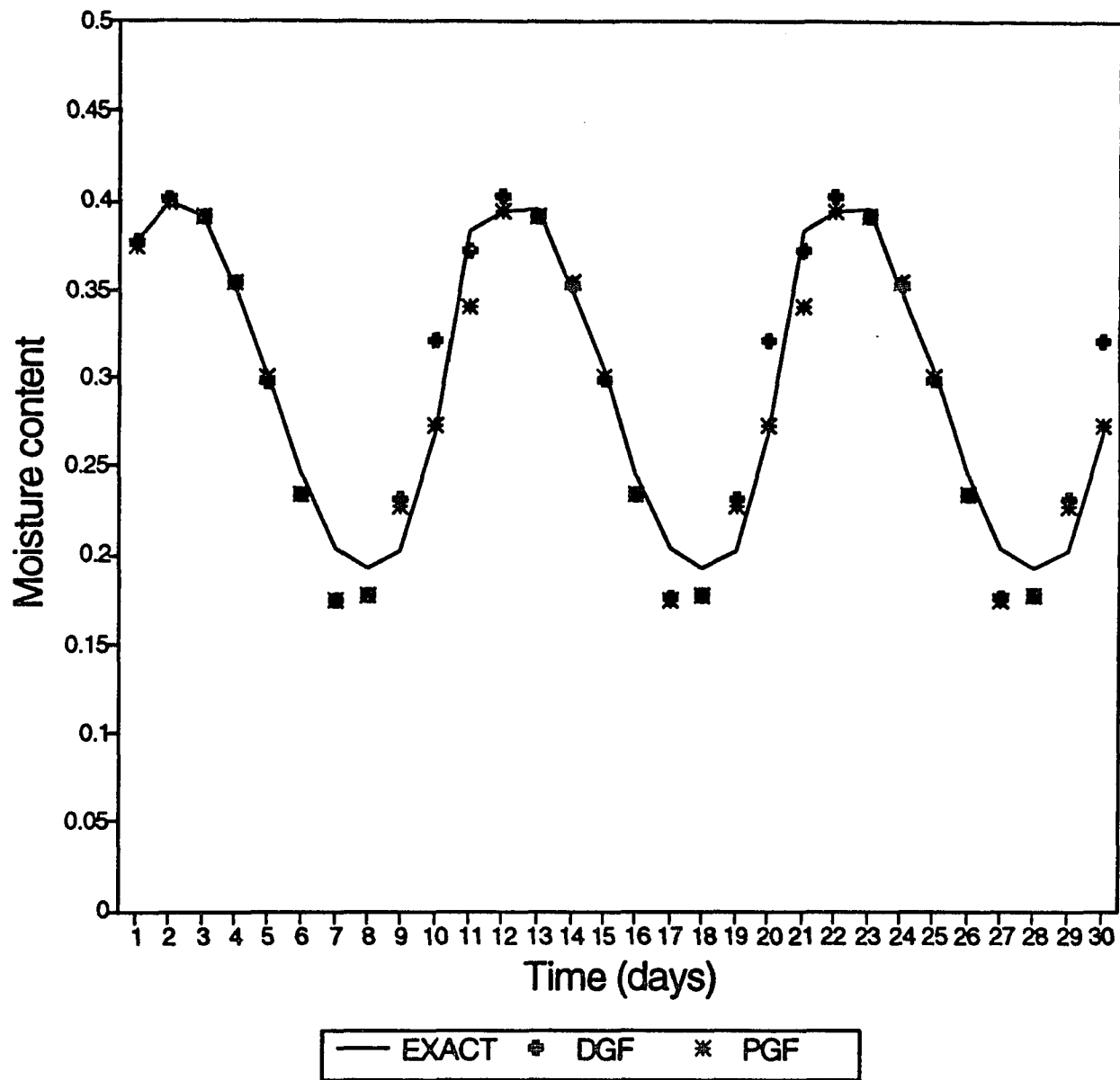


Figure 5-6b. Comparison of Analytical and Boundary Integral solution at Node No. 3.

Green's function and exact solutions are observed.

The finite-difference solution is obtained using the Gauss-Seidel iteration and Gauss-Jordan elimination methods. The direct Green's function boundary integral solution and Gauss-Seidel iteration solution are compared with the analytical solution for internal nodes 1 and 3 as shown in Figures 5-7a and 5-7b respectively. A maximum absolute deviation of 0.04855 from the analytical solution is observed in the case of the Gauss-Seidel iteration method of finite-difference solution at node number 1. However, a maximum absolute deviation of 1.53 percent in moisture content is obtained in the direct Green's function boundary integral solution. The deviations at node number 3 is significant for Gauss-Seidel iteration method of finite-difference solution. A maximum absolute deviation of 0.10779 occurs in the finite-difference solution by Gauss-Seidel iteration method against a maximum absolute deviation of 0.05259 by direct Green's function boundary integral solution at node number 3.

The direct Green's function boundary integral solutions are also compared with the solution obtained by the Gauss-Jordan elimination method as shown in Figures 5-8a and 5-8b. The Gauss-Seidel iteration and Gauss-Jordan elimination methods of finite-difference solution give similar results. The similar results shown by the iteration and elimination methods and the agreeable solution by both the methods with the analytical solution establish the credibility of the finite-difference scheme used for the governing equation. However, the direct Green's function boundary

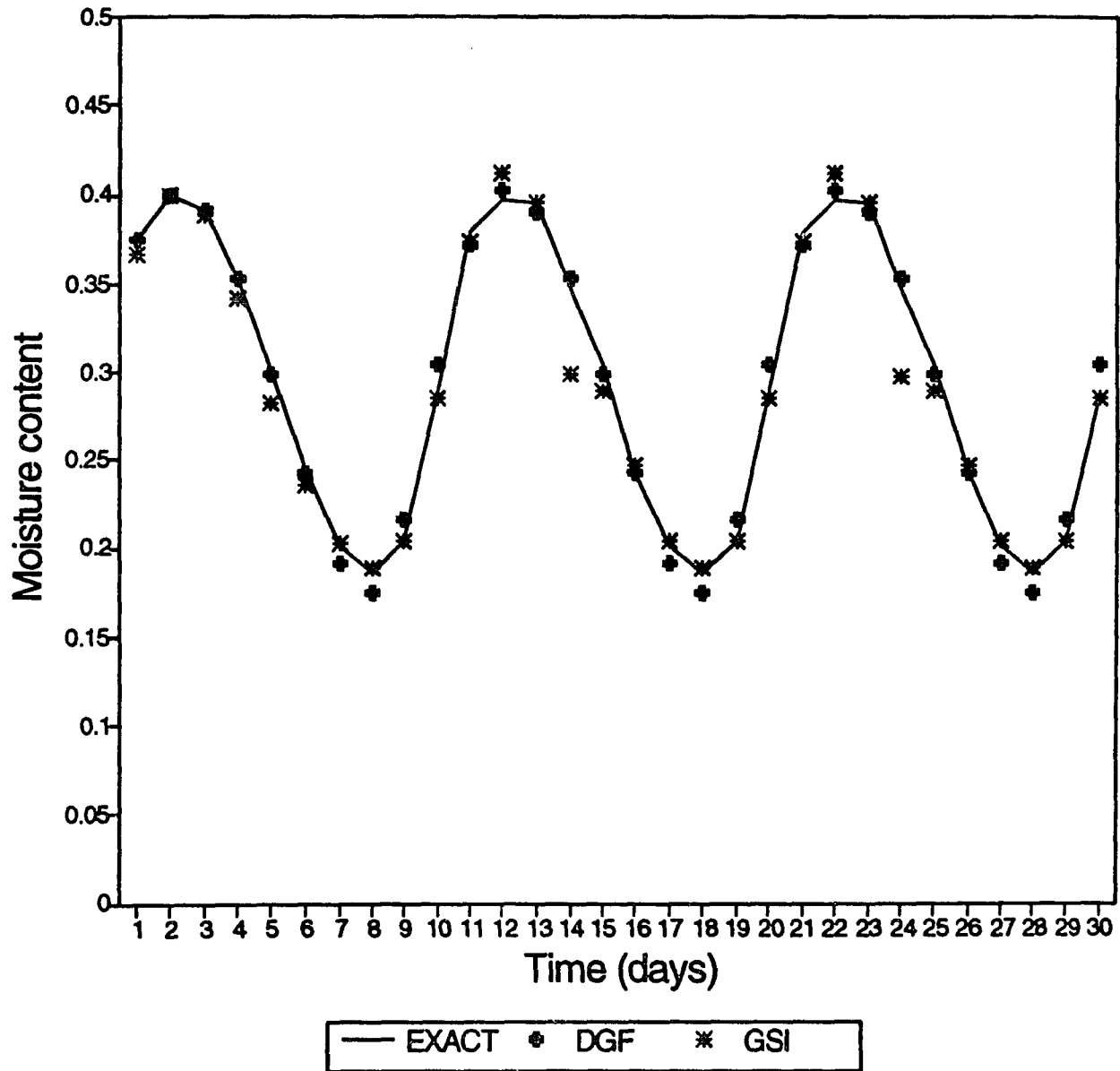


Figure 5-7a. Comparison of Direct Green's Function Boundary Integral Solution and Finite-difference Solution by Gauss-Seidel Iteration Method at Node No. 1.

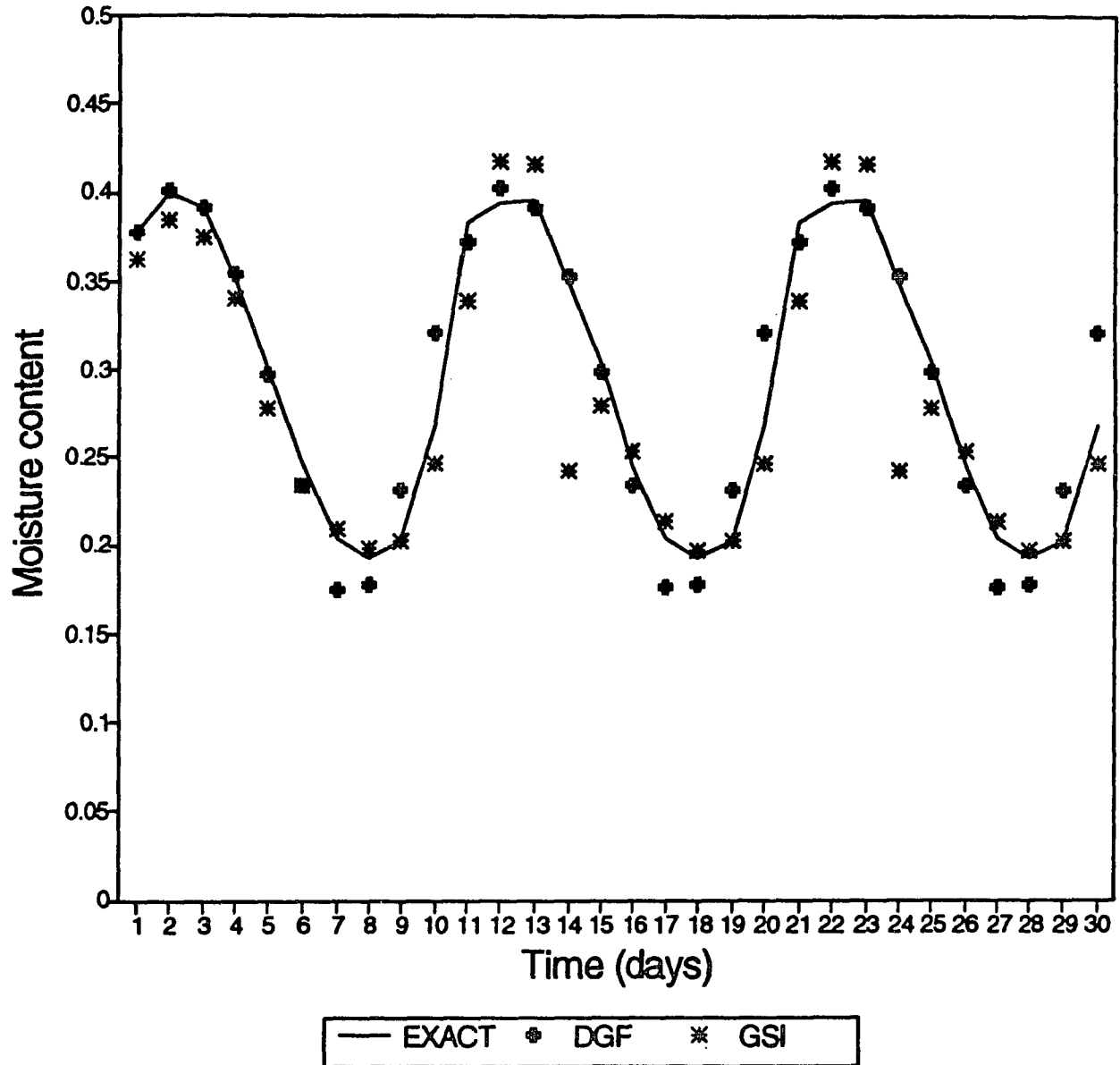


Figure 5-7b. Comparison of Direct Green's Function Boundary Integral Solution and Finite-difference Solution by Gauss-Seidel Iteration Method at Node No. 3.

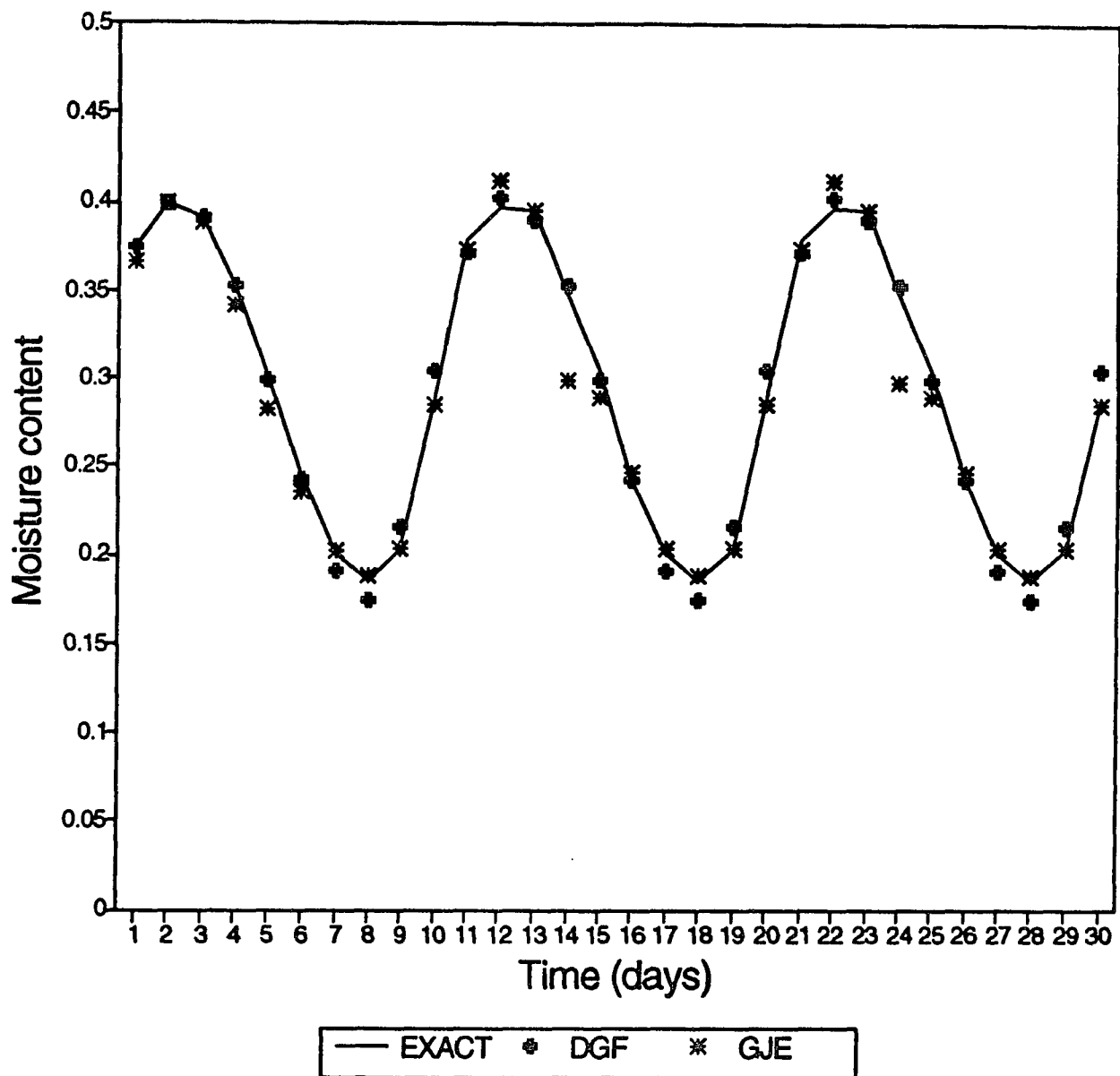


Figure 5-8a. Comparison of Direct Green's Function Boundary Integral Solution and Finite-difference Solution by Gauss-Jordan Elimination method at Node No. 1.

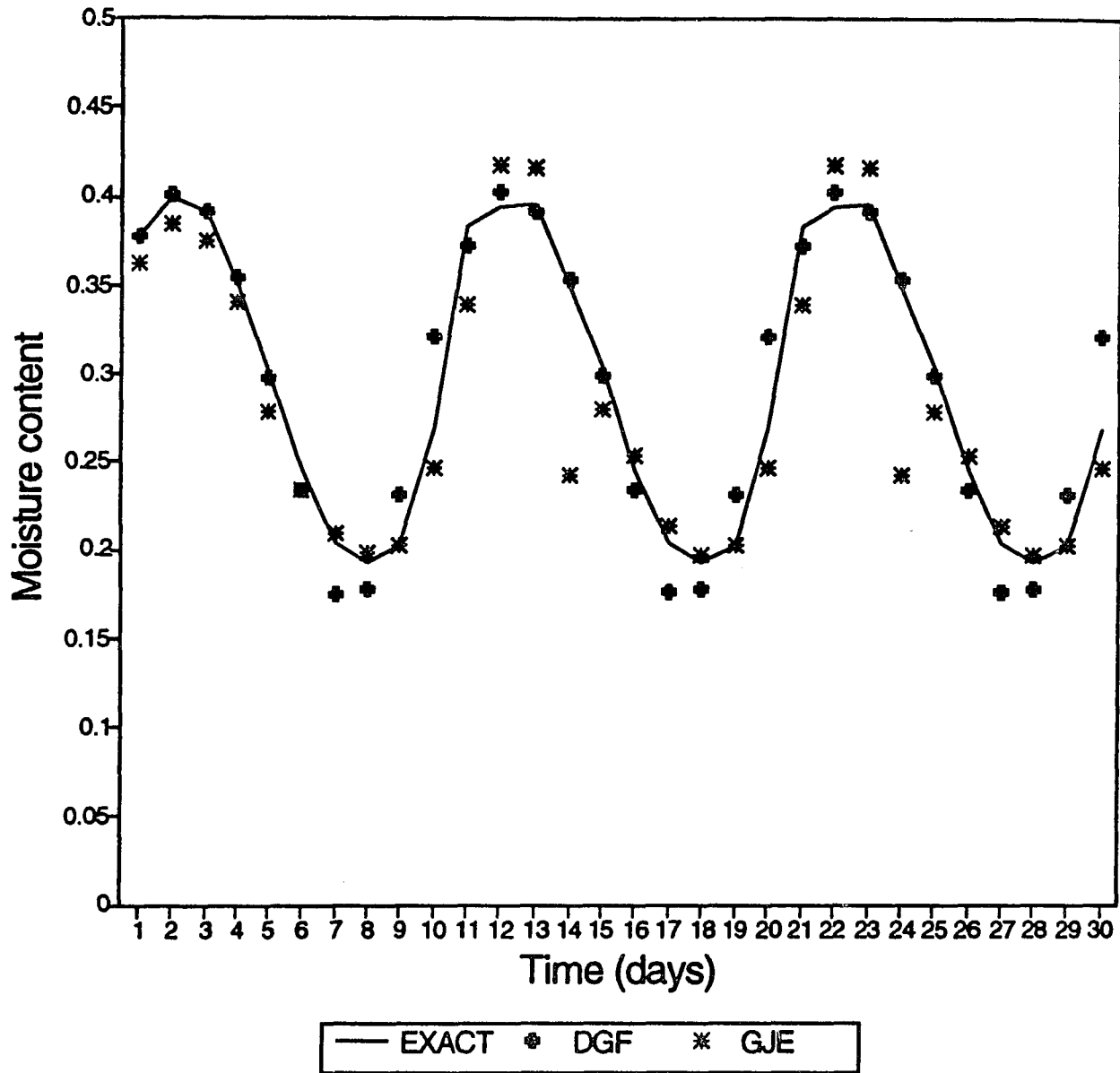


Figure 5-8b. Comparison of Direct Green's Function Boundary Integral Solution and Finite-difference Solution by Gauss-Jordan Elimination method at Node No. 3.

integral method shows more agreeable results than the finite-difference solution obtained by both the Gauss-Seidel iteration and Gauss-Jordan elimination methods as shown in Figures 5-7 and 5-8. This implies that the boundary integral solution gives more accurate results than the finite-difference method for moisture content in the unsaturated zone.

The results of perturbation boundary integral solution and finite-difference solution using Gauss-Seidel iteration method are shown in Figures 5-9a and 5-9b respectively. A maximum absolute deviation of 2.27 percent of moisture content from the analytical solution is found at node number 1 in the case of perturbation Green's function boundary integral solution. At node number 3, the maximum absolute deviation from analytical solution is small in perturbation Green's function solution, compared to the Gauss-Seidel iteration method of finite-difference solution. The maximum absolute deviation of 0.04368 from analytical solution is observed in perturbation Green's function solution against the absolute deviation of 0.10779 in the Gauss-Seidel iteration method of the finite-difference solution.

The perturbation Green's function boundary integral solution is also compared with the finite-difference solution using Gauss-Jordan elimination method. The comparisons with respect to the analytical solution at node numbers 1 and 3 are shown in Figures 5-10a and 5-10b respectively. The deviation in results by Gauss-Jordan elimination method is similar to the deviation in results obtained by Gauss-Seidel iteration method. The computation of gravity term in the perturbation

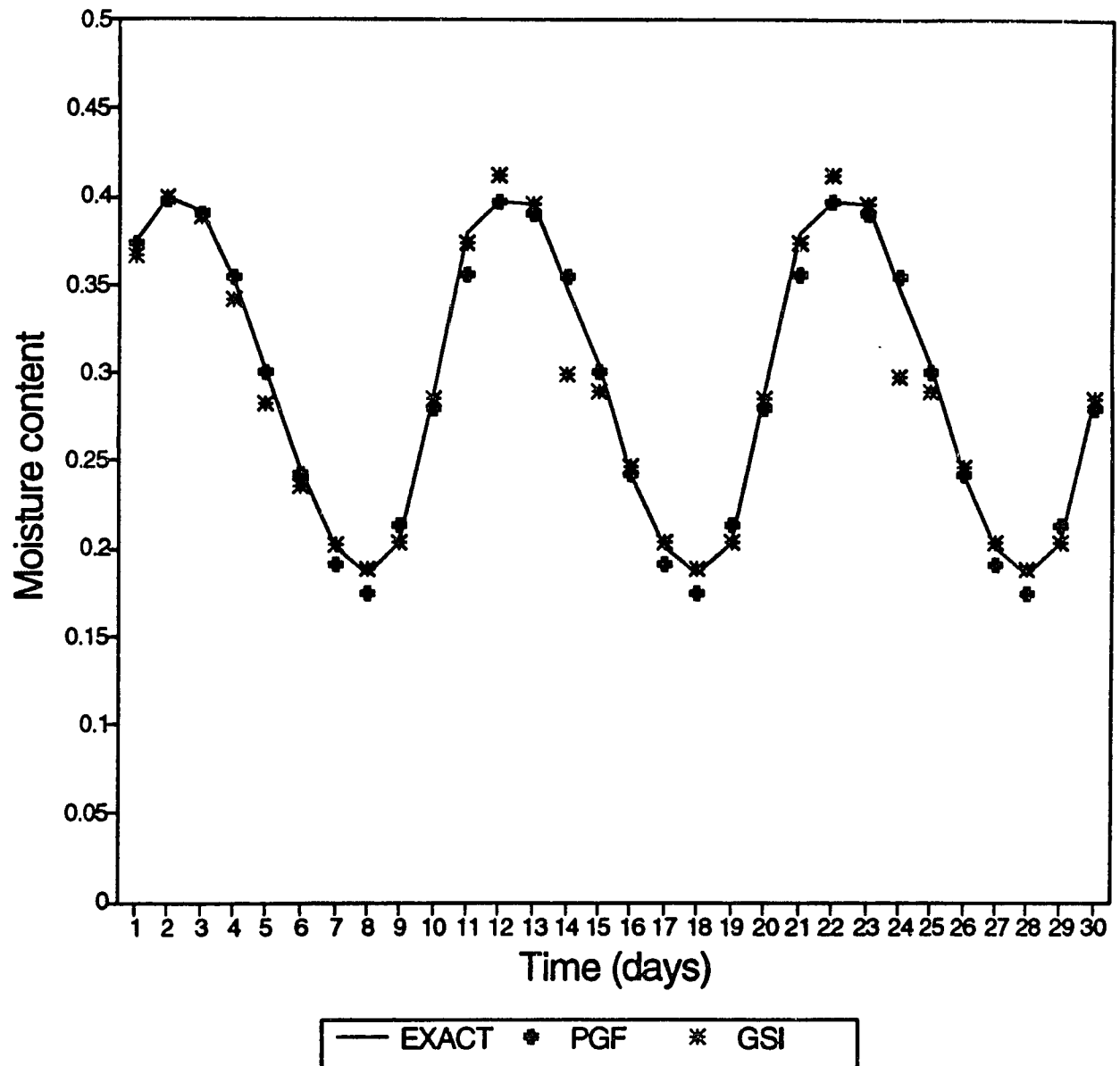


Figure 5-9a. Comparison of Perturbation Green's Function Boundary Integral Solution and Finite-difference Solution by Gauss-Seidel Iteration Method at Node No. 1.

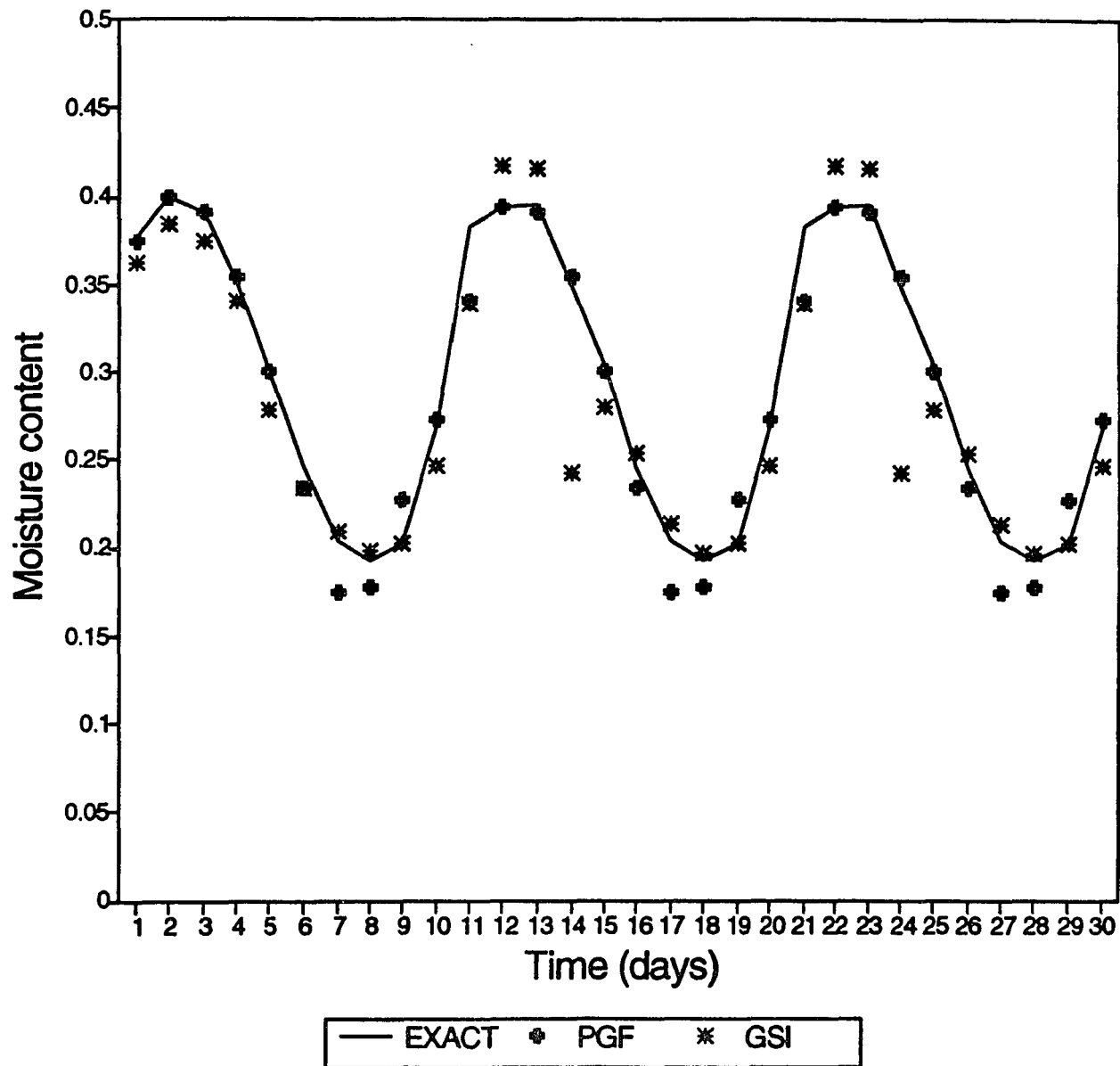


Figure 5-9b. Comparison of Perturbation Green's Function Boundary Integral Solution and Finite-difference Solution by Gauss-Seidel Iteration Method at Node No. 3.

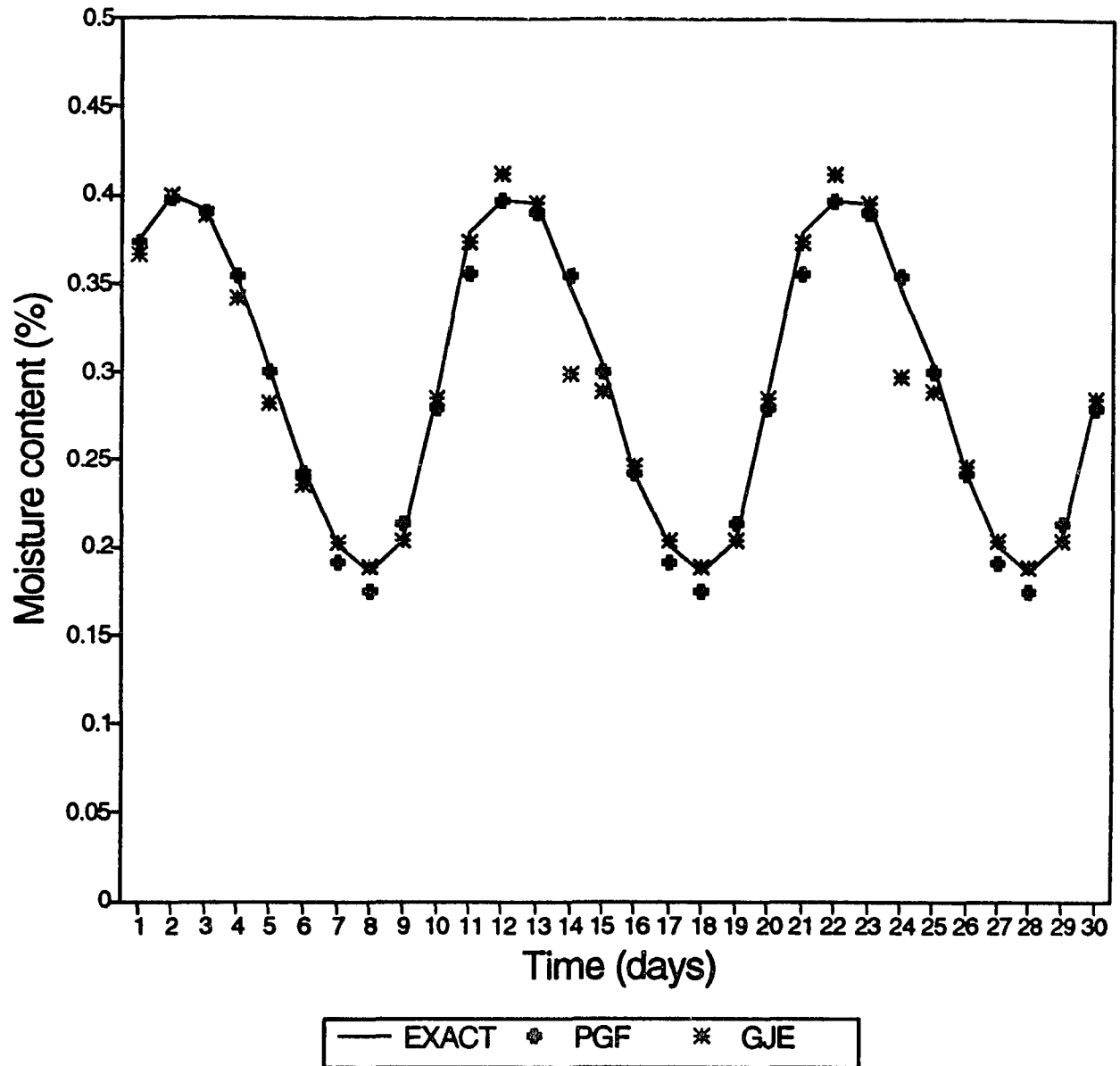


Figure 5-10a. Comparison of Perturbation Green's Function Boundary Integral Solution and Finite-difference Solution by Gauss-Jordan Elimination Method at Node No. 1.

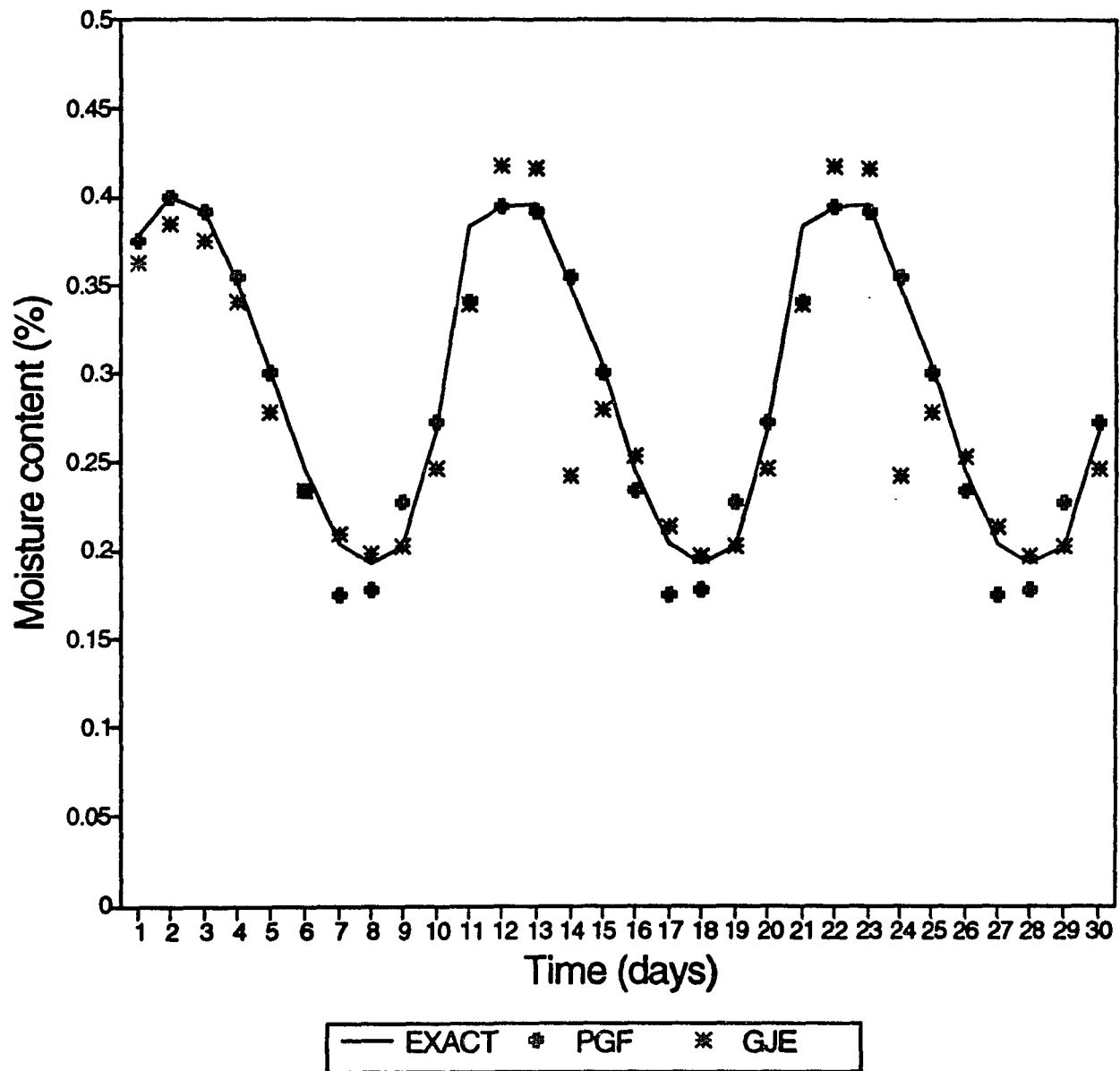
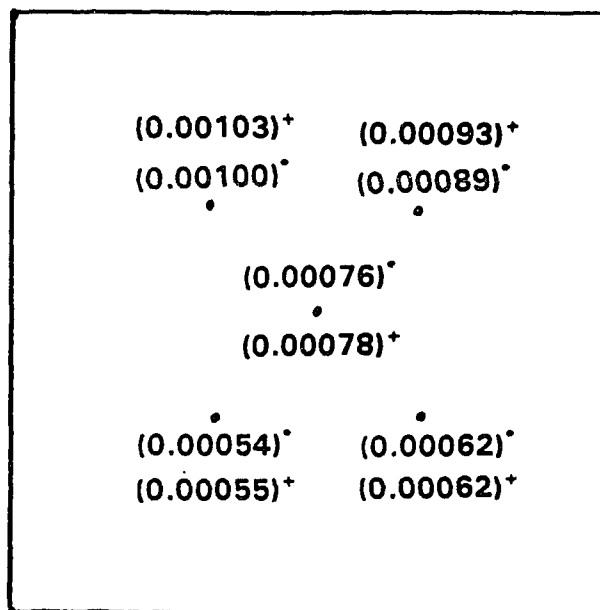


Figure 5-10b. Comparison of Perturbation Green's Function Boundary Integral Solution and Finite-difference Solution by Gauss-Jordan Elimination Method at Node No. 3.

boundary integral solution is done using perturbation series at the current time level. This gives a better representation of the physical system by considering moisture content from the higher order solution of the perturbation series at the current time level to compute the gravity term. It is observed that the perturbation Green's function solution gives more agreeable results compared to the finite-difference solution. Therefore, it can be concluded that the boundary integral method represents the moisture variation in the unsaturated zone of the subsurface more accurately than the finite-difference method.

Unlike the finite-difference method, the velocity at internal points can be computed by boundary integral formulation, once the Dirichlet and Neumann boundary conditions are solved for the boundary nodes. The internal velocities in x- and z-directions are obtained using the expressions (4-29) and (4-30). The internal velocities in x- and z-directions are compared with the analytical solution as shown in Figures 5-11a and 5-11b respectively. Close agreement between the boundary integral and analytical solution shows the ability of the boundary integral method to compute leachate accretion more accurately in the unsaturated zone of a landfill.

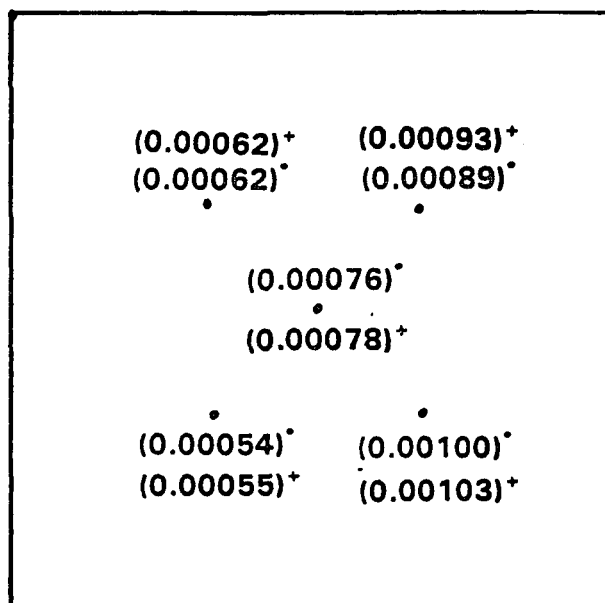
It is valuable to find the differences in the solution obtained by finite-difference and boundary integral techniques from the exact solution. The mean of the absolute values of the deviations between numerical and analytical solutions at each time step is computed for the whole domain. The variation of mean values of the



* Boundary Integral Solution

+ Exact Solution

Figure 5-11a. Internal Flux in x-direction.



* Boundary Integral Solution

+ Exact Solution

Figure 5-11b. Internal Flux in z-direction.

absolute error is shown in Figure 5-12. It is observed that the errors in the direct Green's function and the perturbation Green's function boundary integral solutions are less compared to the finite-difference solution.

A further investigation of the boundary integral solution is done by considering 40 boundary elements as shown in Figure 5-13. The domain integral for the direct Green's function boundary integral solutions are evaluated by discretizing the domain using 56 triangular cells and 40 nodes at the boundary as shown in Figure 5-13. The analytical and boundary integral solutions at node numbers 1 and 5 are shown in Figures 5-14 and 5-15 respectively. The boundary integral solutions are in close agreement with the exact solutions as shown in both the Figures 5-14 and 5-15 for node numbers 1 and 5 respectively.

The deviation of the boundary integral solution from the exact solution decreases by 1.34 percent using 40 boundary elements compared to the results using 20 boundary elements. This shows that the deviations of the boundary integral solution from the exact solution do not decrease significantly by increasing the number of elements from 20 to 40 at the boundary.

The present two-dimensional direct Green's function boundary integral model is also used to simulate the one-dimensional moisture-flow problem in a soil column with the Dirichlet (Γ_1) and Neumann (Γ_2) boundary conditions as shown in Figure 5-16. The soil column is discretized using 44 boundary elements and 80 triangular

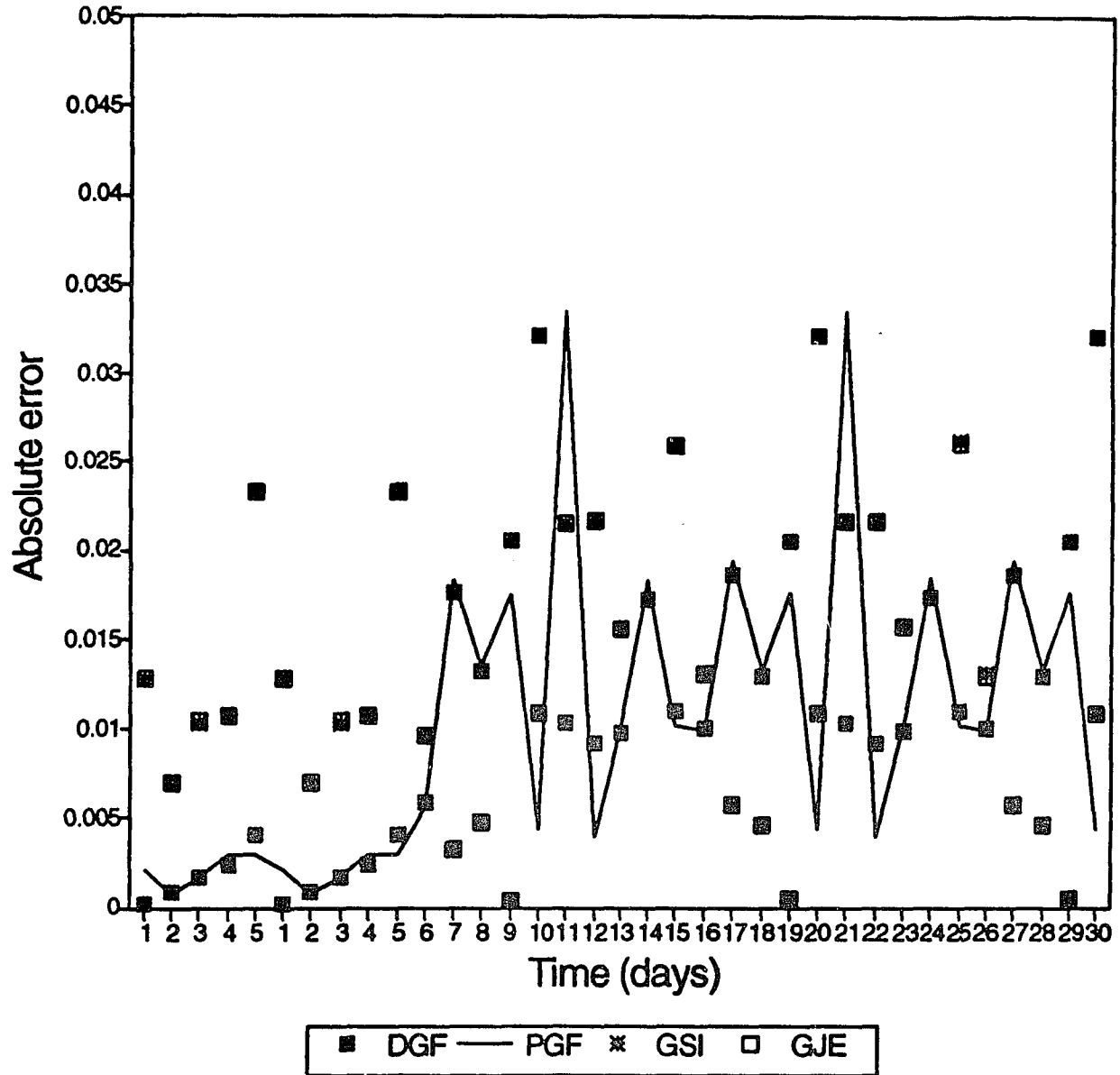
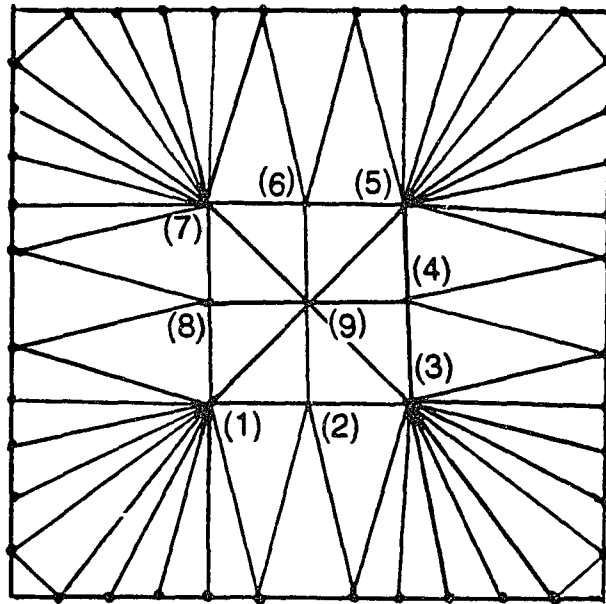


Figure 5-12. Variation of Mean Absolute Error in Boundary Integral and Finite-difference Solution.



(1) : Internal Node No. 1.

Figure 5-13. Computational Domain for Boundary Integral Solution with 40 Elements at the Boundary.

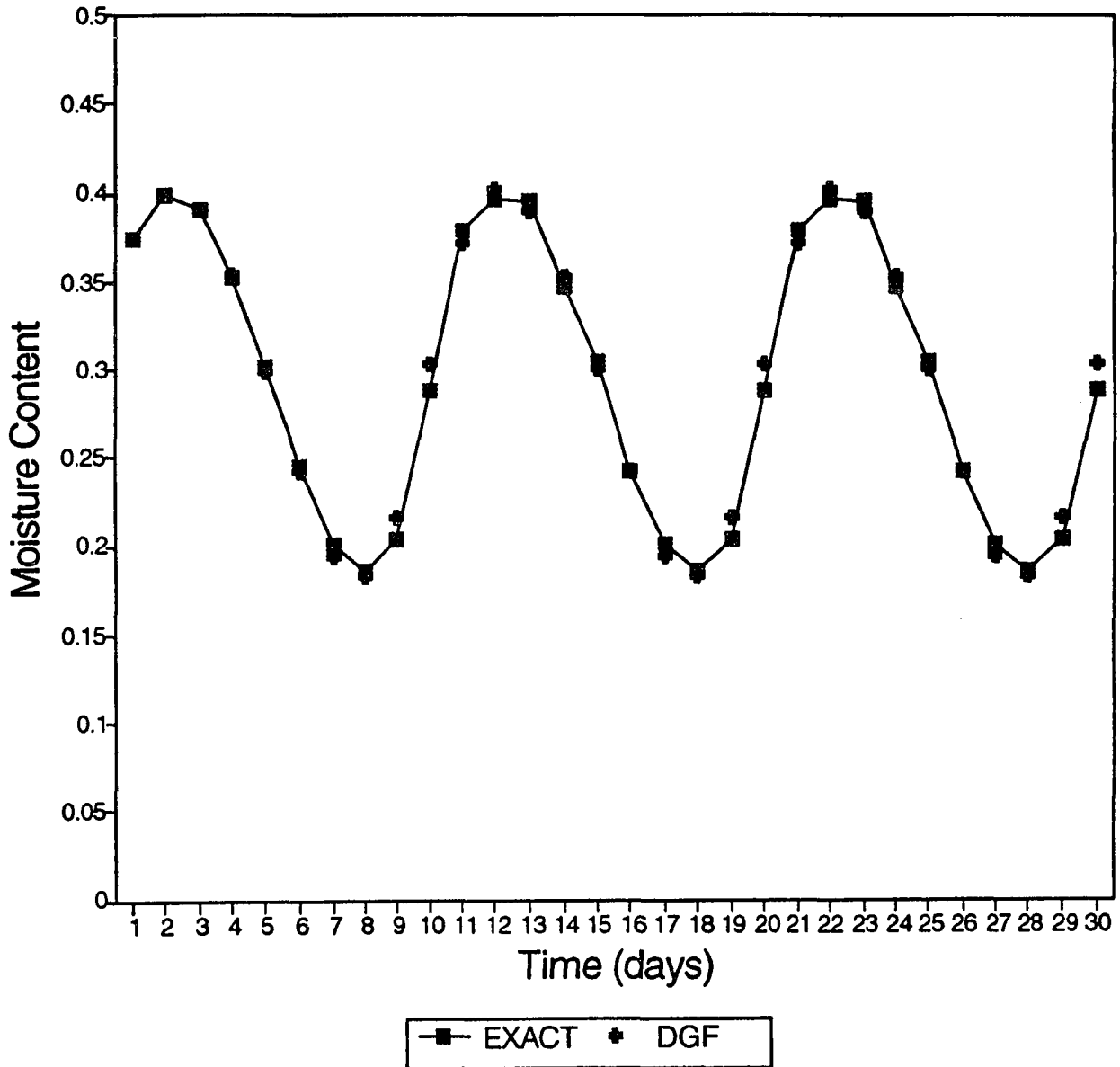


Figure 5-14. Comparison of Analytical and Boundary Integral Solution at Node Number 1 with 40 Elements at the Boundary.

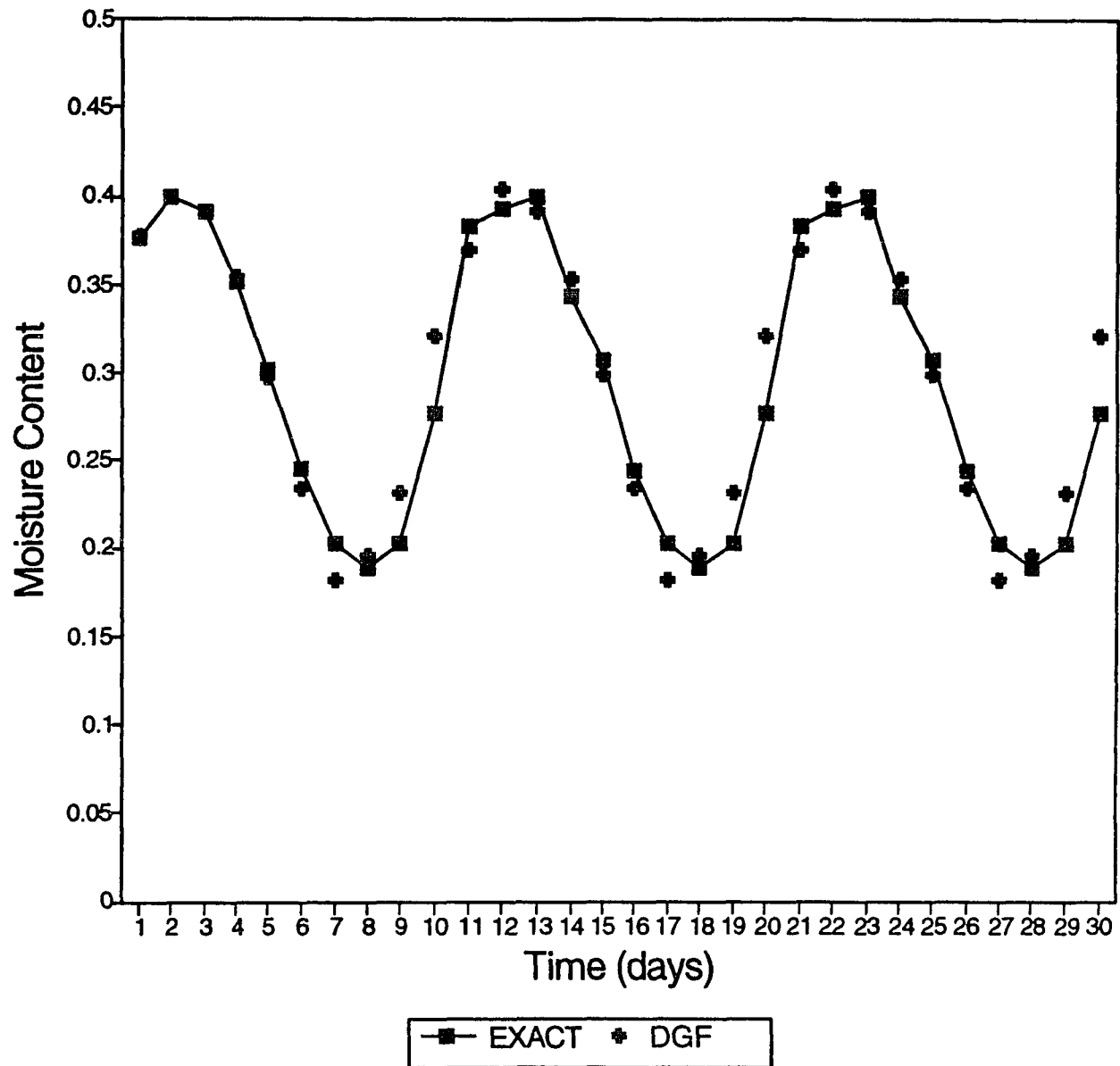
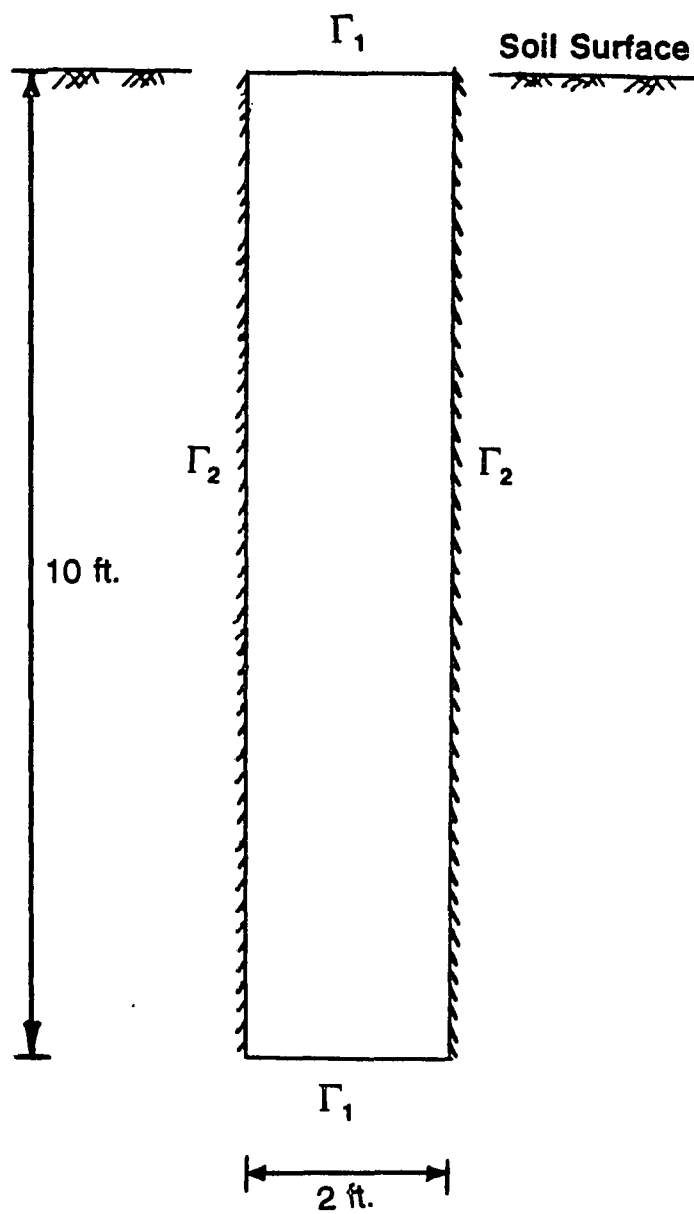


Figure 5-15. Comparison of Analytical and Boundary Integral Solution at Node Number 5 with 40 Elements at the Boundary.



Γ_1 : Dirichlet Boundary Condition

Γ_2 : Neumann Boundary Condition

Figure 5-16. Computational Domain for Comparison of Analytical and Boundary Integral Solution for One-Dimensional Moisture-Flow Problem.

cells as shown in Figure 5-17.

The exact solution in the soil column is obtained using the one-dimensional form of the moisture-flow equation (5-1) by disregarding the gravity term. The one-dimensional analytical solution is expressed in equation (5-1). The analytical and direct Green's function boundary integral solutions are obtained for 19 internal nodes in the computational domain. The comparisons of the direct Green's function boundary integral and exact solutions are made for node numbers 10 and 16 as shown in Figures 5-18 and 5-19 respectively. The variation of the moisture contents predicted by the direct Green's function boundary integral technique for the one-dimensional moisture-flow problem is in close agreement with the variation obtained by the exact solution.

A comparison of the central processing unit time corresponding to the finite-difference and boundary integral solutions is made by running the finite-difference and boundary integral models in an IBM compatible 386 25 MHz personal computer. It is observed that for perturbation Green's function boundary integral solution the central processing unit (CPU) time required is equal to 3 times the CPU time required in direct Green's function boundary integral solution. The more CPU time required in the perturbation Green's function solution is due to the repeated solution required to compute the perturbation quantities.

The boundary integral method of solution is observed to be more time consuming

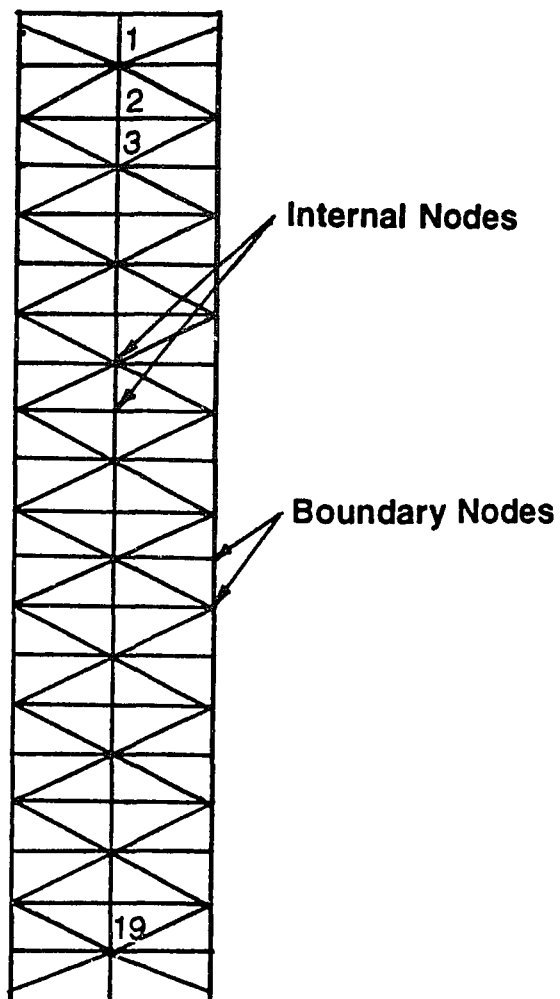


Figure 5-17. Computational Domain for Boundary Integral Solution for One-Dimensional Moisture-Flow Problem.

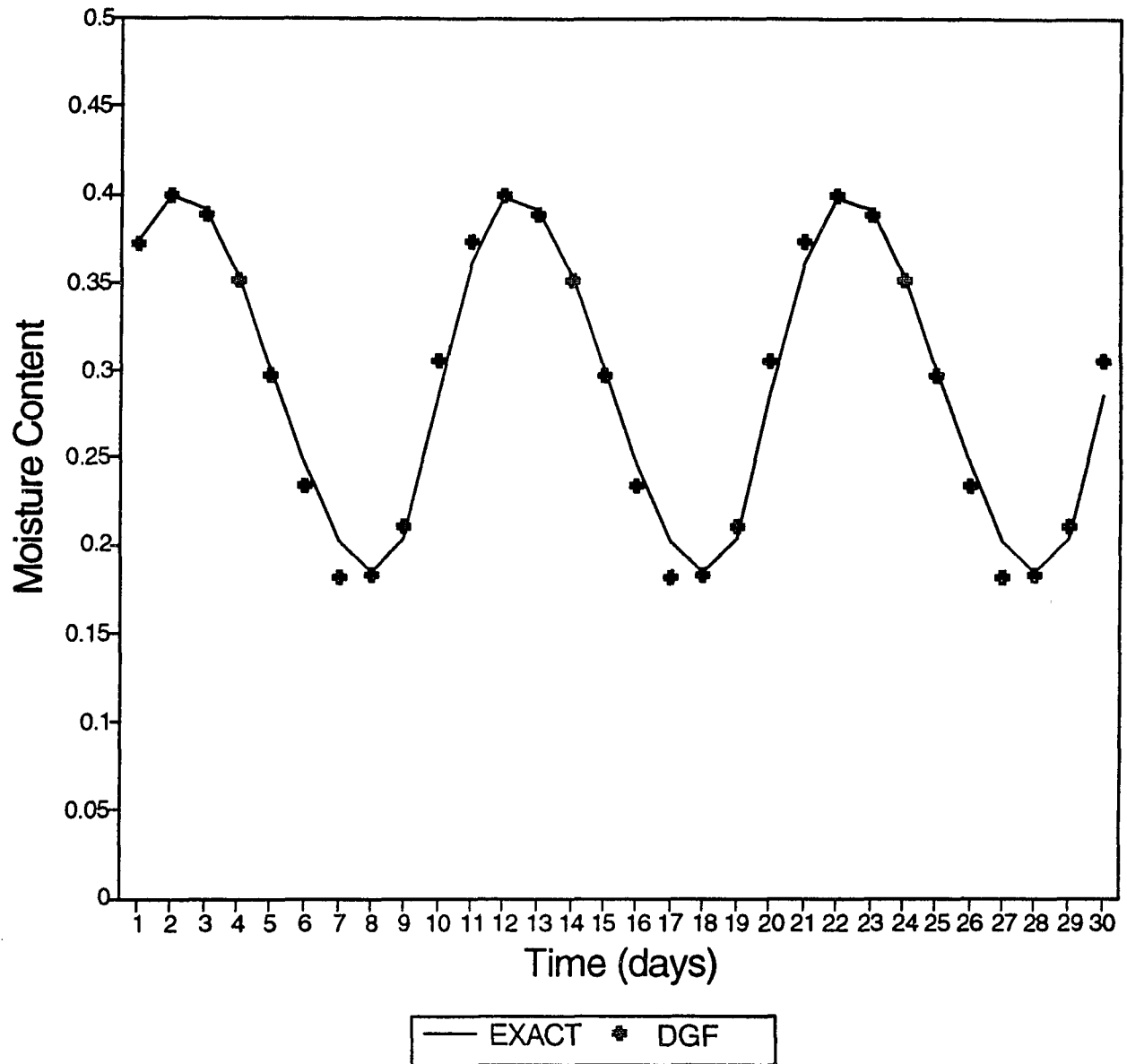


Figure 5-18. Comparison of Analytical and Boundary Integral Solution at Node Number 10 for One-Dimensional Moisture-Flow Problem.

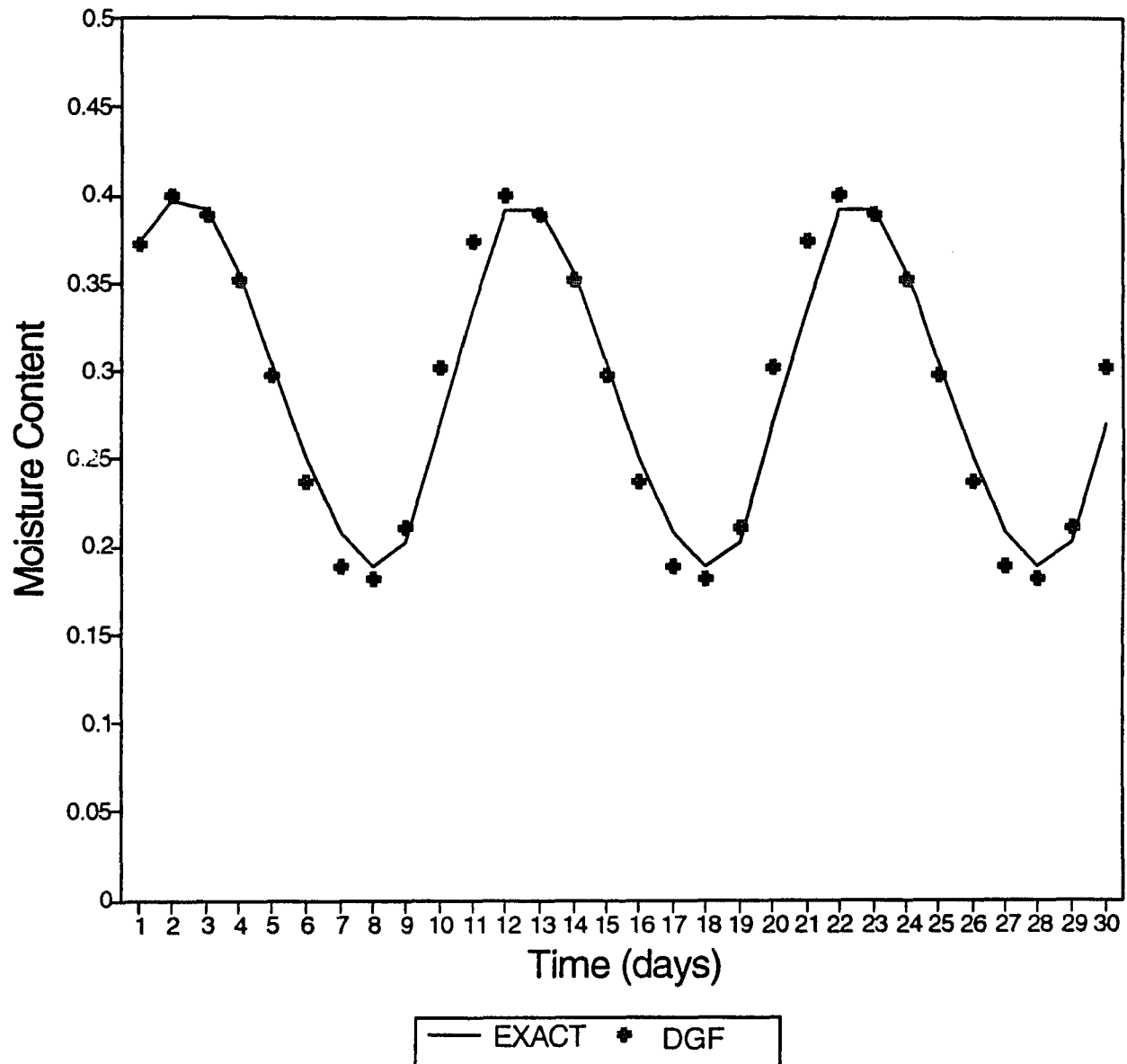


Figure 5-19. Comparison of Analytical and Boundary Integral Solution at Node Number 16 for One-Dimensional Moisture-Flow Problem.

than the finite-difference method of solution. In the direct Green's function boundary integral solution the CPU time required is 172 seconds but in Gauss-Seidel iteration finite-difference solution, the CPU time is only 10 seconds. The more CPU time required in the boundary integral computation is due to more mathematical operations needed to obtain the solution.

The input data file required more computer storage in the boundary integral scheme than a regular finite-difference grid network. The more computer storage requirement in the boundary integral method is due to storing coordinates of the boundary and internal node points. The schematization of the domain can not be avoided for the unsaturated leachate flow problem in the boundary integral method. The domain is discretized into triangular elements and the geometric data of the elements are also stored for the evaluation of the area integrals in the boundary integral formulation. However, in a regular finite-difference scheme of equal nodal areas only the space intervals Δx and Δz are stored as single values and no array is required to store the geometric data of the computational domain. This reduces the computer storage requirement in the finite-difference method.

5.2.2 Saturated Flow Problem

The simplest saturated leachate flow problem can be defined by a saturated

leachate mound in the refuse on the top of a bottom clay layer as shown in Figure 5-20. To compare with the analytical solution, the angle of inclination γ of the bottom clay layer (as shown in Figure 2-4) is considered negligible and is assumed to be zero. The thickness of the bottom clay layer is considered 2 feet with a saturated hydraulic conductivity of 10^{-7} cm/sec. A steady leachate accretion of 0.0025 ft./day is assumed. The hydraulic conductivity of the drainage layer is considered 28 ft./day.

The leachate mound head is computed at each time step for a space interval of $\Delta x = 20$ ft. The mound extended for a horizontal length of 105 ft with 20 nodes in which the leachate mound-head is computed. The two finite-difference solution techniques, Gauss-Seidel iteration and Gauss-Jordan elimination are investigated by comparing with the analytical solution.

The saturated leachate flow model was run for $\Delta t = 1$ day for a total time period of 10 days. The numerical solutions are compared for node number 8 and 10 as shown in Figures 5-21 and 5-22 respectively. Both the Gauss-Seidel iteration and Gauss-Jordan elimination methods give identical results. The solution obtained by both methods are in close agreement with the exact solution. The similarity in the solution obtained by both the iteration and elimination methods establishes the credibility of the saturated flow model for its applicability in real field conditions. Although the comparisons are made using both Gauss-Seidel iteration and Gauss-Jordan elimination methods, yet in the main FILL model Gauss-Jordan elimination

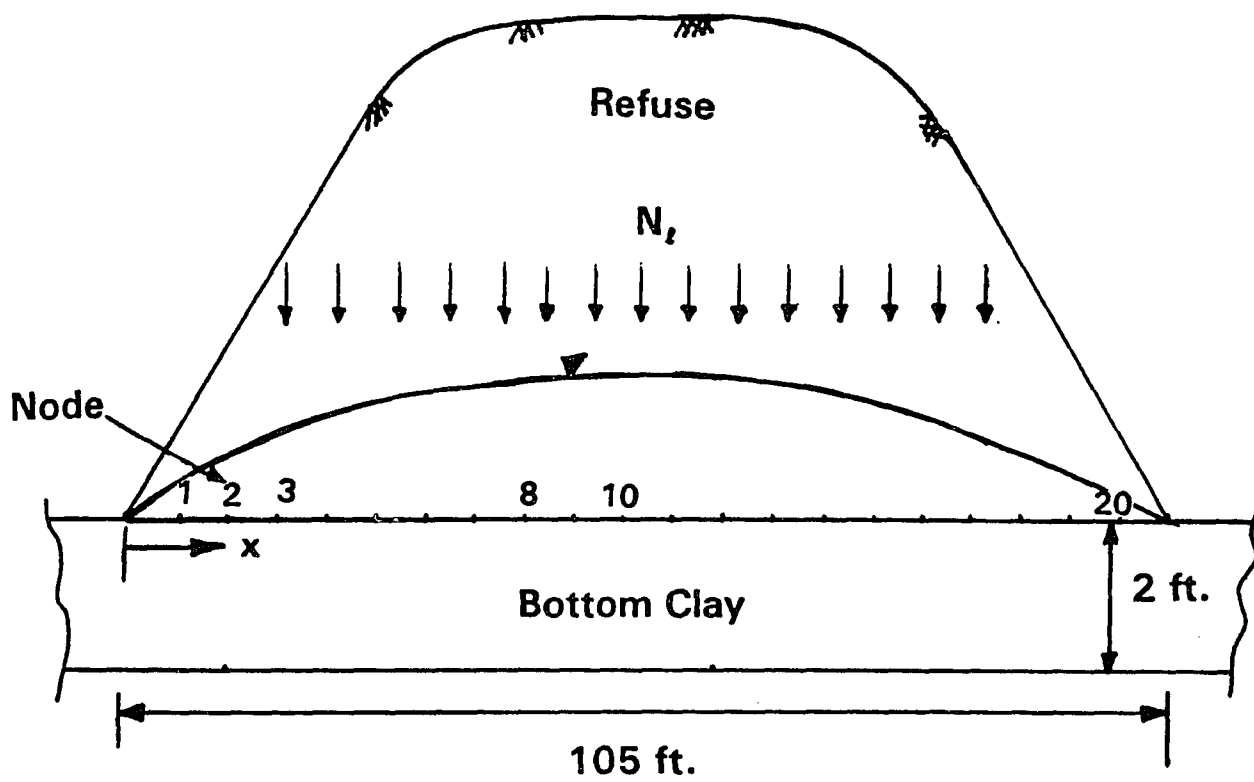


Figure 5-20. Computational Domain for Saturated Flow Problem.

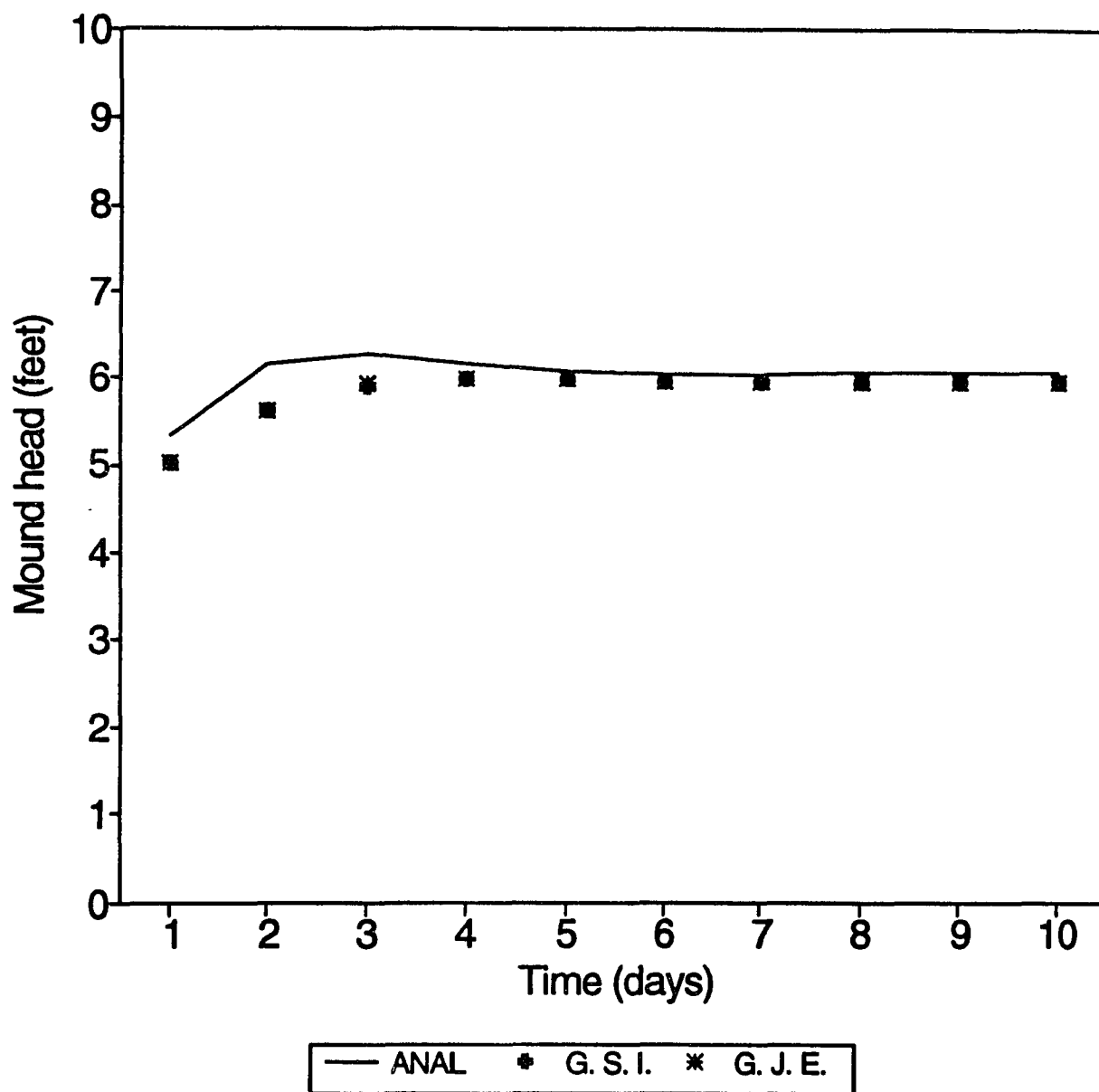


Figure 5-21. Variation of Leachate Mound-head at Node Number 8.

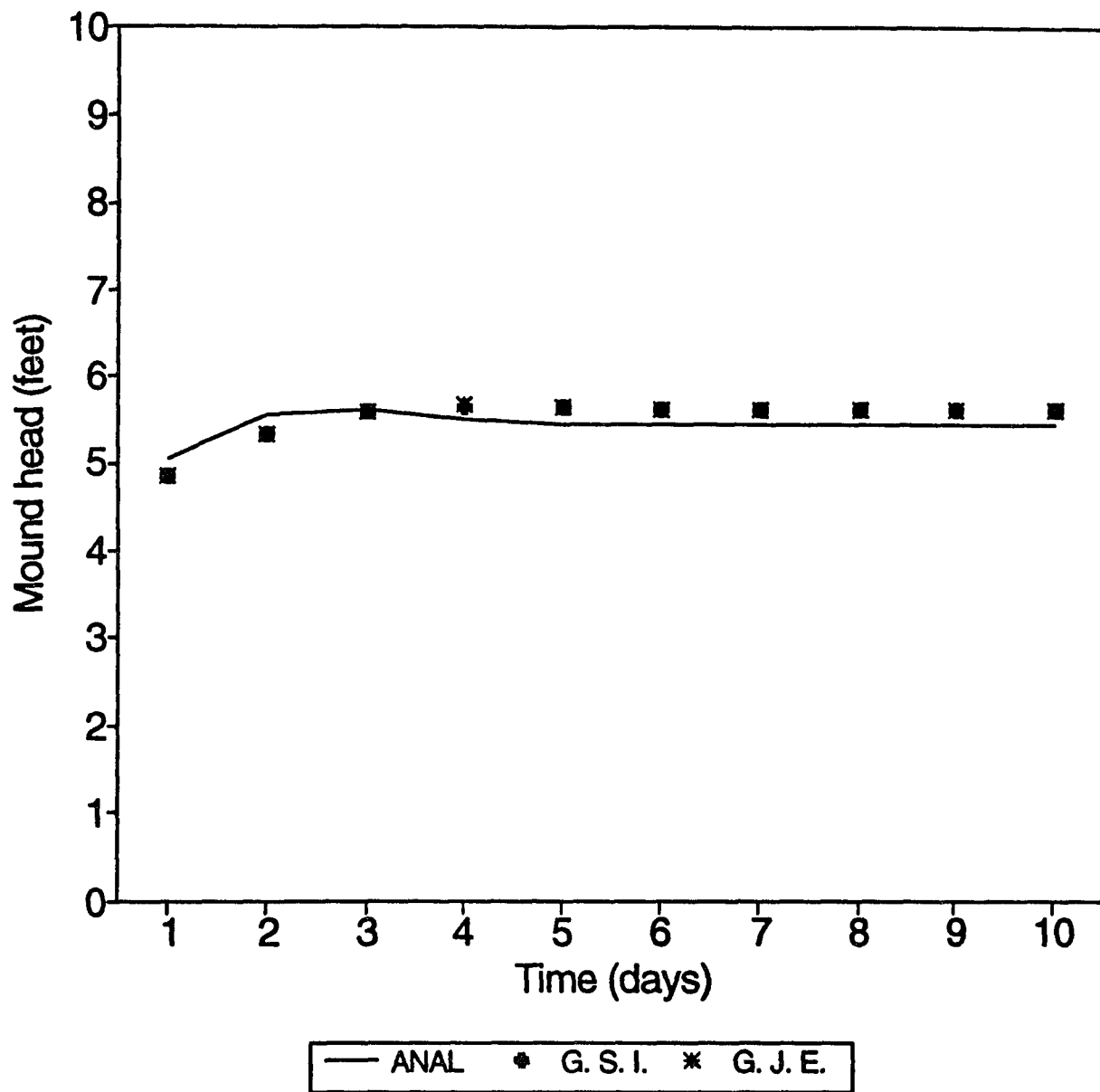


Figure 5-22. Variation of Leachate Mound-head at Node Number 10.

method is used to avoid excessive computing time required for iteration due to large accretion from the unsaturated zone.

5.3 CONCLUSIONS ON COMPARISON OF RESULTS

The conclusions following this chapter can be summarized as follows:

- The boundary integral solution using direct Green's function and perturbation Green's function boundary integral approaches show close agreement with the exact solution for the unsaturated moisture-flow equation. The close agreement of the solution by both approaches establishes the validity of the extension of the boundary integral method for solving flow problems in the unsaturated zone.
- The difference in finite-difference solution using Gauss-Seidel iteration and Gauss-Jordan elimination methods is insignificant. The close agreement of the solution by both iteration and elimination methods establishes the credibility of the finite-difference technique used in the FILL model for moisture-flow equation.
- Both the direct Green's function and the perturbation Green's function boundary integral solution show close agreement to the exact solution

compared to the results obtained by finite-difference techniques using both Gauss-Seidel iteration and Gauss-Jordan elimination methods. The perturbation Green's function solution gives more agreeable results compared to the direct Green's function boundary integral and finite-difference solution.

- Unlike the finite-difference method, extension of the boundary integral method gives satisfactory success in providing information of moisture-flux by directly differentiating the moisture content function in the domain. The boundary integral solution shows close agreement of the moisture-flux compared to the exact solution.
- The finite-difference solution using the Gauss-Seidel iteration and Gauss-Jordan elimination methods for saturated leachate flow equation shows close agreement to the exact solution. Both the Gauss-Seidel iteration and elimination methods give identical results. However, to avoid time-consuming iteration process for large leachate accretion from the unsaturated zone, it is found suitable to use Gauss-Jordan elimination method for real field simulations.

CHAPTER 6

FILL MODEL APPLICATION

The FILL model, developed using finite-difference technique, represents flow in both unsaturated and saturated zones of a landfill. The unsaturated flow model generates the time-history of leachate accretion to the saturated leachate mound. The saturated leachate flow equation is solved to predict leachate mound-head using recharge from the unsaturated zone as the source term.

The FILL model is applied to an example landfill section to demonstrate the effectiveness of the model to predict leachate flow rates in a landfill. The data are chosen from a local landfill section. Before applying the FILL model, a sensitivity test of the input parameters are done to observe the changes in the model results for the extreme values of the parameters. All these results are presented in this chapter.

6.1 SENSITIVITY TEST

Sensitivity analysis is done to see how the model reacts to extreme values of the input parameters. Variations in the value of leachate flow parameters have unequal

effects on the model results. Some of the parameters have greater effects than the others. The important input parameters for the model are saturated hydraulic conductivity, field capacity, porosity, slope of moisture characteristic curve or Campbell's parameter (b), capillary pressure head, temperature, and solar radiation.

A rigorous sensitivity analysis is done by making arbitrary changes in the aforementioned parameters. The model run was made to observe the effect on the model results for 5 percent, 10 percent, 20 percent, 35 percent, 40 percent, and 50 percent change in the values of the input parameters. The results are shown in Figures 6-1 to 6-13. Figure 6-1 shows the effect of the variation of hydraulic conductivity upon recharge to the saturated leachate mound in a landfill. Decrease in the value of hydraulic conductivity by 50 percent causes a decrease in the value of recharge by 6.31 percent as shown in Figure 6-1. The hydraulic conductivity is observed to be very sensitive to the lateral flow computed by the FILL model. Increase in the value of the hydraulic conductivity by 50 percent causes an increase of 39.35 percent in the lateral flow as shown in Figure 6-2.

The change in recharge is moderate due to the change in the value of field capacity as shown in Figure 6-3. In the FILL model, the field capacity is considered as the initial and irreducible moisture content to compute leachate accretion in the unsaturated zone. The nonlinear variation in recharge for change in field capacity is also due to the fact that the leachate accretion in the unsaturated zone of a

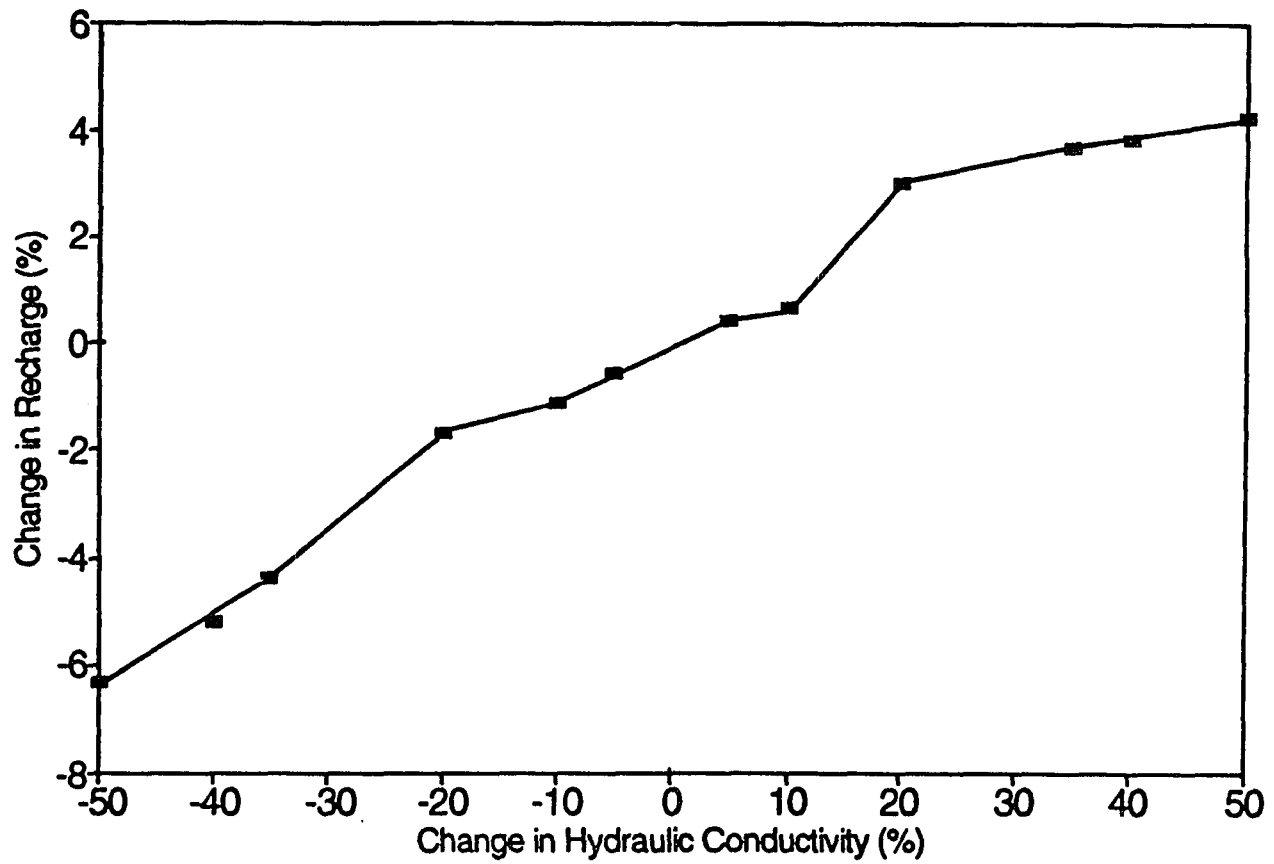


Figure 6-1. Sensitivity of Hydraulic Conductivity to Annual Recharge Computed by the FILL Model.

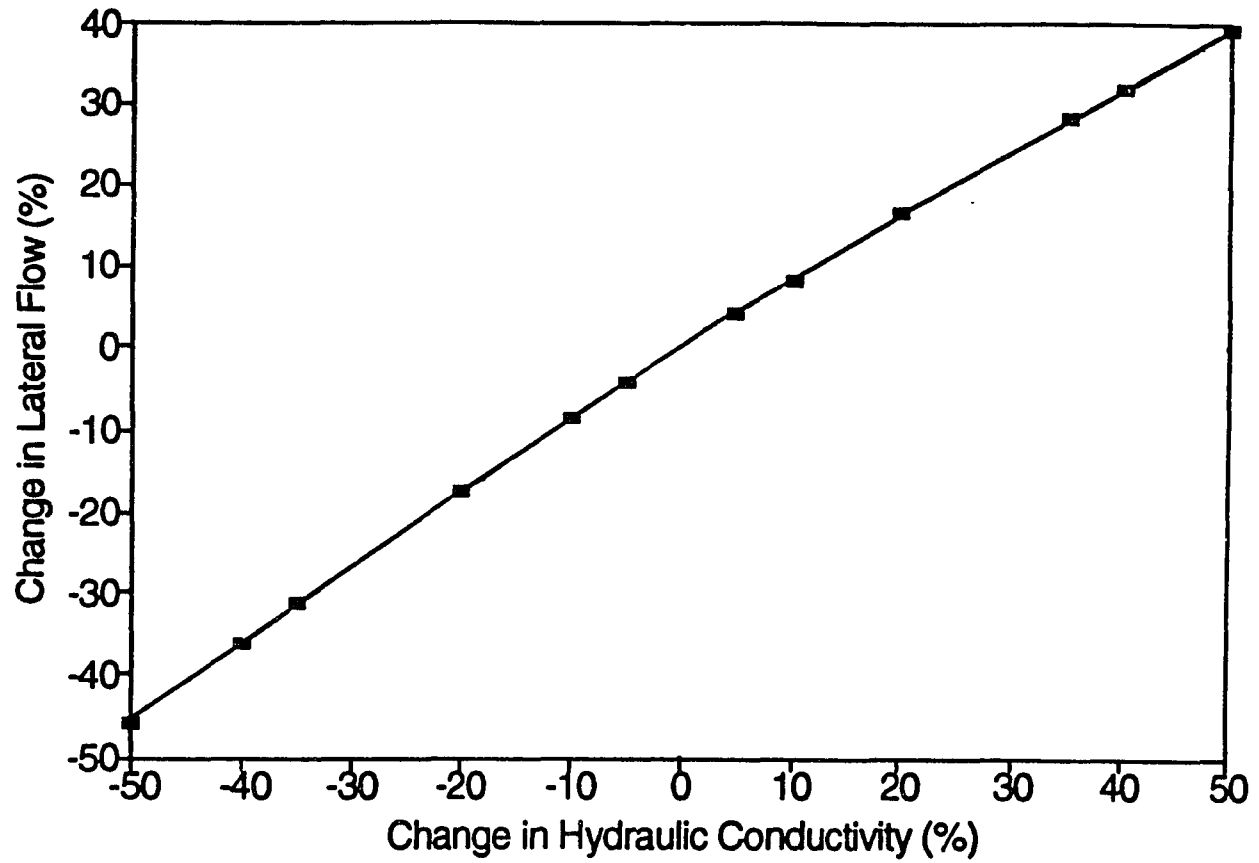


Figure 6-2. Sensitivity of Hydraulic Conductivity to Average Daily Lateral Flow Computed by the FILL Model.

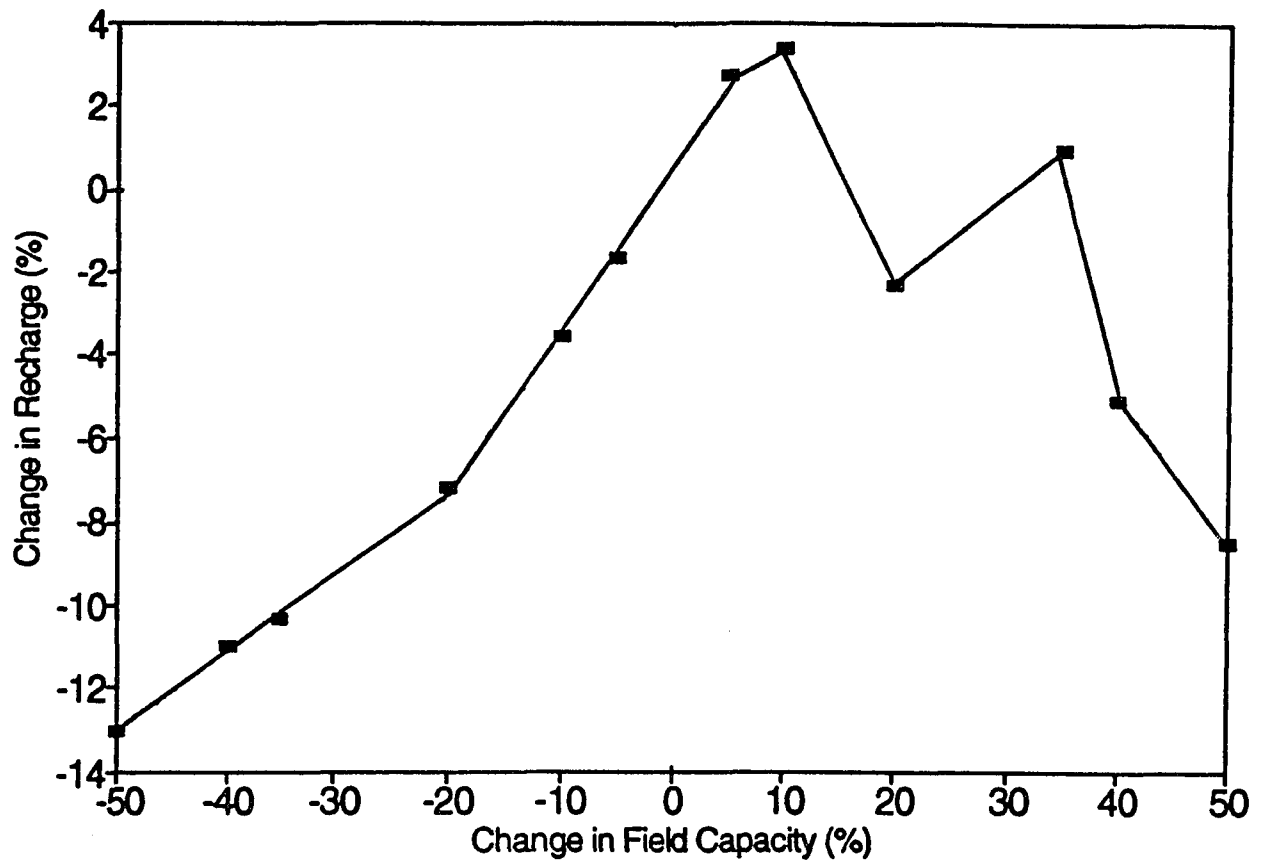


Figure 6-3. Sensitivity of Field Capacity to Annual Recharge
Computed by the FILL Model.

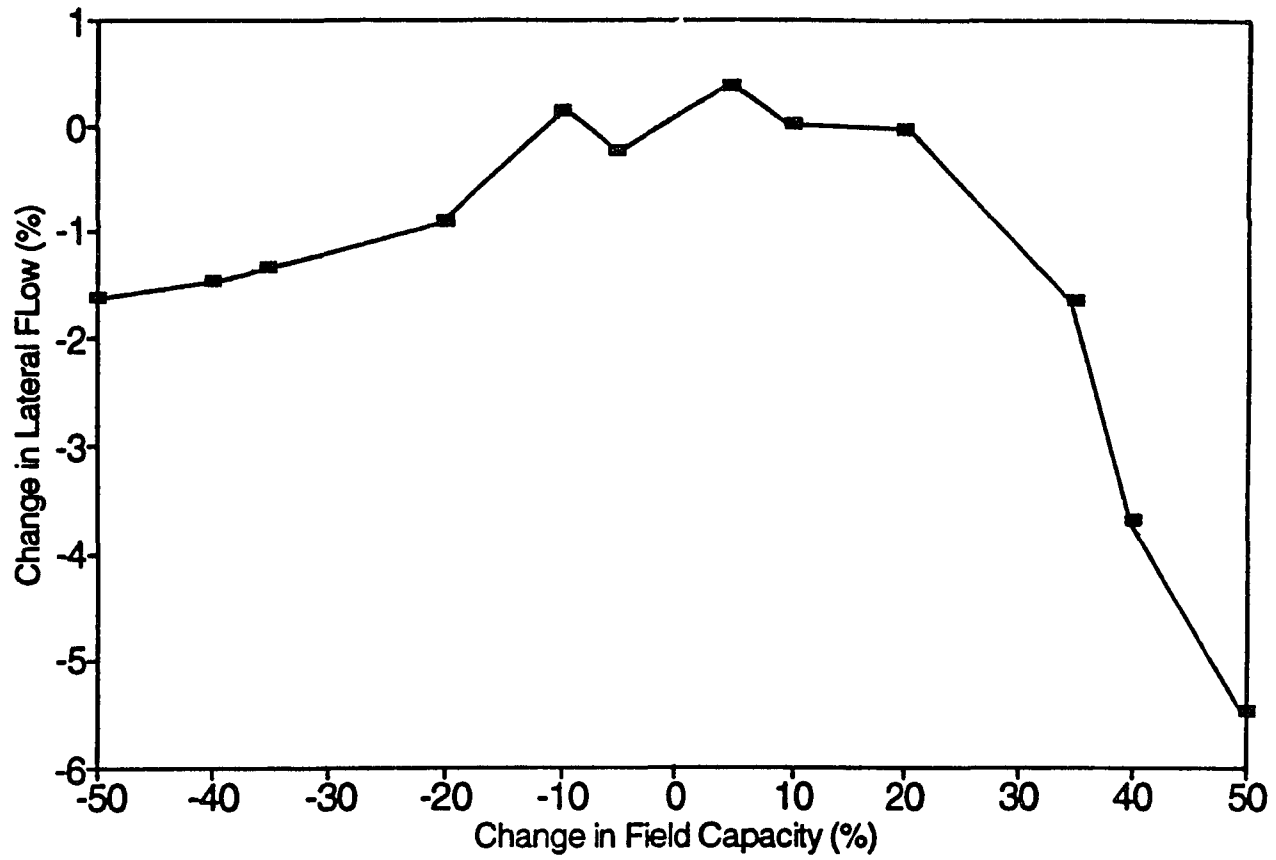


Figure 6-4. Sensitivity of Field Capacity to Average Daily Lateral Flow Computed by the FILL Model.

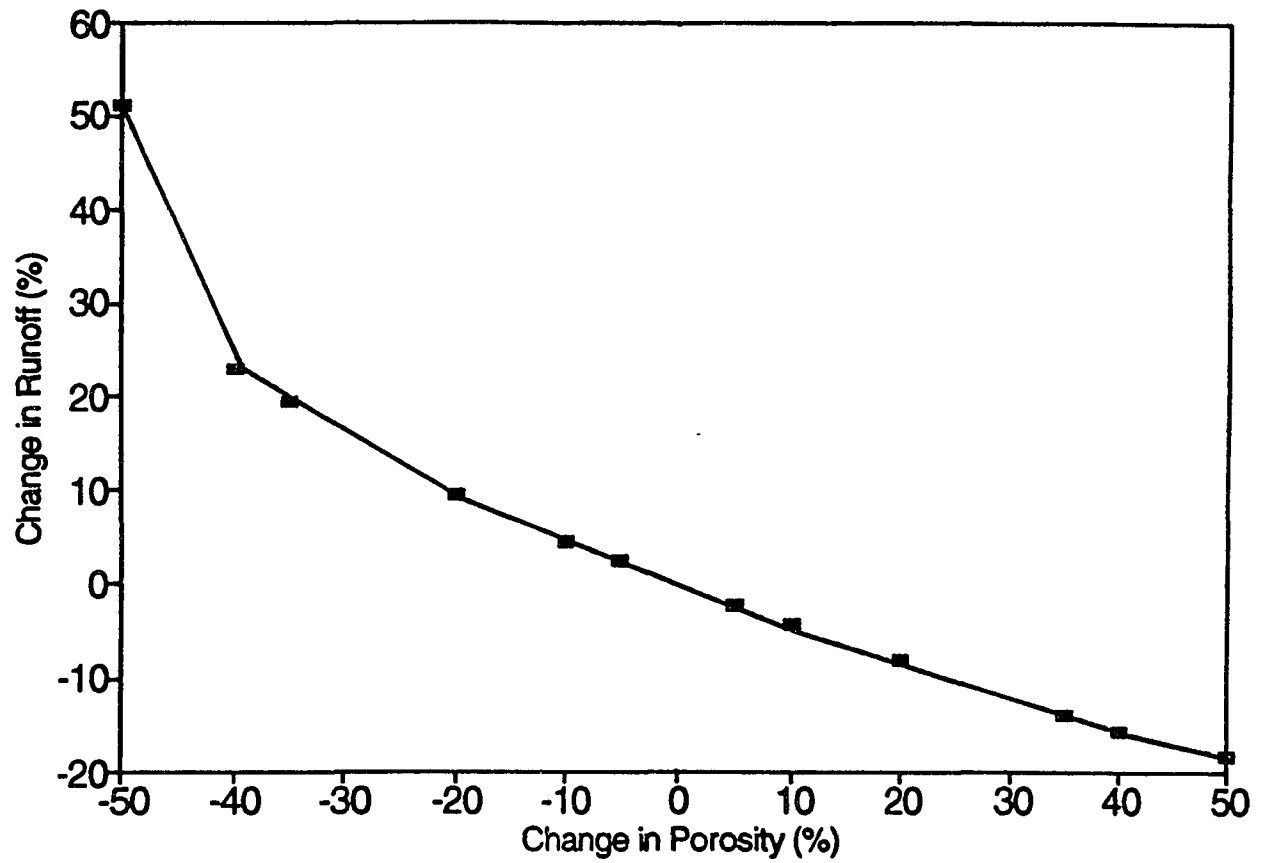


Figure 6-5. Sensitivity of Porosity to Surface Runoff Computed by the FILL Model.

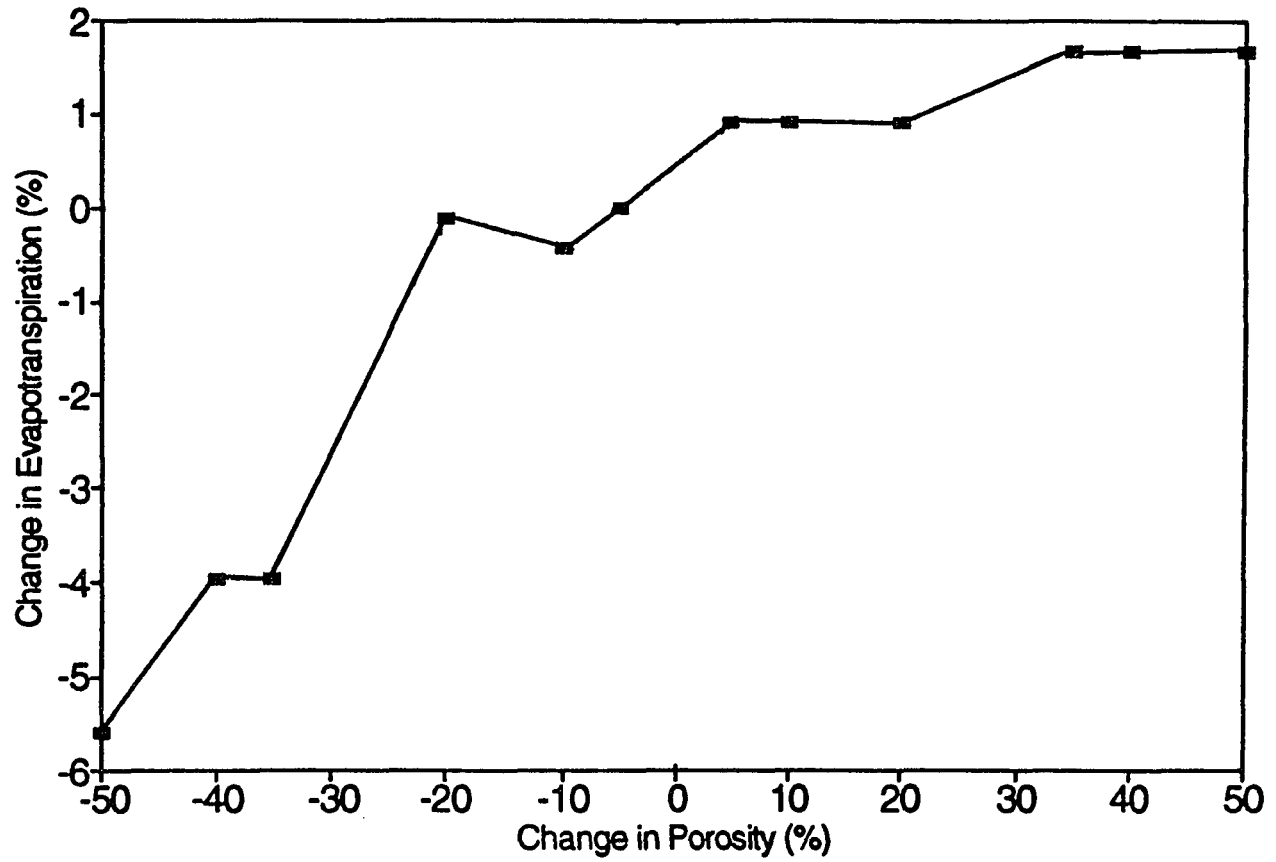


Figure 6-6. Sensitivity of Porosity to Evapotranspiration
Computed by the FILL Model.

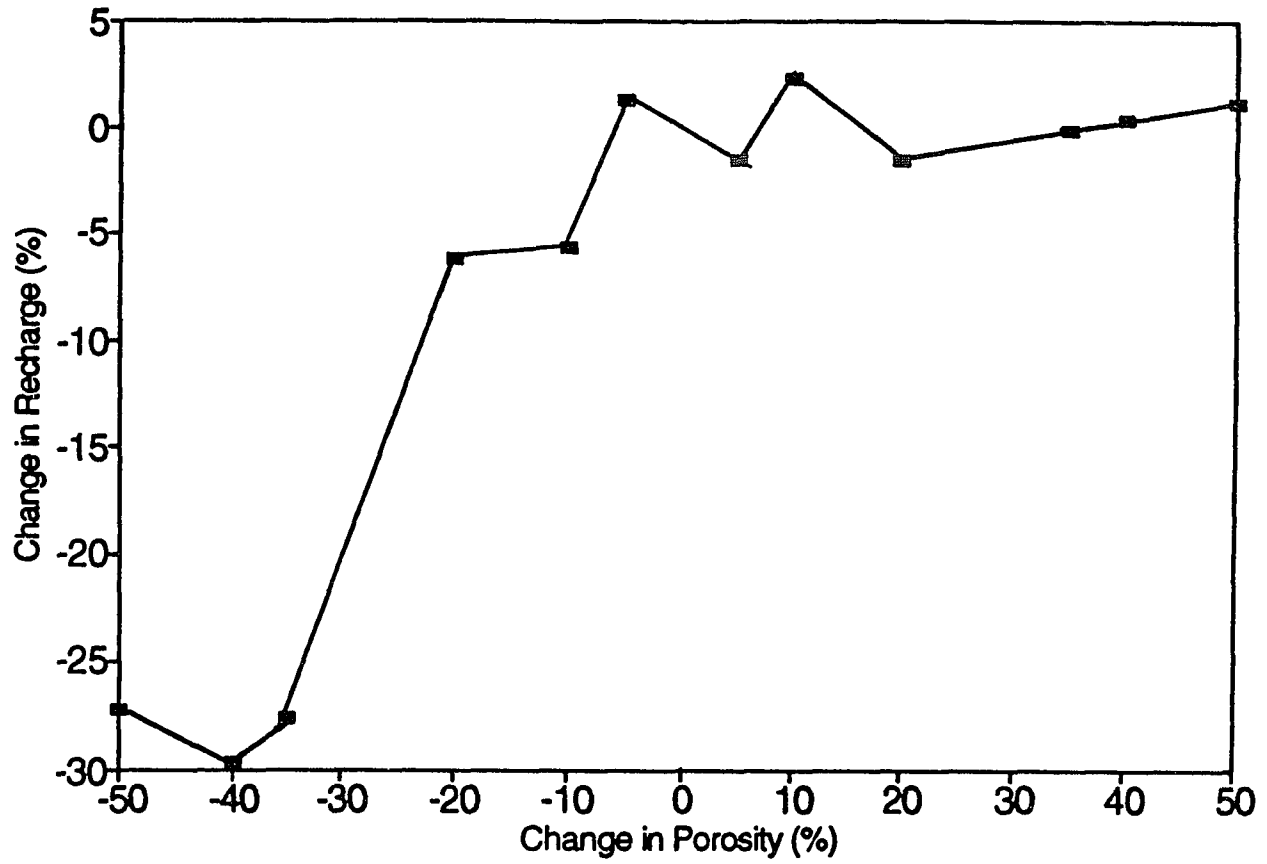


Figure 6-7. Sensitivity of Porosity to Annual Recharge Computed by the FILL Model.

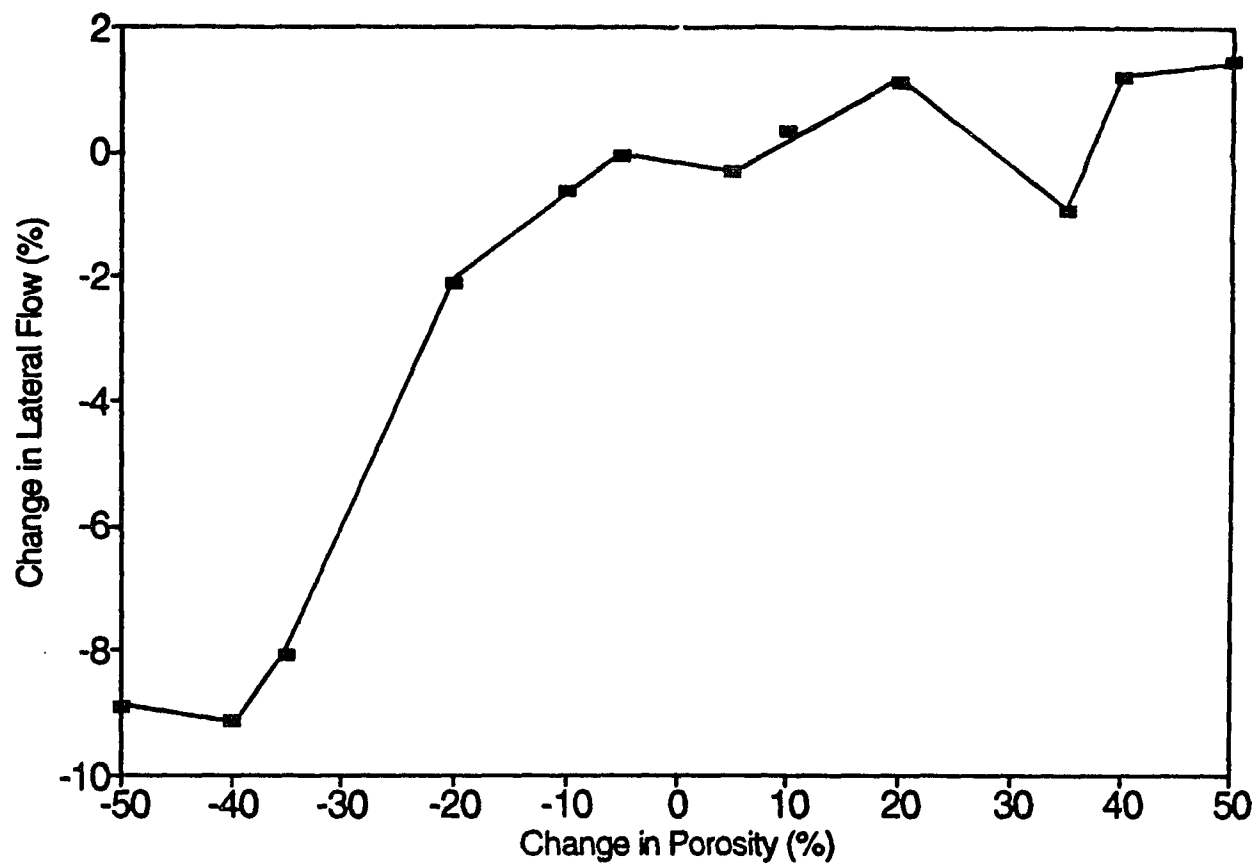


Figure 6-8. Sensitivity of Porosity to Lateral Flow Computed by the FILL Model.

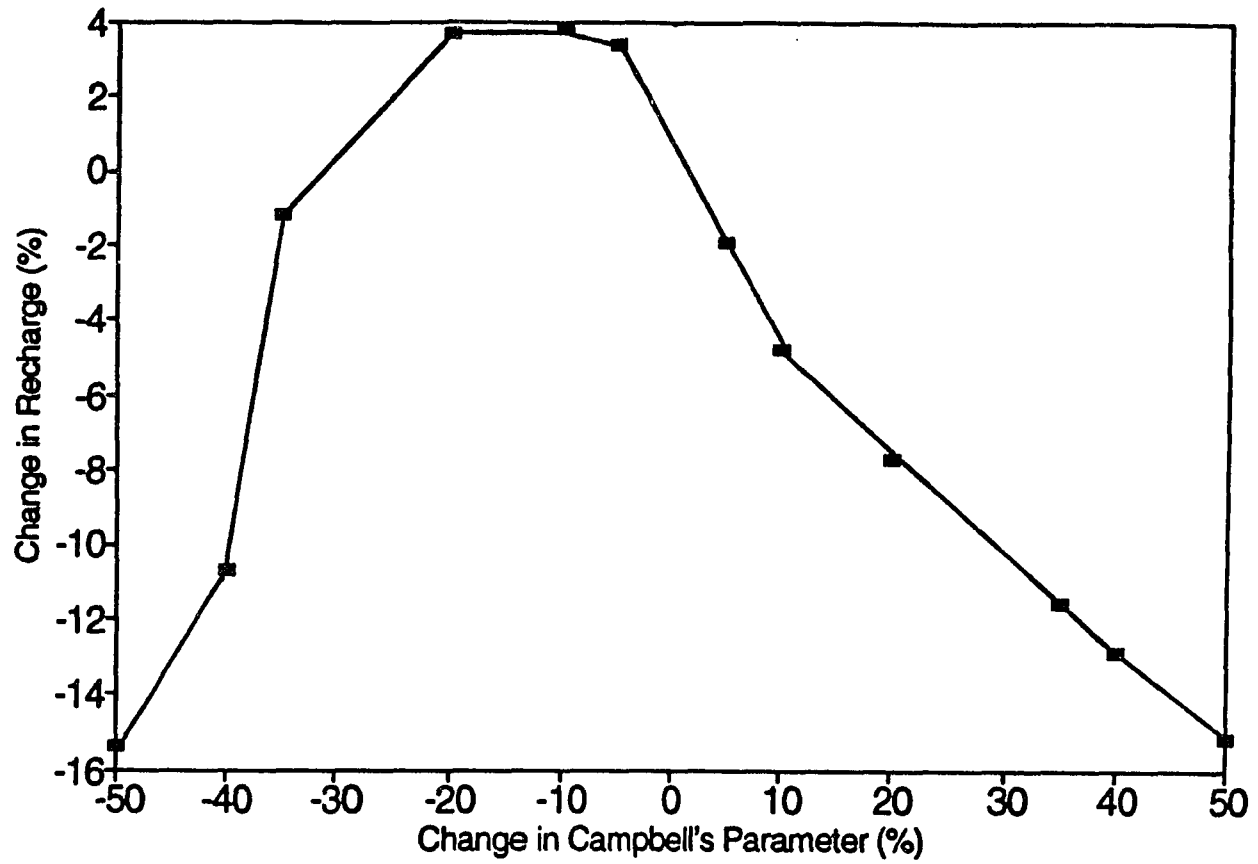


Figure 6-9. Sensitivity of Campbell's Parameter to Annual Recharge Computed by the FILL Model.

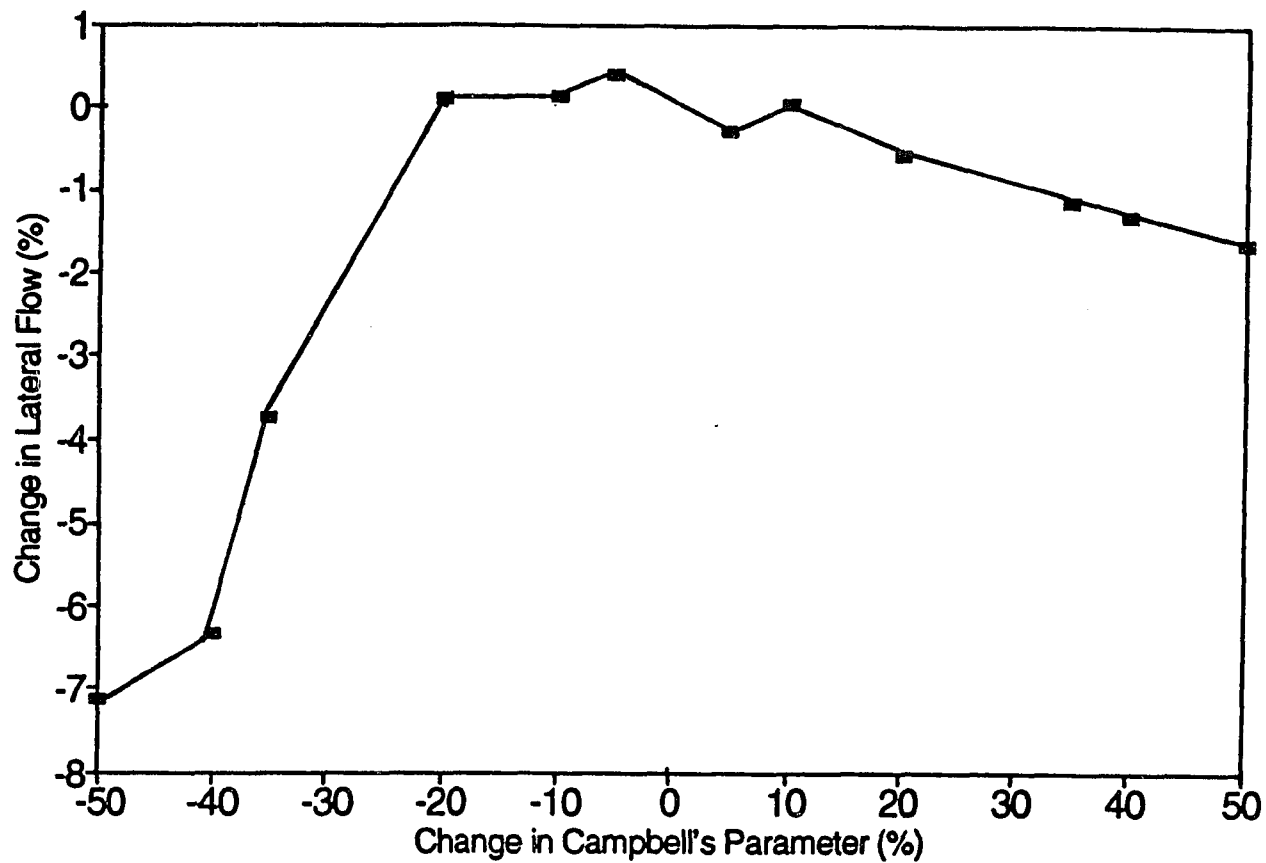


Figure 6-10. Sensitivity of Campbell's Parameter to Average Daily Lateral Flow Computed by the FILL Model.

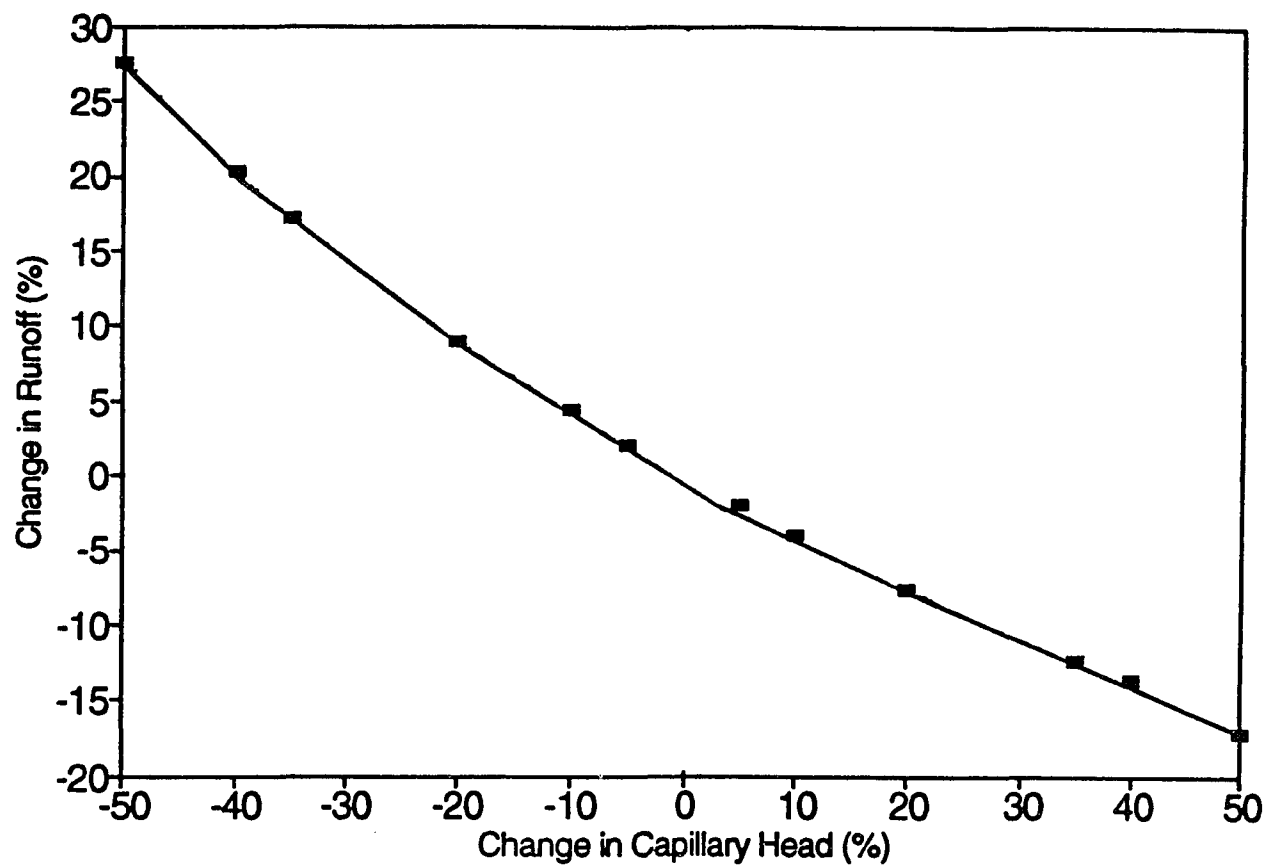


Figure 6-11. Sensitivity of Capillary Pressure Head to Surface Runoff Computed by the FILL Model.

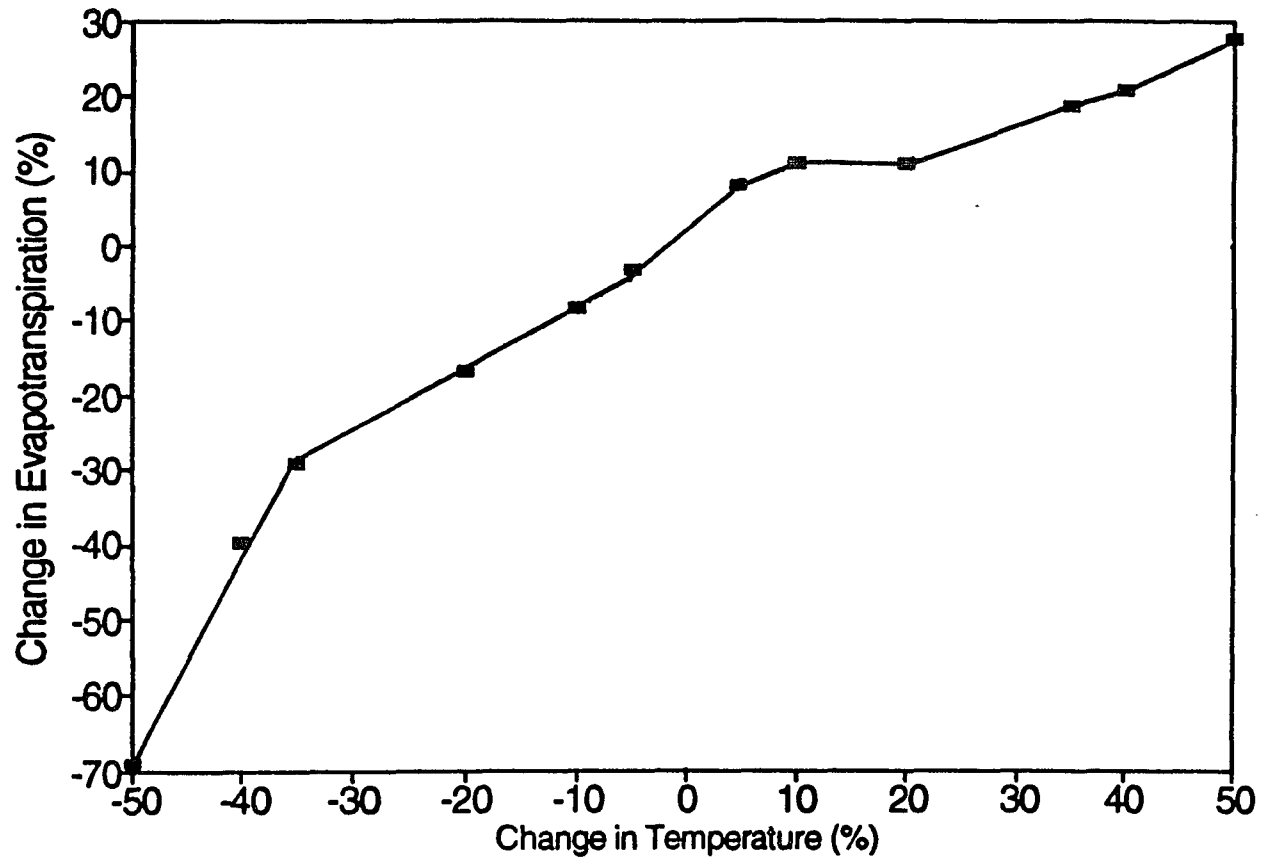


Figure 6-12. Sensitivity of Temperature to Evapotranspiration Computed by the FILL Model.

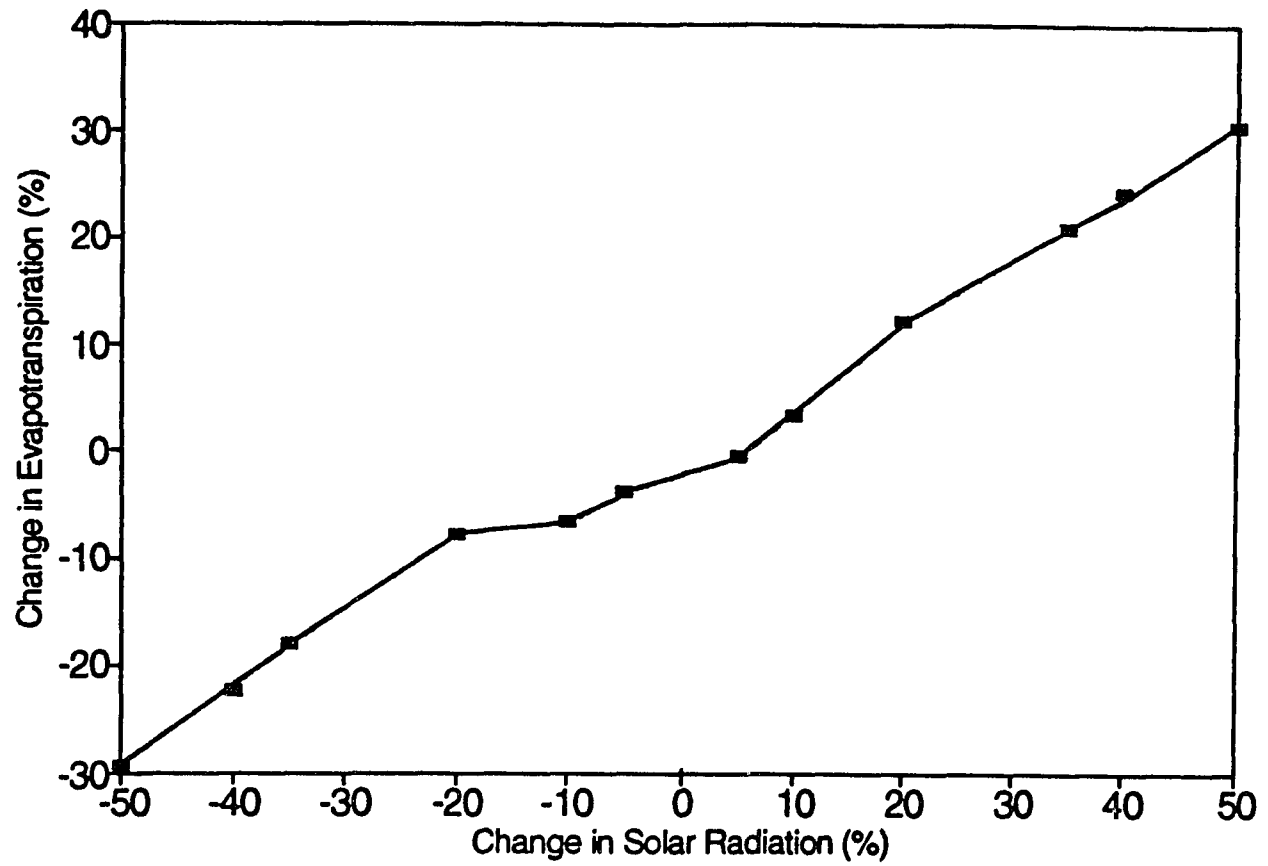


Figure 6-13. Sensitivity of Solar Radiation to Evapotranspiration
Computed by the FILL Model.

landfill is a nonlinear function of moisture content. A 20 percent decrease in the value of field capacity causes a decrease of 7.13 percent in the annual recharge to the saturated leachate mound as shown in Figure 6-3. The effect of field capacity upon lateral flow is insignificant as shown in Figure 6-4. For a 50 percent increase in field capacity, the lateral flow decreases by 5.46 percent as shown in Figure 6-4.

Porosity defines the total volume of pore space in the soil matrix that can be filled up for complete saturation. The infiltration is a function of the porosity of the top soil. The rainfall excess is the amount of precipitation less infiltration that produces surface runoff. Porosity has significant effect upon surface runoff. Decrease in the value of porosity by 50 percent causes an increase of 51.17 percent in the annual surface runoff as shown in Figure 6-5. Evapotranspiration is not very sensitive to porosity. For a 50 percent increase in the value of porosity, evapotranspiration increases by 1.68 percent as shown in Figure 6-6. Porosity is also used to define the hydraulic conductivity and diffusivity in the unsaturated zone and these properties are used to compute recharge to the saturated leachate mound. The recharge is observed to be very sensitive for decrease in the value of the porosity. Decrease in the value of the porosity by 50 percent causes a decrease of 27.16 percent in the annual recharge as shown in Figure 6-7. Porosity is also sensitive to the lateral flow computed by the model. For a 50 percent decrease in the value of the porosity causes 8.9 percent decrease in the lateral flow as shown in Figure 6-8.

Campbell's parameter (b) is used to determine the hydraulic conductivity and diffusivity of the unsaturated porous media. Campbell's parameter contribute to the nonlinearity of unsaturated hydraulic conductivity and diffusivity functions. The change in recharge to the saturated leachate mound and lateral flow due to the change in Campbell's parameter is shown in Figures 6-9 and 6-10 respectively. For a 50 percent increase in the value of the Campbell's parameter (b), the annual recharge decreases by 15.09 percent as shown in Figure 6-9. The sensitivity of Campbell's parameter (b) to the lateral flow is shown in Figure 6-10. For a 50 percent decrease in the value of the Campbell's parameter (b), the lateral flow decreases by 7.12 percent as shown in Figure 6-10. The capillary pressure head is found to be sensitive to compute surface runoff using kinematic wave equation. An increase of 27.59 percent in runoff occurs due to a 50 percent decrease in the capillary pressure head as shown in Figure 6-11.

Temperature and solar radiation values are also found to be sensitive to compute evapotranspiration using modified Penman method as shown in Figures 6-12 and 6-13. For a 50 percent decrease in the daily values of temperature causes a decrease of 68.94 percent in the total evapotranspiration computed by the FILL model. Solar radiation is less sensitive to the evapotranspiration computed by the FILL model compared to the sensitivity shown by the variation in temperature. A decrease of 29.32 percent in the total evapotranspiration occur due to a 50 percent decrease in the daily values of solar radiation.

Influence of Manning's roughness coefficient used to compute the surface runoff has insignificant effect on the model output. Wilting point is used to compute evapotranspiration by modified Penman method. It has also insignificant effect on the computation of evapotranspiration for significant changes in its values.

6.2 MODEL APPLICATION

To demonstrate the effectiveness of the model to predict seasonal variation of leachate flow rates and to estimate the average leachate flow rates, the model was applied to an example landfill section as shown in Figure 6-14. The landfill has a surface area of 376 acres. The characteristics of the landfill matrix (porosity, field capacity, wilting point, and hydraulic conductivity) and local climatological data are used to estimate the leachate flow rates in this landfill section.

The climatological data (temperature and solar radiation) were used to compute evapotranspiration and snowmelt. These data were obtained from the default values of the HELP model. Values of leaf-area index and winter cover factor were also obtained from the default values of the HELP model.

Rainfall data for seven years (1949 to 1955) were obtained from New York City department of Environmental Protection and analyzed to obtain the average rainfall year. The average annual rainfall was obtained; and the annual rainfall for the year

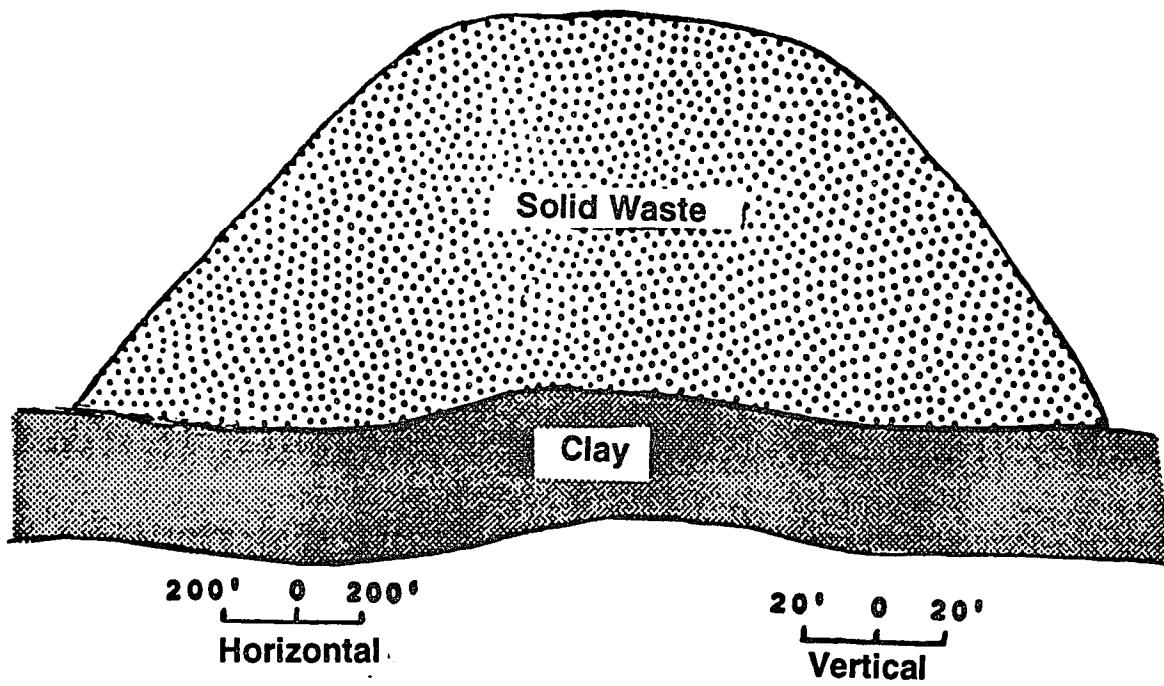


Figure 6-14. A landfill Section for FILL Model Application.

closest to the average annual value was selected as the average rainfall-year. The model was run to determine the variation of leachate flow for an average rainfall condition in the example landfill section.

The hydraulic conductivity value is obtained from the recovery test conducted by Wehran (1983) for the solid wastes at Fresh Kills landfill. The Fresh Kills landfill is located on the western shore of Staten Island, New York. The 3000-acre facility, owned and operated by the City of New York, has been receiving household wastes for disposal since 1947. The landfill has been developed into four distinct mounds, which correspond to the areas designated as Sections 1/9, 2/8, 3/4, and 6/7. The hydraulic conductivity for solid waste in Section 6/7 of Fresh Kills landfill was 0.02 cm/sec using the recovery test conducted by Wehran (1983). This value of hydraulic conductivity is used in the model to compute moisture content and leachate flow rates. A six-inch daily cover is considered on top the solid wastes with an average hydraulic conductivity value of 10^{-4} cm/sec. The model was run by considering the field capacity as the initial condition. The information to run the model for the example section is summarized in Table 6-1. The vertical landfill section was schematized into a network of finite-difference nodal areas as shown in Figure 6-15. There are a total of 64 grid-points.

The monthly variation of precipitation, runoff, evapotranspiration, and net infiltration in the landfill area is shown in Figure 6-16. The monthly values of precipitation for the average rainfall year, shown in Figure 6-16 are computed from the hourly

TABLE 6-1
DATA REQUIRED FOR THE FILL MODEL

Area	376 acres
Hydraulic conductivity:	
Daily cover	10^{-4} cm/sec (0.2835 ft./day)
Waste inside the landfill	0.02 cm/sec
Bottom clay	10^{-7} cm/sec
Capillary pressure head	4.0 inches
Porosity of the waste	0.417
Field capacity of the waste	0.05
Wilting point of the waste	0.02
Campbell's Parameter (b)	1.54

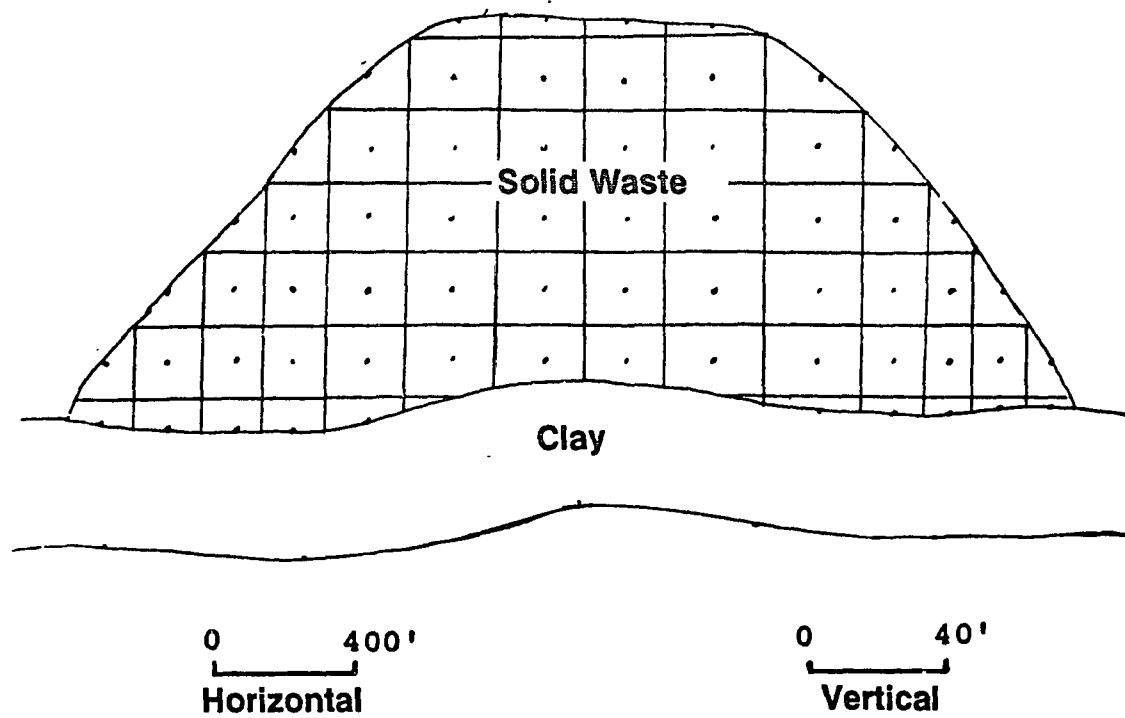


Figure 6-15. Computational Domain for FILL model.

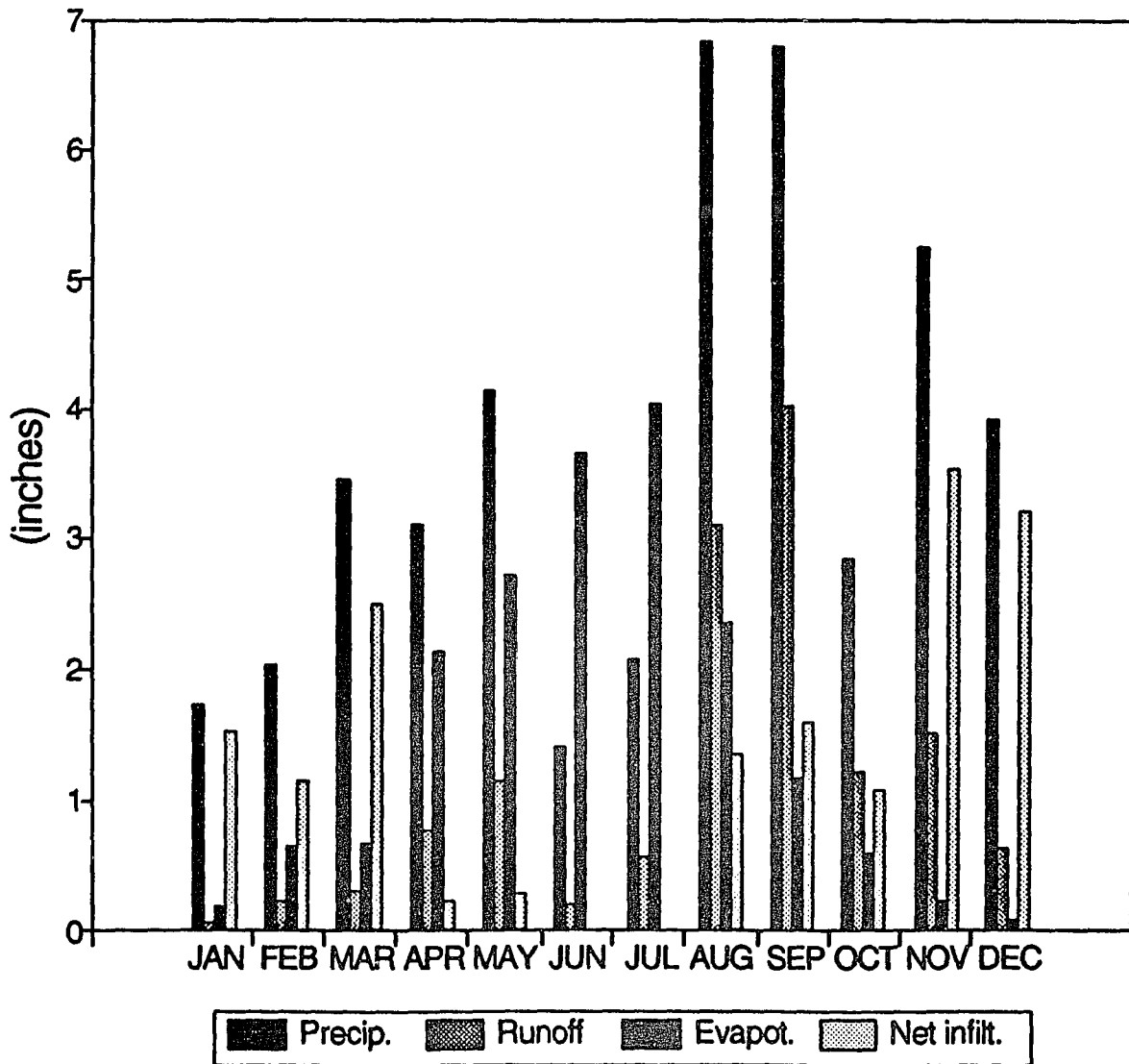


Figure 6-16. Variation of Precipitation, Runoff, Evapotranspiration, and Net Infiltration for an Average Rainfall Year.

rainfall values collected from New York City Department of Sanitation. Based on the hourly values, the runoff submodel runs for 24 hours a day when the temperature is above freezing (32°F) and when there is rainfall and/or snowmelt on the landfill surface. Figure 6-16 shows that a peak runoff value of 4.01 inches occurs during September when the rainfall is 6.78 inches. The evapotranspiration submodel computes the daily values of evaporation and transpiration using the modified Penman method described by Schroeder et al. (1984). In this estimation, evapotranspiration depends on surface runoff, infiltration, and moisture content of the soil. It is observed that the higher evapotranspiration occurs during July compared to the rest of the period in the average rainfall-year. The total annual surface runoff and evapotranspiration are computed at 13.7 inches and 18.45 inches respectively in the year when the rainfall was 43.6 inches. A higher amount of surface runoff is due to considering the side slope along the landfill surface.

The net infiltration is the precipitation less evapotranspiration and surface runoff. The variation of average daily values of net infiltration is shown in Figure 6-17. The total annual net infiltration of 11.45 inches occur, which is 26.26 percent of the total annual rainfall during the average rainfall-year. The recharge is the leachate accretion that reaches the saturated leachate mound. This is computed from the moisture content by solving the unsaturated leachate flow equation. The variation of average daily recharge is shown in Figure 6-18. The variation of recharge is consistent with the variation of net infiltration into the landfill as shown in Figures 6-17 and 6-18. Figure 6-17 shows that there is no net infiltration during June and

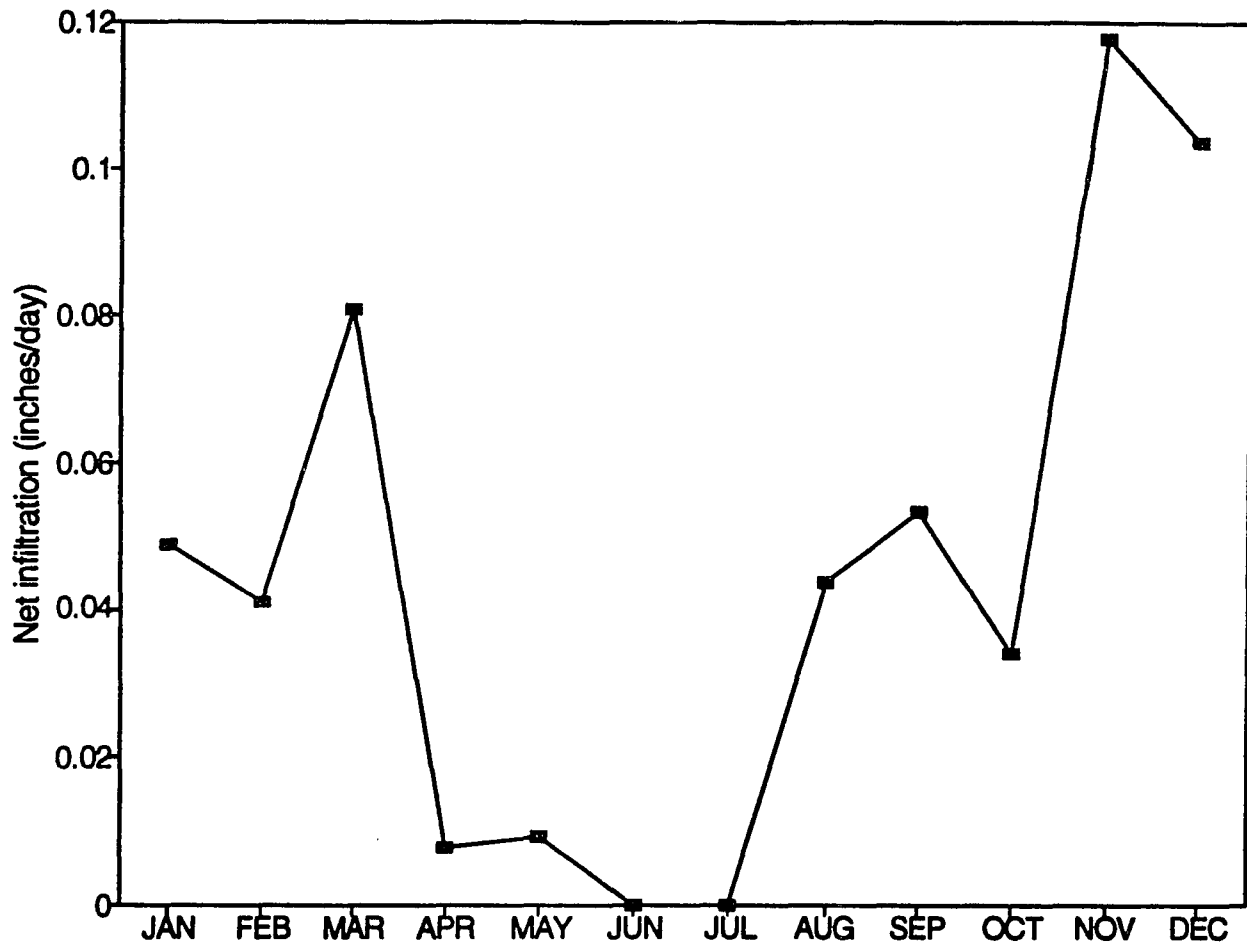


Figure 6-17. Variation of Daily Net Infiltration Computed by the FILL Model.

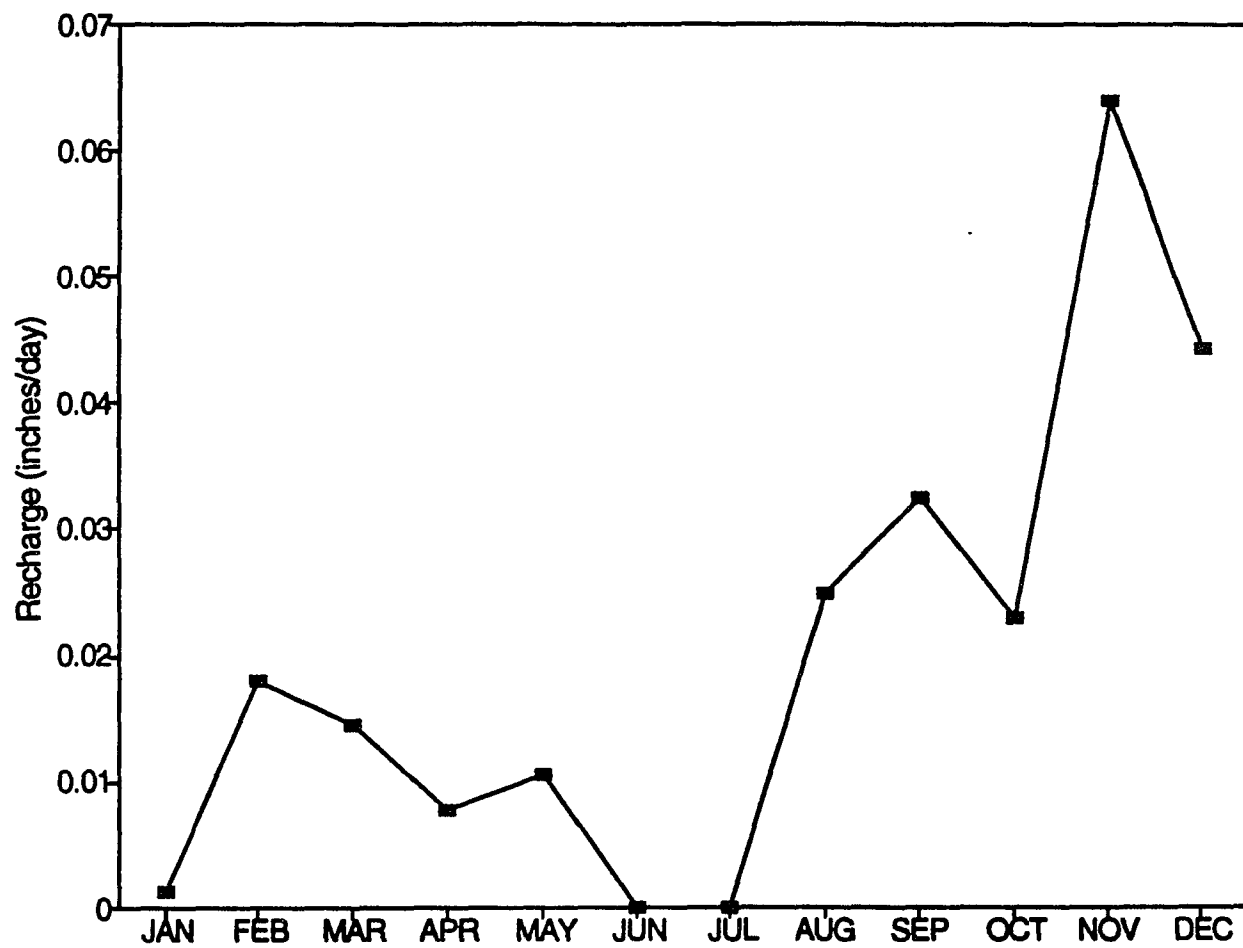


Figure 6-18. Variation of Daily Recharge Computed by the FILL Model.

July. During this time, the evapotranspiration remains higher than the total rainfall and no water is added to the retained moisture content inside the landfill. Due to the higher amount of evapotranspiration, the soil moisture content remains close to the field capacity and no recharge occurs as shown in Figure 6-17.

The outflows from the saturated leachate mound are the lateral flow towards the perimeter or the collection system and the vertical flow through the clay at the bottom of a landfill. The variation of average daily lateral flow and vertical flow is shown in Figures 6-19 and 6-20 respectively. An average lateral flow of 0.4233 mgd occurs during the average rainfall-year in the example landfill section. The vertical component of leachate flow is small compared to the lateral flow. The small value of daily vertical flow is due to the small hydraulic conductivity (10^{-7} cm.sec) of the bottom clay that restricts leakage through the bottom of the landfill. An average vertical flow of 0.0428 mgd is obtained from the model results.

The total daily flow is the sum of the lateral and vertical flow in mgd and the variation of total flow with time for the example landfill section is shown in Figure 6-21. The total average flow varies from 0.3728 mgd to 0.5977 mgd with an average daily flow of 0.466 mgd.

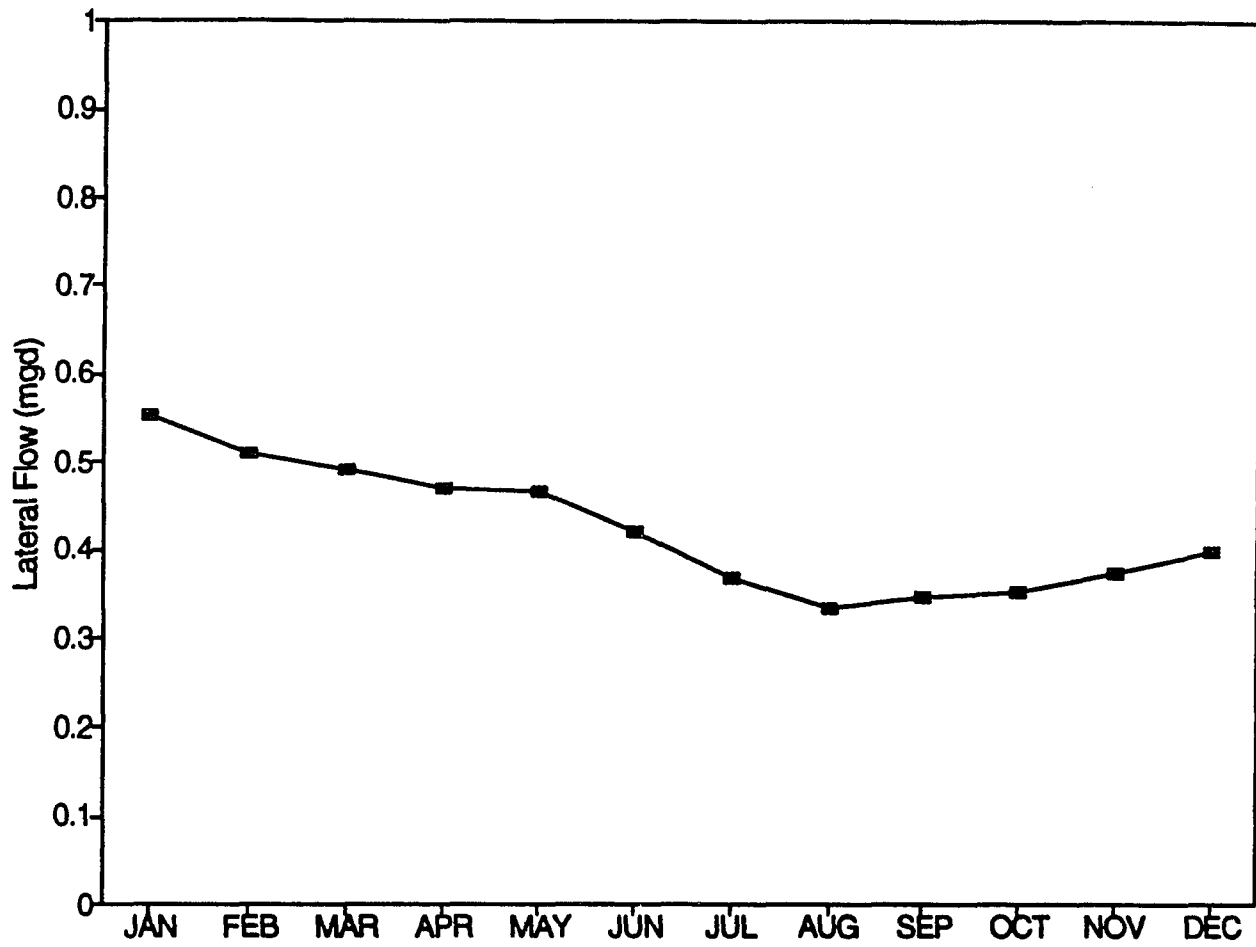


Figure 6-19. Variation of Average Daily Lateral Flow Computed by the FILL Model.

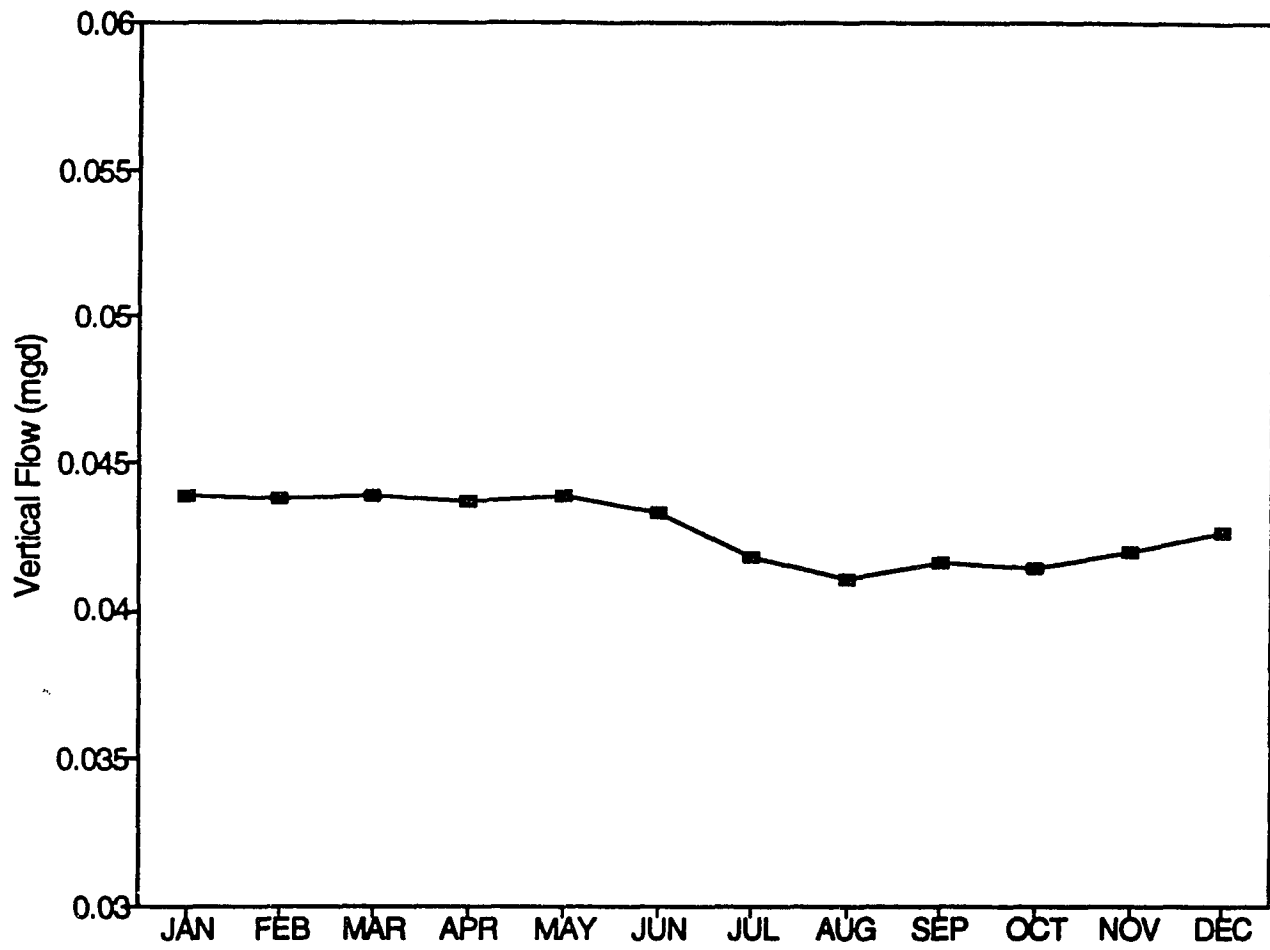


Figure 6-20. Variation of Average Daily Vertical Flow Computed by the FILL Model.

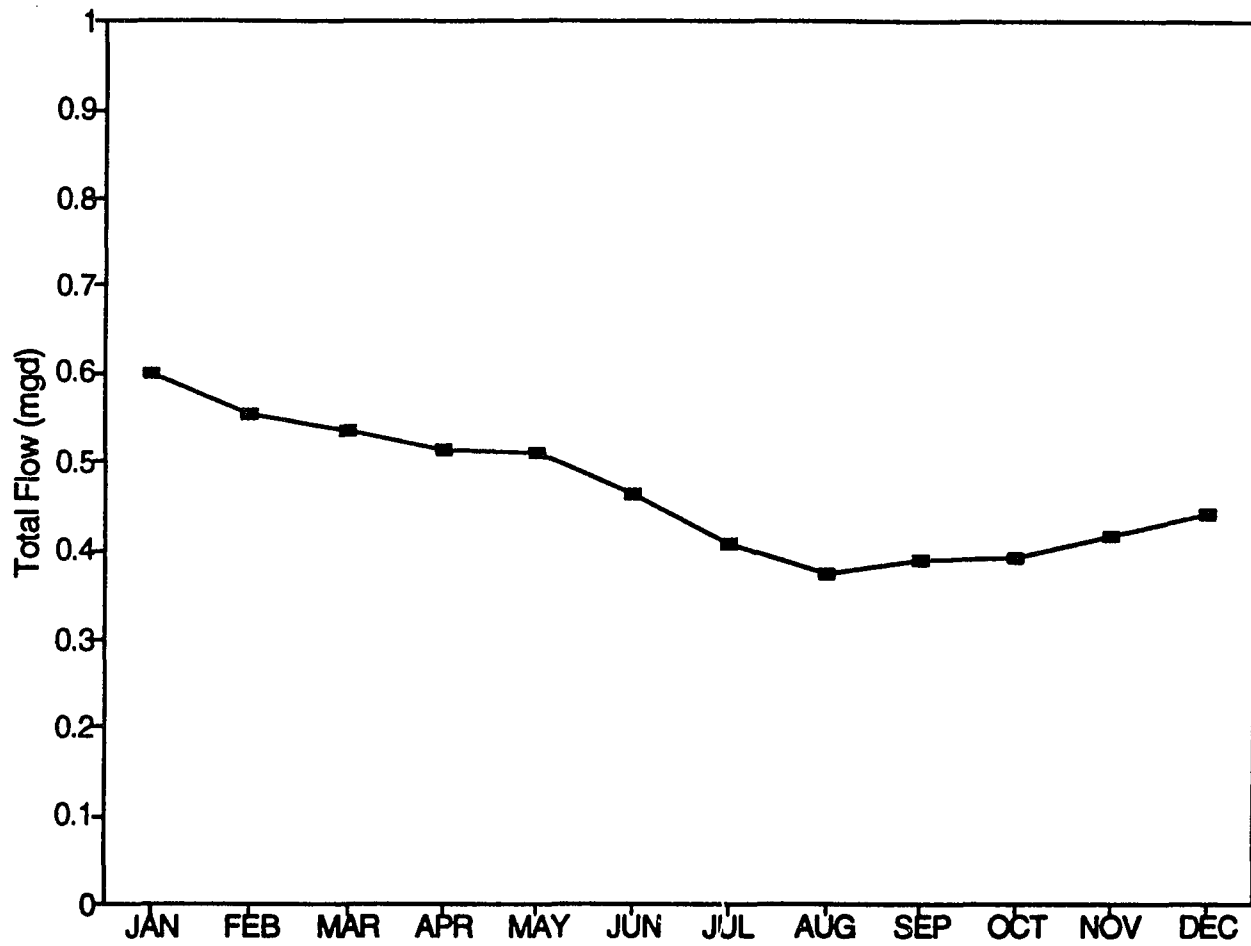


Figure 6-21. Variation of Average Daily Total Flow Computed by the FILL Model.

CHAPTER 7

DISCUSSION, CONCLUSIONS AND RECOMMENDATIONS

7.1 DISCUSSION AND CONCLUSIONS

In this dissertation, the leachate flow process into, through, and out of a solid waste landfill have been investigated. The leachate flow problems are described in both the unsaturated and saturated zones of a landfill. The unsaturated zone constitute a major portion of the landfill matrix. On-site leachate accretion occurs primarily in the unsaturated zone of a landfill. Saturated leachate mound exists at the bottom of a landfill. Leachate mound is caused by the restriction of moisture flow by a natural impermeable clay layer or an artificial clay liner beneath the landfill.

The governing equations in both the unsaturated and saturated zones in a landfill are described by partial differential equations obtained by considering mass balance and Darcy's Law. The governing equations are expressed in terms of moisture content in the unsaturated zone and leachate mound-head in the saturated zone of a landfill.

The unsaturated leachate accretion is described by a two-dimensional unsteady

state moisture-flow equation in a vertical domain. The unsteady state flow condition is considered to achieve realistic results for the time-varying field conditions. The properties of the medium, such as unsaturated hydraulic conductivity and diffusivity change with time in the unsaturated zone due to change in the moisture content. Also, the boundary condition (boundary condition depends on precipitation, evapotranspiration, and surface runoff) gives a time-varying net infiltration into the landfill that produces an unsteady leachate accretion. The saturated leachate mound-head variation in the two-dimensional vertical domain is formulated by considering leachate accretion from the unsaturated zone and flow due to hydraulic gradient and landfill bottom slope. The FILL model incorporates the leachate flow process from the upper boundary through the unsaturated and saturated zones by solving the governing equations using the implicit finite-difference method.

Leachate flow modeling requires accurate solution of the governing partial differential equations with appropriate boundary and initial conditions. To achieve that, numerical solution techniques are investigated in this work. The finite-difference solutions for both the unsaturated and saturated zones of a landfill are compared with the exact solutions. To gain more confidence, two numerical solution techniques, Gauss-Seidel iteration and Gauss-Jordan elimination methods are used and tested for accuracy by comparing with the exact solutions.

High-quality numerical modeling also depends on adopting efficient and suitable solution scheme. The most common numerical schemes used to solve the

unsaturated leachate flow equation are developed using finite-difference and finite element methods. To study the accuracy and suitability of numerical schemes in providing information in the unsaturated zone of a landfill, two boundary integral approaches, direct Green's function and perturbation Green's function, are used and compared with the finite-difference technique.

Accurate analysis was carried out by first developing analytical solutions for the governing equations. Some simplifying assumptions are made and the method of separation of variables is used to derive the analytical solution for the unsaturated and saturated leachate flow equation. Both Gauss-Seidel iteration and Gauss-Jordan elimination methods give agreeable solution compared to the analytical solution for the unsaturated moisture-flow equation. The difference in solution using Gauss-Seidel iteration and Gauss-Jordan elimination methods is insignificant. The close agreement of the solution by both iteration and elimination methods establishes the credibility of the finite-difference model used for moisture-flow equation. The solution of the finite-difference scheme for saturated leachate flow equation is obtained using Gauss-Seidel iteration and Gauss-Jordan elimination methods. These two methods of finite-difference solutions are also tested by comparing with the exact solution and close agreement is observed.

Extension of the boundary integral method to solve the unsaturated moisture-flow equation gives satisfactory success compared to the finite-difference method in providing information such as the moisture content and the moisture flux. The use

of time-dependent Green's function helped solve the unsteady unsaturated moisture-flow equation in both space and time. The two boundary integral approaches, direct Green's function and perturbation Green's function solutions show close agreement to the exact solution in comparison with the finite-difference solution.

In the direct Green's function formulation, the gravity term $\partial K(\theta)/\partial z$ is computed by considering the moisture content from the previous time step. In the perturbation Green's function formulation, the gravity term is perturbed and it is obtained from the preceding level of solution in the perturbation series. The perturbation series solution helps represent the physical system realistically by evaluating the gravity term at the current time level. But the perturbation series solution is done at the expense of more computing time than the direct Green's function solution.

Comparison of the finite-difference and the boundary integral approaches are done by computing the differences in the solution from the analytical solution. The absolute errors in boundary integral solutions remain less compared to the finite-difference method of solution. The maximum absolute deviation from the exact solution is higher in the finite-difference solution by both the Gauss-Seidel iteration and Gauss-Jordan elimination methods than that in the boundary integral solution. An average absolute error of 0.01523 is obtained in Gauss-Seidel iteration method of finite-difference solution. The results obtained by the Gauss-Jordan elimination method is similar to the result obtained by the Gauss-Seidel iteration method in the

finite-difference solution. An average absolute error of 0.01522 is observed in the Gauss-Jordan elimination method of finite-difference solution. In the direct Green's function solution, an average absolute error of 0.01336 is obtained. An average absolute error of 0.01165 is obtained in the perturbation Green's function solution. The perturbation Green's function solution is found to give more satisfactory results than the rest of the methods used in this study.

The two-dimensional boundary integral solutions are also suitable for simulating the one-dimensional moisture-flow problems. The close agreement of the variation of the moisture contents predicted by the one-dimensional analytical and the two-dimensional boundary integral solutions shows the capability of the boundary integral techniques to solve the moisture-flow problems for a soil column where flow is occurring predominantly in the vertical direction.

The boundary integral solution is capable of generating moisture fluxes in the unsaturated flow-domain of a landfill. The internal fluxes are compared with the analytical solution and show close agreement. The internal fluxes are used to determine the leachate accretion to the saturated leachate mound. The ability of the boundary integral method to directly predict the internal fluxes is an advantage that is useful to know for solving contaminant transport problem in the unsaturated zone.

In the finite-difference method, only the solution of the dependent variable is

obtained at discrete points in the domain. The internal velocities can not be predicted by directly differentiating the function because the derivative is approximated in the finite-difference grid. However, in the boundary integral method, the function is obtained as a spatially continuous solution and the formulation of the dependent variable allows to compute the velocity by directly differentiating the function. The internal fluxes are compared with the exact solutions and close results are observed. The close agreement of the internal fluxes computed by the boundary integral method and the analytical solution is observed. The close agreement shows the ability of the boundary integral method to estimate the leachate accretion and recharge to the saturated leachate mound using the internal moisture-flux in the domain.

The performance of the finite-difference scheme for saturated leachate flow equation is investigated by comparing with the analytical solution. To establish credibility in the performance of the finite-difference scheme for saturated leachate flow equation, two numerical solution techniques, Gauss-Seidel iteration and Gauss-Jordan elimination are used and the results are compared with the analytical solution. It is observed that the numerical and exact solutions are in close agreement. The numerical solution by Gauss-Seidel iteration and Gauss-Jordan elimination methods show identical results. Gauss-Seidel iteration method requires iteration and it is observed that the number of iteration increases in the main model for large leachate accretion from the unsaturated zone. To avoid time consuming iteration process, Gauss-Jordan elimination method is found suitable

to be used in the FILL model.

Sensitivity analysis was done to examine the influence of the input parameters on the solution obtained by the model. The sensitivity analysis was done in an example computational landfill domain. Extremely small variations are observed for field capacity and wilting point to compute evapotranspiration. Porosity is observed to have a significant effect on surface runoff and annual recharge but an insignificant effect on evapotranspiration. The recharge to the saturated leachate mound depends on the medium properties such as diffusivity and unsaturated hydraulic conductivity. Both of the properties are functions of porosity. Hydraulic conductivity is also a sensitive parameter. A 50 percent decrease in its value causes a decrease of 45.61 percent in the lateral flow. Recharge also increases with the increase in hydraulic conductivity. A 4.25 percent increase in annual recharge is observed due to 50 percent increase in the hydraulic conductivity value.

The FILL model was applied to an example landfill section to illustrate its effectiveness in providing information of flow components, such as surface runoff, evapotranspiration, infiltration, and leachate flow rates in both the lateral and vertical directions. Surface runoff is computed based on the hourly rainfall data. A significant amount of surface runoff occurs from the landfill surface considering side slope and less permeable daily soil cover (hydraulic conductivity 10^{-4} cm/sec) on the landfill surface. An annual amount of 13.7 inches of surface runoff occurs

from the example landfill section for an annual precipitation of 43.6 inches.

The evapotranspiration submodel computes the daily values of evaporation and transpiration using the modified Penman method. The evapotranspiration depends on infiltration and moisture content of the top soil. The modified Penman method also considers the effect of climatological data such as temperature and solar radiation. The soil data such as field capacity, wilting point, and porosity are also used to compute the actual evapotranspiration using the modified Penman method. A total amount of 18.45 inches of evapotranspiration is obtained for the example landfill area.

The recharge is the amount of leachate that is added to the saturated leachate mound from the unsaturated zone in a landfill. The main purpose of the solution for the unsaturated flow equation is to compute leachate accretion generating recharge to the saturated leachate mound. The variation of net infiltration and recharge are compared. The pattern of the variation of net infiltration into the landfill is similar to the pattern of the variation of recharge predicted for a period of one year. An annual amount of 7.29 inches of recharge to the leachate mound occurs in a period of one year for a total rainfall of 43.6 inches.

The FILL model provides the time-varying leachate flow rates which has the practical applications in the appropriate design for treatment facilities. The model provides time-varying drainage rate from the leachate mound called lateral flow

and leakage through the bottom clay called vertical flow. An average daily lateral flow of 0.4233 mgd is obtained for the example landfill area. The vertical flow values are observed to vary only insignificantly with time. An average vertical flow of 0.0428 mgd is computed by the FILL model. The insignificant variation of the vertical flow is due to the lower value of the barrier hydraulic conductivity of 10^{-7} cm/sec in the landfill.

The finite-difference formulation for both unsaturated and saturated zones are evaluated by comparing with the analytical solution. Satisfactory performance and credibility of the model results are achieved for both the unsaturated and saturated flow problem for idealized flow conditions. In real field applications, the major uncertainties in making more accurate simulations are in the material properties such as the capillary pressure head and the coefficients related to the soil moisture diffusivity $D(\theta)$ and hydraulic conductivity $K(\theta)$. These uncertainties can only be resolved by experimental studies in landfill refuse materials.

7.2 RECOMMENDATIONS FOR FURTHER STUDY

The present single barrier layer configuration of the FILL model may be extended to include operational barrier layers considering both the unsaturated and saturated flows. The model can also be modified to consider simultaneously the final cover hydraulics, unsaturated-saturated flow through the intermediate barrier layers and

the leachate flow at the bottom of a landfill. The same governing equations for both the unsaturated and saturated flows can be solved for the zones above the intermediate barrier layers. Such a model will be more useful for practical applications in providing simultaneously the information of the landfill final cover efficiency, leachate accretion above the intermediate barrier layers, leakage through all of the barrier layers, and leachate flow rates at the bottom of a landfill.

The present boundary integral formulations have been employed to solve unsteady leachate flow in the unsaturated zone of a landfill. In the saturated zone, the same boundary integral approaches can be employed to compute the leachate mound-head. To solve the unsaturated-saturated flow problem simultaneously, the possibility is to use zone concept with a moving interface between the unsaturated and saturated zone of a landfill. The development of such a model using boundary integral method is meaningful to realistically represent the physical system describing simultaneously the unsaturated and the saturated leachate flow phenomena in a landfill.

The finite-difference and boundary element models for the unsaturated flow in a landfill were compared with the analytical solutions. Experimental investigations can also be performed by constructing landfill cells to generate data of surface runoff, evapotranspiration, moisture content, and leachate accretion. Model results can be compared with the experimental data for further verification of the numerical models.

APPENDIX A

DERIVATION OF ANALYTICAL SOLUTION FOR THE UNSATURATED MOISTURE FLOW EQUATION

The simplified form of the moisture-flow equation in the unsaturated zone is:

$$\frac{\partial \theta}{\partial t} = \hat{D} \left(\frac{\partial^2 \theta}{\partial x^2} + \frac{\partial^2 \theta}{\partial z^2} \right) - p \quad (\text{A-1})$$

Here p defines the gravity $\partial K(\theta^k)/\partial z$ as expressed in equation (5-1).

The boundary conditions as illustrated in Figure 5-1 are as follows:

$$\theta(x, 0, t) = \bar{A}$$

$$\theta(0, z, t) = \bar{B}$$

$$\frac{\partial \theta}{\partial x}(L_1, z, t) = \bar{C}$$

$$\frac{\partial \theta}{\partial z}(x, L_2, t) = \bar{D}$$

The initial condition is :

$$\theta(x, z, 0) = \theta_0$$

The solution for $\theta(x,z,t)$ can be set in the following form:

$$\theta(x, z, t) = v(x, z, t) + \bar{A} + \frac{z^2}{2L_2} \bar{D} \quad (\text{A-2})$$

Here $v(x,z,t)$ is another variable as a function of both space and time to be determined.

From equation (A-2):

$$\frac{\partial \theta}{\partial x} = \frac{\partial v}{\partial x}$$

$$\frac{\partial^2 \theta}{\partial x^2} = \frac{\partial^2 v}{\partial x^2} \quad (\text{A-3})$$

$$\frac{\partial \theta}{\partial x}(L_1, z, t) = \frac{\partial v}{\partial x}(L_1, z, t) = \bar{C}$$

$$\frac{\partial \theta}{\partial z} = \frac{\partial v}{\partial z} + \frac{\bar{D}z}{L_2} = \bar{D}$$

$$\frac{\partial^2 \theta}{\partial z^2} = \frac{\partial^2 v}{\partial z^2} + \frac{\bar{D}}{L_2} \quad (\text{A-4})$$

$$\frac{\partial \theta}{\partial z}(x, L_2, t) = \frac{\partial v}{\partial z}(x, L_2, t) + \bar{D} = \bar{D}$$

$$\frac{\partial v}{\partial z}(x, L_2, t) = 0$$

$$\frac{\partial \theta}{\partial t} = \frac{\partial v}{\partial t} \quad (\text{A-5})$$

Substituting (A-3), (A-4) and (A-5) into (A-1):

$$\frac{\partial v}{\partial t} = \hat{D} \left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial z^2} \right) + \hat{p} \quad (\text{A-6})$$

where,

$$\hat{p} = \frac{\hat{D}\bar{D}}{L_2} - p$$

The initial and boundary conditions for the dependent variable $v(x,z,t)$ in equation (A-6) are as follows:

$$v(x, z, 0) = \theta \quad (\text{A-7a})$$

$$v(x, 0, t) = 0 \quad (\text{A-7b})$$

$$v(0, z, t) = \hat{B} \quad (\text{A-7c})$$

$$\frac{\partial v}{\partial x}(L_1, z, t) = \bar{C} \quad (\text{A-7d})$$

$$\frac{\partial v}{\partial z}(x, L_2, t) = 0 \quad (\text{A-7e})$$

where,

$$\hat{B} = \bar{B} - \bar{A} - \frac{z^2}{2L_2} \bar{D}$$

$$\hat{\theta} = \theta_o - \bar{A} - \frac{z^2}{2L_2} \bar{D}$$

The solution for $v(x,z,t)$ can be written as:

$$v(x, z, t) = \sum_{n=0}^{\infty} F_n(x, t) \sin\left\{\frac{(2n+1)\pi z}{2L_2}\right\} \quad (\text{A-8})$$

Here $F_n(x,t)$ is a coefficient that depends on horizontal dimension x and time t , which is to be determined.

From equation (A-8):

$$\frac{\partial v}{\partial t} = \sum_{n=0}^{\infty} \frac{\partial F_n(x, t)}{\partial t} \sin\left\{\frac{(2n+1)\pi z}{2L_2}\right\}$$

$$\frac{\partial v}{\partial x} = \sum_{n=0}^{\infty} \frac{\partial F_n(x, t)}{\partial x} \sin\left\{\frac{(2n+1)\pi z}{2L_2}\right\} \quad (\text{A-9})$$

$$\frac{\partial^2 v}{\partial x^2} = \sum_{n=0}^{\infty} \frac{\partial^2 F_n(x, t)}{\partial x^2} \sin\left\{\frac{(2n+1)\pi z}{2L_2}\right\} \quad (\text{A-10})$$

$$\frac{\partial v}{\partial z} = \sum_{n=0}^{\infty} F_n(x, t) \frac{(2n+1)\pi}{2L_2} \cos\left\{\frac{(2n+1)\pi z}{2L_2}\right\}$$

$$\frac{\partial^2 v}{\partial z^2} = -\sum_{n=0}^{\infty} F_n(x, t) \left\{ \frac{(2n+1)\pi}{2L_2} \right\}^2 \sin\left\{ \frac{(2n+1)\pi z}{2L_2} \right\} \quad (\text{A-11})$$

The sine series expression for the nonhomogeneous term in equation (A-6) can be expressed as:

$$\hat{p} = \sum_{n=0}^{\infty} t_n(x, t) \sin\left\{ \frac{(2n+1)\pi z}{2L_2} \right\} \quad (\text{A-12})$$

From which:

$$t_n(x, t) = \frac{2}{L_2} \int_0^{L_2} \hat{p} \sin\left\{ \frac{(2n+1)\pi z}{2L_2} \right\} dz \quad (\text{A-13})$$

Substituting equations (A-10) to (A-13) in equation (A-6), the following equation is obtained:

$$\frac{\partial F_n}{\partial t} = \hat{D} \left[\frac{\partial^2 F_n}{\partial x^2} - \left\{ \frac{(2n+1)\pi}{2L_2} \right\}^2 F_n \right] + t_n(x, t) \quad (\text{A-14})$$

From equation (A-8):

$$v(0, z, t) = \sum_{n=0}^{\infty} F_n(0, t) \sin\left\{ \frac{(2n+1)\pi z}{2L_2} \right\}$$

Therefore,

$$F_n(0, t) = \frac{2}{L_2} \int_0^{L_2} \theta \sin\left\{\frac{(2n+1)\pi z}{2L_2}\right\} dz = \beta_n \quad (\text{A-15})$$

Again from equation (A-8):

$$v(x, z, 0) = \sum_{n=0}^{\infty} F_n(x, 0) \sin\left\{\frac{(2n+1)\pi z}{2L_2}\right\}$$

From which:

$$F_n(x, 0) = \frac{2}{L_2} \int_0^{L_2} v(x, z, 0) \sin\left\{\frac{(2n+1)\pi z}{2L_2}\right\} dz$$

By equation (A-7a):

$$F_n(x, 0) = \frac{2}{L_2} \int_0^{L_2} \theta \sin\left\{\frac{(2n+1)\pi z}{2L_2}\right\} dz = \gamma_n \quad (\text{A-16})$$

From equation (A-9):

$$\frac{\partial v}{\partial x}(L_1, z, t) = \sum_{n=0}^{\infty} \frac{\partial F_n(L_1, t)}{\partial x} \sin\left\{\frac{(2n+1)\pi z}{2L_2}\right\}$$

Therefore, using equation (A-7d):

$$\frac{\partial F_n(L_1, t)}{\partial x} = \frac{2}{L_2} \int_0^{L_2} \bar{C} \sin\left\{\frac{(2n+1)\pi z}{2L_2}\right\} dz = \hat{C} \quad (\text{A-17})$$

Hence the initial and boundary condition for the problem defined by equation (A-14) are as follows:

$$F_n(x, 0) = \gamma_n$$

$$F_n(0, t) = \beta_n$$

$$\frac{\partial F_n(L_1, t)}{\partial x} = \hat{C}$$

The solution for $F_n(x, t)$ can be written as:

$$F_n(x, t) = u_n(x, t) + \beta_n + \frac{x^2}{2L_1} \hat{C} \quad (\text{A-18})$$

From equation (A-18):

$$F_n(0, t) = u_n(0, t) + \beta_n = \beta_n$$

Therefore,

$$u_n(0, t) = 0$$

Similarly,

$$\frac{\partial u_n}{\partial x}(L_1, t) = 0$$

Also, from equation (A-18):

$$\frac{\partial^2 F_n}{\partial x^2} = \frac{\partial^2 u_n}{\partial x^2} + \frac{\hat{C}}{L_1} \quad (\text{A-19})$$

and

$$\frac{\partial F_n}{\partial t} = \frac{\partial u_n}{\partial t} \quad (\text{A-20})$$

Substituting equations (A-18) to (A-20) in equation (A-14):

$$\frac{\partial u_n}{\partial t} = \hat{D} \left(\frac{\partial^2 u_n}{\partial x^2} - \left\{ \frac{(2n+1)\pi}{2L_2} \right\}^2 u_n \right) + \hat{t}_n(x, t) \quad (\text{A-21})$$

where,

$$\hat{t}_n(x, t) = \alpha_c - \alpha_\beta - \hat{u} \frac{x^2}{2L_1} + t_n(x, t) \quad (\text{A-22})$$

$$\alpha_c = \frac{\hat{D}\hat{C}}{L_1}$$

$$\alpha_\beta = \hat{D}\lambda_1^2\beta_n$$

$$\hat{u} = \hat{D}\lambda_1^2\hat{C}$$

$$\lambda_1 = \frac{(2n+1)\pi}{2L_2}$$

The solution for $u_n(x,t)$ in equation (A-21) can be written as:

$$u_n(x, t) = \sum_{m=0}^{\infty} C_{mn}(t) \sin\left\{\frac{(2m+1)\pi x}{2L_1}\right\} \quad (\text{A-23})$$

Here $C_{mn}(t)$ is a coefficient to be determined that depends only on time.

From equation (A-23):

$$u_n(x, 0) = \sum_{m=0}^{\infty} C_{mn}(0) \sin\left\{\frac{(2m+1)\pi x}{2L_1}\right\}$$

Therefore,

$$C_{mn}(0) = \frac{2}{L_1} \int_0^{L_1} u_n(x, 0) \sin\left\{\frac{(2m+1)\pi x}{2L_1}\right\} dx \quad (\text{A-24})$$

From equation (A-18):

$$F_n(x, 0) = \gamma_n = u_n(x, 0) + \frac{x^2 \hat{C}}{2L_1}$$

From which:

$$u_n(x, 0) = \gamma_n - \beta_n - \frac{x^2}{2L_1} \hat{C} = \tilde{\gamma} \quad (\text{A-25})$$

From equations (A-24) and (A-25):

$$C_{mn}(0) = \frac{2}{L_1} \int_0^{L_1} \tilde{\gamma} \sin\left\{\frac{(2m+1)\pi x}{2L_1}\right\} dx = \tilde{\gamma} \quad (\text{A-26})$$

Differentiating equation (A-23) twice with respect to x :

$$\frac{\partial^2 u_n}{\partial x^2} = -\sum_{m=0}^{\infty} C_{mn}(t) \left\{\frac{(2m+1)\pi}{2L_1}\right\}^2 \sin\left\{\frac{(2m+1)\pi x}{2L_1}\right\} \quad (\text{A-27})$$

Substituting equations (A-23) and (A-27) in equation (A-21):

$$\frac{dC_{mn}(t)}{dt} + \alpha C_{mn}(t) = t_m(x, t) \quad (\text{A-28})$$

where,

$$\alpha = \hat{D}(\lambda_1^2 + \lambda_2^2)$$

$$\lambda_1 = \frac{(2n+1)\pi}{2L_2}$$

$$\lambda_2 = \frac{(2m+1)\pi}{2L_1}$$

The nonhomogeneous term $t_m(x, t)$ can be written as:

$$t_m(x, t) = \frac{2}{L_1} \int_0^{L_1} \hat{t}_n(x, t) \sin\lambda_2 x dx \quad (\text{A-29})$$

The solution for the first order ordinary differential equation (A-28) can be written

as:

$$C_{mn}(t) = \frac{1}{e^{\hat{a}t}} \int_0^t e^{\hat{a}t} t_m(x, t) dt + \frac{g}{e^{\hat{a}t}} \quad (\text{A-30})$$

where,

$g = \text{a constant.}$

From equation (A-30):

$$C_{mn}(0) = g \quad (\text{A-31})$$

From equations (A-26) and (A-31):

$$g = \hat{\gamma}$$

Evaluating the integral in equation (A-17):

$$\hat{C} = \frac{4\bar{C}}{(2n+1)\pi}$$

Evaluating the integral in equation (A-15):

$$\beta_n = \frac{4}{(2n+1)\pi} \left\{ \bar{B} - \bar{A} - \frac{\bar{D}}{L_2} \left[\frac{L_2 \text{Sin}(L_2 \lambda_1)}{\lambda_1} - \frac{1}{\lambda_1^2} \right] \right\}$$

Evaluating the integral in equation (A-16):

$$\gamma_n = \frac{4}{(2n+1)\pi} \left\{ \theta_o - \bar{A} - \frac{\bar{D}}{L_2} \left[\frac{L_2 \sin(L_2 \lambda_1)}{\lambda_1} - \frac{1}{\lambda_1^2} \right] \right\}$$

Evaluating the integral in equation (A-26):

$$\hat{\gamma} = \frac{4}{(2m+1)\pi} \left\{ \gamma_n - \beta_n - \frac{\hat{C}}{L_1} \left[\frac{L_1 \sin(L_1 \lambda_2)}{\lambda_2} - \frac{1}{\lambda_2^2} \right] \right\}$$

Evaluating the integral in equation (A-13):

$$t_n(x, t) = \frac{4\alpha_D}{(2n+1)\pi} - \frac{4p}{(2n+1)\pi} = \bar{t}_n$$

where,

$$\alpha_D = \frac{\hat{D}\bar{D}}{L_2}$$

From equation (A-22):

$$\hat{t}_n(x, t) = \alpha_c - \alpha_\beta - \hat{u} \frac{x^2}{2L_1} + \bar{t}_n$$

Evaluating the integral in equation (A-29), $t_m(x,t)$ can be expressed as:

$$t_m(x, t) = \frac{4}{(2m+1)\pi} [t_{mn1} - t_{mn2}] = t_{mn} \quad (\text{A-32})$$

where,

$$t_{mn1} = \bar{t}_n + \alpha_c - \alpha_\beta$$

$$t_{mn2} = \frac{\hat{u}}{L_1} \left\{ \frac{L_1 \sin(L_1 \lambda_2)}{\lambda_2} - \frac{1}{\lambda_2^2} \right\}$$

From equations (A-30) and (A-32):

$$C_{mn}(t) = \frac{t_{mn}}{\hat{a}} + \hat{\gamma} e^{-at} \quad (\text{A-33})$$

Substituting equations (A-8), (A-18), (A-23), and (A-33) in equation (A-2), the moisture content $\theta(x, z, t)$ can be expressed as:

$$\theta = \sum_{n=0}^{\infty} \left[\sum_{m=0}^{\infty} \left\{ \frac{t_{mn}}{\hat{a}} + \hat{\gamma} e^{-at} \right\} \sin \lambda_2 x + \beta_n + \frac{x^2}{2L_1} \hat{C} \right] \sin \lambda_1 z + \bar{A} + \frac{z^2}{2L_2} \bar{D} \quad (\text{A-34})$$

Differentiating equation (A-34), the fluxes in x- and z-directions are obtained as:

$$\frac{\partial \theta}{\partial x} = \sum_{n=0}^{\infty} \left\{ \sum_{m=0}^{\infty} \left[\frac{t_{mn}}{\hat{a}} + \hat{\gamma} e^{-at} \right] \lambda_2 \cos(\lambda_2 x) + \frac{\hat{C}x}{L_1} \right\} \sin(\lambda_1 z)$$

$$\frac{\partial \theta}{\partial z} = \sum_{n=0}^{\infty} \left\{ \sum_{m=0}^{\infty} \left[\frac{t_{mn}}{\alpha} + \gamma e^{-\alpha t} \right] \sin(\lambda_2 x) + \beta_n + \frac{x^2}{2L_1} \hat{C} \right\} \lambda_1 \cos(\lambda_1 z) + \frac{z\bar{D}}{L_2}$$

APPENDIX B

DERIVATION OF ANALYTICAL SOLUTION FOR THE SATURATED LEACHATE FLOW EQUATION

The simplified form of the leachate flow equation in the unsaturated zone is expressed as:

$$\frac{\partial h}{\partial t} = \alpha \frac{\partial^2 h}{\partial x^2} + \beta h - \bar{p} \quad (\text{B-1})$$

where,

$$\alpha = \frac{T}{n_e}$$

$$\beta = -\frac{K}{n_e d}$$

$$\bar{p} = \frac{-N_l + K_c}{n_e}$$

The initial and boundary conditions for the saturated flow problem described by equation (B-1) are as follows:

$$h(x, 0) = h_1 \quad (\text{B-2})$$

$$h(L, t) = h_L \quad (\text{B-3})$$

$$\frac{\partial h}{\partial x}(0, t) = i_o \quad (\text{B-4})$$

where,

L = length of the saturated leachate mound from the origin at $x = 0$.

The solution for $h(x,t)$ described by equation (B-1) can be written as:

$$h(x, t) = u(x, t) + i_o \left(1 - \frac{x}{L}\right) + \frac{x^2}{L^2} h_L \quad (\text{B-5})$$

Here $u(x,t)$ is the correction as a function of both space and time which is to be determined.

Using the initial and boundary conditions described by equations (B-2) to (B-4) in the expression (B-5), the initial and the homogeneous boundary conditions for the dependent variable $u(x,t)$ is obtained as follows:

$$u(x, 0) = u_1(x) \quad (\text{B-6})$$

$$\frac{\partial u}{\partial x}(0, t) = 0 \quad (\text{B-7})$$

$$u(L, t) = 0 \quad (\text{B-8})$$

where,

$$u_1(x) = h_1 - i_o x + \frac{i_o x^2}{L} - \frac{x^2}{L^2} h_L$$

Using the expression (B-5) in equation (B-1), the governing equation for the dependent variable $u(x,t)$ can be obtained as:

$$\frac{\partial u}{\partial t} = \alpha \frac{\partial^2 u}{\partial x^2} + \beta u + q \quad (\text{B-9})$$

where,

$$q = -\frac{2\alpha i_o}{L} + \frac{2\alpha h_L}{L^2} - \bar{p} + \beta i_o x \left(1 - \frac{x}{L}\right) + \frac{\beta x^2}{L^2} h_L \quad (\text{B-10})$$

The solution for $u(x,t)$ in equation (B-9) can be written as:

$$u(x,t) = \sum_{n=0}^{\infty} C_n(t) \cos\left\{\frac{(2n+1)\pi x}{2L}\right\} \quad (\text{B-11})$$

The cosine series expression for the nonhomogeneous term q can be written as:

$$q = \sum_{n=0}^{\infty} \phi_n(x,t) \cos\left\{\frac{(2n+1)\pi x}{2L}\right\} dx \quad (\text{B-12})$$

Therefore,

$$\phi_n(x, t) = \frac{2}{L} \int_0^L q \cos\left\{\frac{(2n+1)\pi x}{2L}\right\} dx = \hat{\phi}_n \quad (\text{B-13})$$

Using equation (B-11) and (B-12) in equation (B-9), the following ordinary differential equation for $C_n(t)$ is obtained as:

$$\frac{dC_n(t)}{dt} + \hat{\nu}C_n(t) = \hat{\phi}_n(x, t) \quad (\text{B-14})$$

where,

$$\hat{\nu} = \hat{\lambda} - \beta$$

$$\hat{\lambda} = \alpha \left\{ \frac{(2n+1)\pi}{2L} \right\}^2$$

The solution for the first order ordinary differential equation (B-14) is expressed as:

$$C_n(t) = \frac{1}{e^{\hat{\nu}t}} \int_0^t e^{\hat{\nu}t} \hat{\phi}_n dt + \frac{g_s}{e^{\hat{\nu}t}} \quad (\text{B-15})$$

where,

$$g_s = \text{a constant.}$$

From equation (B-15):

$$C_n(0) = g_s$$

From equation (B-11), $C_n(t)$ can be expressed as:

$$C_n(t) = \frac{2}{L} \int_0^L u(x, t) \cos\left\{\frac{(2n+1)\pi x}{2L}\right\} dx$$

Therefore,

$$C_n(0) = \frac{2}{L} \int_0^L u(x, 0) \cos\left\{\frac{(2n+1)\pi x}{2L}\right\} dx \quad (\text{B-16})$$

Substituting equation (B-6) in equation (B-16) and evaluating the integral:

$$C_n(0) = g_s = \frac{2}{L} \left[h_1 u_0 - i_0 u_1 + \frac{i_0}{L} u_2 - \frac{h_L}{L^2} u_2 \right] \quad (\text{B-17})$$

where,

$$u_0 = \frac{2L}{(2n+1)\pi} (-1)^n$$

$$u_1 = \frac{2L^2}{(2n+1)\pi} \left[(-1)^n - \frac{2}{(2n+1)\pi} \right]$$

$$u_2 = \frac{2L}{(2n+1)\pi} \left[(-1)^n L^2 - \frac{8(-1)^n}{\pi^2} \left(\frac{L}{2n+1} \right)^2 \right]$$

Substituting equation (B-10) in equation (B-13) and evaluating the integral, the following expression is obtained as:

$$\phi_n = \frac{4\alpha}{L^2} u_0 \left[-i_0 + \frac{h_L}{L} \right] + \frac{2}{L} \left[-\bar{p} u_0 + \beta i_0 u_1 - \beta i_0 \frac{u_2}{L} + \beta h_L \frac{u_2}{L^2} \right]$$

Finally, the expression for the leachate mound head $h(x,t)$ can be written from equations (B-5) and (B-11) as:

$$h(x,t) = \sum_{n=0}^{\infty} C_n(t) \cos\left\{ \frac{(2n+1)\pi x}{2L} \right\} + i_0 x \left(1 - \frac{x}{L} \right) + \frac{x^2}{L^2} h_L$$

where,

$$C_n(t) = \frac{\phi_n}{\hat{v}} + g_s e^{-\hat{v}t}$$

APPENDIX C

FORTRAN PROGRAM FOR FILL MODEL

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C
C ESTIMATION OF UNSTEADY STATE LEACHATE FLOW THROUGH LANDFILLS FOR
C UNSATURATED-SATURATED CONDITON IN A TWO-DIMENSIONAL DOMAIN
C
C MAIN PROGRAM
C
  IMPLICIT REAL *8(A-H,O-Z)
  COMMON /STOR/ BCONL,BCONR
  COMMON /MELT/ TOTMET,SNOWF
  COMMON /RUN/ HF,PINFIL,T,RUNOF,TOTRUN,THETAP,TAVINF,TAVRUN,TRR,HR
  COMMON /GAUMAIN/ ALBEDO,PSYCRO,AS,WP,FC,BC,SNOWP,ACEVAT,
  +DIFFUS,HCON
  COMMON /MFLO/ AQV,FLOWL,FLOWR,FLOWVV,FLOWLV,FLOWLI,BEFFICY,S
  +MEFICY,EFFICY
  COMMON /MONFO/ GRAD,BIG
  COMMON /SILA/ FLOMLL,FLOMLR,FLOWMV,BIGLAT,SMALAT,AVMOST,
  +RAINM,PERCM,RUNM,EVAM,RECHM,BIGVER,SMAVER
  COMMON /BIND/ FLOML,TOTML,TOMFLO,BIGEST,SMALST,TOTFL,TOTFV,
  +AVDALF,AVMOLF,AVMOVF,AVDAVF,AVDL,AVML,DAIFTO,TOTLL,TOTLR,
  +AVDLL,AVDLR,AVDA+LL,AVDALR,AVMLL,AVMLR,ANRAIN,
  +ANPERC,ANRUN,ANEVAT,AVDAIN,AVDAIC,AVD
  +AUN,AVDAET,AVEFCY,ANVPE,ANVERT
  PARAMETER(NOD=100,NOC=50,NEX=50,NOR=20)
  REAL *8 FLOWLV(1),THETA(NOD),THETAP(NOD),TEMP(2,365),PRECIP(2,365
  +),DIFFUS(NOD,4),HCON(NOD,4),SLOPE(NEX),SOLRAD(2,365),WIND(365),THI
  +CK(NOC),SLOPBR(NOC)
  REAL *8 FLOWL(1),FLOWR(1),FLOML(1,2,12),TOTML(2,12),TOMFLO(2,12),T
  +OTFL(1),TOTFV(1),AVDALF(1),AVMOLF(1),AVMOVF(1),AVDAVF(1),FLOMLL(1,
  +2,12),RAIN(2,365),RAINM(2,12),PERCM(2,12),RUNM(2,12),EVAM(2,12)
  REAL *8 HIN(NOC),XLAI(365),WCF(365),PINFIL(NEX),RUNOF(NEX),HR(NEX)
  +,BC(NEX),GRAD(1,NOC),CBIG(1),CAPINF(2,12),FLORUN(2,12),HDD(NOC)
  REAL *8 HDIN(NOC),FLOMLR(1,2,12),FLOWMV(1,2,12),BIGLAT(1),SMALAT(1
  +),HDP(NOC),TOTLL(1),TOTLR(1),AVDALL(1),AVDALR(1),AVMOST(1,2,12),FL
  +OWLI(1),BEFFICY(2),SMEFICY(2)
  REAL *8 THETAM(NOD),HORAIN(24),HF(NOC,24),TOTRUN(NOC,24),RECHM(1,2
  +,12),BIGVER(1),SMAVER(1),BIG(1),BBIG(1,12),COFINT(NOD,4),VDISTB(NO
  +D,1),HCOL(1,NEX)
  INTEGER INTNO(NOD),LINKNT(NOD),LINKNB(NOD),LINKNL(NOD),LINKNR(NOD)
  +,NEXUP(NEX),ICONF(NOD),NODE(NOD),IGROW(365),NRCONF(NOD),NCCONF(NOD
  +),ICRNO(NOR,NOC),IBRCON(NOD),IBET(NOD),INEAR(NOD)
  INTEGER LCONL(NOC),ICOL(NOC),LINKNU(NEX),LINKND(NEX),LCONR(NOC),MO
  +ND(NOD)
  CHARACTER*3 MONTH(12)
  EXTERNAL IN3PUT,SNOW,CONFIG,RUNOFF,SELECT,FDMITER,SATELIM,IN4PUT,F

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```

+LOW,IN2PUT
OPEN(2,FILE='LAND1.DAT',STATUS='OLD')
OPEN(3,FILE='RESULT.OUT',STATUS='UNKNOWN')
OPEN(5,FILE='LAND2.DAT',STATUS='OLD')
OPEN(6,FILE='FILLINP',STATUS='OLD')
OPEN(7,FILE='LAND3.DAT',STATUS='OLD')
CALL IN3PUT(NINTER,NEXTUP,NEXBOT,NOROW,NOCOL,NOBAR,DT,
+TDAYS,NYEARS,NMONTH,DTR,MONTH,AS,ALBEDO,PSYCRO,SNOWP,WP,
+ROUGH,FORCOV,TEMPD,EXPOC,SLOPMC,COFM,HC,FC,THETAS,
+HCONS,CONDB,THETIN,THETRR,DX,DZ,SAIE,BCOND,EFPORO,HCONCOV,DRC)
NONODE=NINTER+NEXTUP
NODTOT=NONODE+NEXBOT
CALLIN4PUT(ICONF,NRCONF,NCCONF,IBRCON,INEAR,LINKNT,LINKNB,LINKNR,
+LINKNL,SLOPBR,HDIN,NOCOL,ICOL,NEXTUP,LINKND,LINKNU,SLOPE,TDAYS,DT,
+NYEARS,TEMP,SOLRAD,PRECIP,RAIN,XLAI,WCF,IGROW,WIND,THICK,NODTOT,LC
+ONL,LCONR,COFINT)
TRR=0.0
DO 670 I=1,NOBAR
DO 670 J=1,NYEARS
DO 670 K=1,NMONTH
FLOMLL(I,J,K)=0.0
FLOMLR(I,J,K)=0.0
FLOWMV(I,J,K)=0.0
AVMOST(I,J,K)=0.0
RECHM(I,J,K)=0.0
670 CONTINUE
DO 675 I=1,NYEARS
DO 675 J=1,NMONTH
TOTML(I,J)=0.0
RAINM(I,J)=0.0
PERCM(I,J)=0.0
RUNM(I,J)=0.0
EVAM(I,J)=0.0
CAPINF(I,J)=0.0
FLORUN(I,J)=0.0
675 CONTINUE
DO 680 I=1,NOBAR
TOTFL(I)=0.0
TOTLL(I)=0.0
TOTLR(I)=0.0
TOTFV(I)=0.0
680 CONTINUE
BIGEST=0.0
SMALST=0.0
DTTIME=TDAYS-DT
CALL CONFIG(ICONF,NODTOT,IBRCON,NOROW,NOCOL,LINKNT,LINKNB,
+NOBAR,NONODE,NRCONF,NCCONF,NODE,NEXUP,INTNO,IBET,ICRNO)
EDIFF=0.0
DO 413 J=1,NODTOT
DO 413 I=1,NOBAR

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VDISTB(J,I)=0.0
413 CONTINUE
DO 410 I=1,NOROW
DO 410 J=1,NOCOL
NN=ICRNO(I,J)
IF(NN.EQ.0)GOTO 410
DO 411 K=1,NOBAR
KK=2*K+1
IF((IBET(NN).EQ.KK).OR.(IBET(NN).EQ.K))IB=K
411 CONTINUE
NNB=LINKNB(NN)
IF(ICONF(NN).EQ.4)THEN
VDISTB(NN,IB)=0.0
ELSE
IF(ICONF(NNB).EQ.4)THEN
VDISTB(NN,IB)=COFINT(NN,3)*DZ
ELSE
VDISTB(NN,IB)=COFINT(NN,3)*DZ+VDISTB(NNB,IB)
ENDIF
ENDIF
410 CONTINUE
DO 412 I=1,NEXTUP
NN=NEXUP(I)
HCOL(IB,I)=VDISTB(NN,IB)
412 CONTINUE
II=0
16 II=II+1
T=0.0
15 T=T+DT
write(*,*)'t=',t
J1=T/DT
DO 415 I=1,NOROW
DO 415 J=1,NOCOL
NN=ICRNO(I,J)
IF(NN.EQ.0)GOTO 415
DO 420 K=1,NOBAR
KK=2*K+1
IF((IBET(NN).EQ.KK).OR.(IBET(NN).EQ.K))IB=K
420 CONTINUE
IF(J1.GT.1)THEN
IF(HDP(J).GE.VDISTB(NN,IB))THEN
MOND(NN)=1
ELSE
MOND(NN)=0
ENDIF
ELSE
IF(HDIN(J).GE.VDISTB(NN,IB))THEN
MOND(NN)=1
ELSE
MOND(NN)=0

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ENDIF
ENDIF
415 CONTINUE
CALL IN2PUT(HORAIN)
IF(J1.EQ.254)THEN
DARAN=0.0
DO 29 J=1,24
DARAN=DARAN+HORAIN(J)
29 CONTINUE
ELSE
ENDIF
DO 26 I=1,NEXTUP
HIN(I)=0.0
26 CONTINUE
95 IF(TEMP(II,J1).LE.32.0)THEN
SNOWF=SNOWP+RAIN(II,J1)
TOTMET=0.0
SNOWP=SNOWF
ELSE
IF(SNOWP.GT.0.0)THEN
CALL SNOW(RAIN(II,J1),TEMP(II,J1),SOLRAD(II,J1),WIND(J1),FORCOV,EX
+POC,ALBEDO,SNOWP,COFM,TEMPD)
ELSE
TOTMET=0.0
SNOWF=0.0
SNOWP=SNOWF
ENDIF
ENDIF
IF((TEMP(II,J1).LE.32.0).OR.((PRECIP(II,J1)+TOTMET).EQ.0.0))THEN
DO 21 I=1,NEXTUP
HR(I)=0.0
DO 687 K=1,24
HF(I,K)=0.0
TOTRUN(I,K)=0.0
687 CONTINUE
RUNOF(I)=0.0
PINFIL(I)=0.0
TAVINF=0.0
TAVRUN=0.0
FLONET=0.0
21 CONTINUE
ELSE
DO 31 I=1,NONODE
IF((II.EQ.1).AND.(J1.EQ.1))THEN
THETAP(I)=THETIN
ELSE
ENDIF
31 CONTINUE
CALL RUNOFF(TOTMET,PRECIP,HIN,SLOPE,ROUGH,DTR,THETAS,HCONS,HORAIN,
+HC,FC,II,J1,NEXTUP,NEXUP,LINKNB,ICONF,LCONL,LCONR,NOCOL,FLONET,DX,

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+TEMP,NODE,NONODE,BCOND,COFINT,HCONCOV)
ENDIF
CALL FDMITER(THETIN,SLOPMC,HCONS,SAIE,THETAS,LINKNT,LINKNR,LINKNB,
+LINKNL,ICONF,THETA F,DT,DX,DZ,J1,NONODE,NEXTUP,NEXUP,II,NODE,PRECIP
+(II,J1),HR,TEMP(II,J1),RAIN(II,J1),SOLRAD(II,J1),IGROW(J1),XLAI(J1
+),WCF(J1),TAVRUN,THETAP,COFINT,MOND)
AVIN=RAIN(II,J1)-TAVRUN-ACEVAT
CALL SATELIM(NODTOT,NOBAR,NOROW,NOCOL,ICRNO,ICONF,IBET,
+LINKNB,NEXTUP,NEXUP,THETAS,BC,HDIN,ICOL,LCONL,LCONR,HCONS,
+CONDB,EFPORO,BCONL,BCONR,DT,DX,DZ,THICK,TAVRUN,ACEVAT,
+PRECIP(II,J1),NODE,NONODE,SLOPBR,THETA F,II,J1,GRAD,THETAM,
+MOND,LINKNT,HDD,HDP,TAVSOR,HCOL,SLOPMC,
+SAIE,THETAP,VDISTB,HCON,DIFFUS,FC,LINKNR,LINKNL,DRC)
DO 139 I=1,NOBAR
IF(J1.EQ.1)CBIG(I)=0.0
139 CONTINUE
CALL FLOW(NOCOL,NOBAR,HDD,SLOPBR,HCONS,J1,II,RAIN(II,J1),THICK,CON
+DB)
AVTHET=0.0
DO 902 I=1,NONODE
NN=NODE(I)
AVTHET=AVTHET+THETA F(I)
902 CONTINUE
DO 903 I=1,NEXBOT
AVTHET=AVTHET+THETAS
903 CONTINUE
AVTHET=AVTHET/NODTOT
IF((J1.EQ.1).OR.(J1.EQ.182).OR.(J1.EQ.365))THEN
WRITE(3,1000)J1
1000 FORMAT(5X,'TIME T = ',I3,1X,'DAY(S)')
WRITE(3,750)
750 FORMAT(5X,'TABLE : DEPTH OF LEACHATE MOUND',//,5X,'_____
+_____','/)
WRITE(3,751)
751 FORMAT(5X,'BARRIER',7X,'COLUMN NO.',5X,'DEPTH OF MOUND')
WRITE(3,752)
752 FORMAT(38X,'(FT)',//,5X,'_____',7X,'_____',5X,'_____
+_' '/')
DO 93 J=1,NOCOL
WRITE(3,94)J,HDD(J)
94 FORMAT(13X,I3,8X,F10.2)
93 CONTINUE
ELSE
ENDIF
CALL SELECT(NO BAR,FLOWL,FLOWR,AQV,J1,II,AVTHET,RAIN,AVIN,TAVRUN,AC
+EVAT,TAVSOR,BIG,BBIG,TAVINF,FLONET,CAPINF,FLORUN)
IF(T.LE.DTTIME)GOTO 15
JJ=J1
IF(II.LT.NYEARS)GOTO 16
CALL FIND(NYEARS,NMONTH,NO BAR,FLOMLL,FLOMLR,FLOWMV,BIGLAT,

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+SMALAT,TDAYS,RAINM,RUNM,EVAM,RECHM,FLORUN,CAPINF,ANFLOW,BALFLO)
WRITE(3,791)
791 FORMAT(1X,'TABLE : MONTHLY VALUES OF PRECIPITATION, RUNOFF, INPUT
+,    ',9X,'AND EVAPOTRANSPIRATION IN INCHES',/,1X,'_____
+_____ ',/)
DO 756 I=1,NYEARS
WRITE(3,761)(MONTH(K),K=1,6)
761 FORMAT(17X,2(4X,A3),4(5X,A3),/)
WRITE(3,762)(RAINM(I,K),K=1,6)
762 FORMAT(1X,'PRECIPITATION',2X,6(2X,F6.2),/)
WRITE(3,811)(RUNM(I,K),K=1,6)
811 FORMAT(1X,'RUNOFF',9X,6(2X,F6.2),/)
WRITE(3,812)(PERCM(I,K),K=1,6)
812 FORMAT(1X,'NET INFILT. ',3X,6(2X,F6.2),/)
WRITE(3,813)(EVAM(I,K),K=1,6)
813 FORMAT(1X,'EVAPOTRANSP.',3X,6(2X,F6.2),/)
WRITE(3,763)(MONTH(K),K=7,12)
763 FORMAT(17X,2(4X,A3),4(5X,A3),/)
WRITE(3,792)(RAINM(I,K),K=7,12)
792 FORMAT(1X,'PRECIPITATION',2X,6(2X,F6.2),/)
WRITE(3,816)(RUNM(I,K),K=7,12)
816 FORMAT(1X,'RUNOFF',9X,6(2X,F6.2),/)
WRITE(3,817)(PERCM(I,K),K=7,12)
817 FORMAT(1X,'NET INFILT. ',3X,6(2X,F6.2),/)
WRITE(3,818)(EVAM(I,K),K=7,12)
818 FORMAT(1X,'EVAPOTRANSP.',3X,6(2X,F6.2),/)
756 CONTINUE
WRITE(3,757)
757 FORMAT(1X,'TABLE : MONTHLY VALUES OF AVERAGE MOISTURE CONTENTS
+IN THE UNSATURATED ZONE',/,1X,'_____
+_____ ',/)
DO 759 I=1,NOBAR
WRITE(3,758)I
758 FORMAT(10X,'AVERAGE MOISTURE CONTENTS UPON BARRIER NO. ',I2,/,10X,
+'_____ ',/)
DO 759 J=1,NYEARS
WRITE(3,801)J
801 FORMAT(///,1X,'MONTHLY VALUES FOR THE YEAR ',I2,':',/)
WRITE(3,773)(MONTH(K),K=1,6)
773 FORMAT(1X,2(7X,A3),4(8X,A3),/)
WRITE(3,774)(AVMOST(I,J,K),K=1,6)
774 FORMAT(1X,6(3X,F8.3),/)
WRITE(3,776)(MONTH(K),K=7,12)
776 FORMAT(1X,2(7X,A3),4(8X,A3),/)
WRITE(3,777)(AVMOST(I,J,K),K=7,12)
777 FORMAT(1X,6(3X,F8.3))
759 CONTINUE
WRITE(3,781)
781 FORMAT(///,1X,'TABLE : MONTHLY VALUES OF RECHARGE ONTO THE LEACH
+ATE MOUND',/,9X,'IN INCHES',/,1X,'_____

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```

+ _____',//)
DO 784 I=1,NOBAR
WRITE(3,783)I
783 FORMAT(9X,'RECHARGE VALUES ONTO THE MOUND UPON BARRIER NO. ',I2,/
+9X,' _____')
DO 784 J=1,NYEARS
WRITE(3,721)J
721 FORMAT(///,1X,'MONTHLY VALUES FOR THE YEAR ',I2,' : ',//)
WRITE(3,722)(MONTH(K),K=1,6)
722 FORMAT(1X,2(7X,A3),4(8X,A3),/)
WRITE(3,723)(RECHM(I,J,K),K=1,6)
723 FORMAT(1X,6(3X,F8.2),/)
WRITE(3,724)(MONTH(K),K=7,12)
724 FORMAT(1X,2(7X,A3),4(8X,A3),/)
WRITE(3,726)(RECHM(I,J,K),K=7,12)
726 FORMAT(1X,6(3X,F8.2))
784 CONTINUE
WRITE(3,790)
790 FORMAT(////,1X,'TABLE : MONTHLY VALUES OF LATERAL AND VERTICAL F
+LOW',/)
WRITE(3,795)
795 FORMAT(1X,' _____
+ _____',/)
DO 755 I=1,NOBAR
WRITE(3,760)I
760 FORMAT(21X,'FLOW DUE TO BARRIER NO. ',I2)
WRITE(3,800)
800 FORMAT(21X,' _____',/)
DO 755 J=1,NYEARS
WRITE(3,765)J
765 FORMAT(///,1X,'MONTHLY VALUES FOR THE YEAR ',I2,' : ',//)
WRITE(3,771)(MONTH(K),K=1,6)
771 FORMAT(18X,2(4X,A3),4(5X,A3),/)
WRITE(3,770)(FLOMLL(I,J,K),K=1,6)
770 FORMAT(1X,'LATERAL FLOW',/,1X,'(LEFT DIRECTION)',1X,6(2X,F6.2),/)
WRITE(3,772)(FLOMLR(I,J,K),K=1,6),(FLOML(I,J,K),K=1,6),(FLOWMV(I,J
+,K),K=1,6)
772 FORMAT(1X,'LATERAL FLOW',/,1X,'(RIGHT DIRECTION)',6(2X,F6.2),/,1X
+,'TOTAL MONTHLY',/,1X,'LATERAL FLOW',5X,6(2X,F6.2),/,1X,'VERTICAL
+ FLOW',4X,6(2X,F6.4),//)
WRITE(3,775)(MONTH(K),K=7,12)
775 FORMAT(18X,2(4X,A3),4(5X,A3),/)
WRITE(3,780)(FLOMLL(I,J,K),K=7,12)
780 FORMAT(1X,'LATERAL FLOW',/,1X,'(LEFT DIRECTION)',1X,6(2X,F6.2),/)
WRITE(3,785)(FLOMLR(I,J,K),K=7,12),(FLOML(I,J,K),K=7,12),(FLOWMV(I
+,J,K),K=7,12)
785 FORMAT(1X,'LATERAL FLOW',/,1X,'(RIGHT DIRECTION)',6(2X,F6.2),/,1X
+,'TOTAL MONTHLY',/,1X,'LATERAL FLOW',1X,6(2X,F6.2),/,1X,'VERTICAL
+ FLOW',4X,6(2X,F6.4))
755 CONTINUE

```

```

WRITE(3,805)
805 FORMAT(/,1X,'
+ _____',1X,'LATERAL FLOW IN CU.FT. PER UNIT WIDTH OF LANDF
+ILL',1X,'VERTICAL FLOW IN CU.FT. PER UNIT AREA OF LANDFILL',///
+/)
WRITE(3,810)
810 FORMAT(1X,'TABLE : TOTAL MONTHLY LATERAL FLOW FROM
+THE LANDFILL',/)
WRITE(3,815)
815 FORMAT(1X,'
+ _____')
DO 685 I=1,NYEARS
WRITE(3,635)I
635 FORMAT(/,28X,'YEAR ',I3)
WRITE(3,820)
820 FORMAT(28X,' _____',//)
WRITE(3,655)(MONTH(J),J=1,6)
655 FORMAT(3X,6(7X,A3),/)
WRITE(3,690)(TOTML(I,J),J=1,6)
690 FORMAT(1X,6(2X,F8.4),//)
WRITE(3,700)(MONTH(J),J=7,12)
700 FORMAT(3X,6(7X,A3),/)
WRITE(3,695)(TOTML(I,J),J=7,12)
695 FORMAT(1X,6(2X,F8.4),//)
685 CONTINUE
WRITE(3,825)
825 FORMAT(1X,'
+ _____',1X,'LATERAL FLOW IN CU.FT. PER UNIT WIDTH OF LANDFILL',/
+///)
WRITE(3,830)
830 FORMAT(1X,'TABLE : AVERAGE DAILY AND MONTHLY FLOWS UPON
+EACH BARRIER',//,1X,'
+ _____',/)
DO 720 I=1,NOBAR
WRITE(3,715)I
715 FORMAT(1X,'FLOW DUE TO BARRIER ',I2,/,1X,' _____',
+/)
WRITE(3,725)AVDALL(I),AVDALR(I),AVDALF(I),AVMOLF(I),AVDAVF(I),AVMO
+VF(I)
725 FORMAT(1X,'AVERAGE DAILY LATERAL FLOW (LEFT DIRECTION)',1X,F10.4,/
+/,1X,'AVERAGE DAILY LATERAL FLOW (RIGHT DIRECTION)',5X,F10.4,/,1X
+,'AVERAGE DAILY LATERAL FLOW',10X,F10.4,/,5X,'AVERAGE MONTHLY LAT
+ERAL FLOW',10X,F10.4,/,1X,'AVERAGE DAILY VERTICAL FLOW',10X,F10.
+4,/,1X,'AVERAGE MONTHLY VERTICAL FLOW',10X,F10.4,/,1X,' _____
+ _____',///)
720 CONTINUE
WRITE(3,730)
730 FORMAT(1X,'TABLE : AVERAGE, PEAK, AND MINIMUM VALUES OF FLOW',/1
+4X,'UPON ALL OF THE BARRIERS',//,1X,'
+ _____',/)

```

```

WRITE(3,650)BIGEST,SMALST,AVDLL,AVDLR,AVDL
650 FORMAT(1X,'PEAK DAILY LATERAL FLOW ',6X,F10.4,//,1X,'MINIMUM DAIL
+Y LATERAL FLOW ',4X,F10.4,//,1X,'AVERAGE DAILY LATERAL FLOW (LEFT
+DIRECTION)',3X,F10.4,//,1X,'AVERAGE DAILY LATERAL FLOW (RIGHT DIRE
+CTION)',3X,F10.4,//,1X,'AVERAGE DAILY LATERAL FLOW',10X,F10.4,/)
WRITE(3,651)AVMLL,AVMLR,AVML
651 FORMAT(1X,'AVERAGE MONTHLY LATERAL FLOW (LEFT DIRECTION)',3X,F10.4
+,,//,1X,'AVERAGE MONTHLY LATERAL FLOW (RIGHT DIRECTION)',3X,F10.4/
+/,1X,'AVERAGE MONTHLY LATERAL FLOW ',1X,F10.4,//,1X,'_____
+_____','//)
113 STOP
END

```

C

C THIS SUBROUTINE READS DATA FROM THE DATA FILE

C

```

SUBROUTINE IN3PUT(NINTER,NEXTUP,NEXBOT,NOROW,NOCOL,NOBAR,DT,
+TDAYS,NYEARS,NMONTH,DTR,MONTH,AS,ALBEDO,PSYCRO,SNOWP,WP,
+ROUGH,FORCOV,TEMPD,EXPOC,SLOPMC,COFM,HC,FC,THETAS,HCONS,
+CONDB,THETIN,THETRR,DX,DZ,SAIE,BCOND,EFPORO,HCONCOV,DRC)
IMPLICIT REAL *8(A-H,O-Z)
COMMON /STOR/ BCONL,BCONR
CHARACTER*3 MONTH(12)
READ(2,*)NINTER,NEXTUP,NEXBOT,NOROW,NOCOL,NOBAR
READ(2,*)DT,TDAYS,NYEARS,NMONTH,DTR,DX,DZ
READ(2,212)(MONTH(I),I=1,NMONTH)
212 FORMAT(12A3)
READ(2,*)AS,ALBEDO,PSYCRO,SNOWP,WP,EFPORO
READ(2,*)ROUGH,FORCOV,TEMPD,EXPOC,SLOPMC,COFM,HC
185 FORMAT(6F10.3)
READ(2,*)BCONL,BCONR,FC,THETAS,HCONS,HCONCOV,CONDB,THETIN
READ(2,*)THETRR,SAIE,BCOND,DRC
RETURN
END

```

C

C THIS SUBROUTINE READS DATA FROM THE DATA FILE

C

```

SUBROUTINE IN4PUT(ICONF,NRCONF,NCCONF,IBRCON,INEAR,LINKNT,LINKNB,L
+INKNR,LINKNL,SLOPBR,HDIN,NOCOL,ICOL,NEXTUP,LINKND,LINKNU,SLOPE,TDA
+YS,DT,NYEARS,TEMP,SOLRAD,PRECIP,RAIN,XLAI,WCF,IGROW,WIND,THICK,NOD
+TOT,LCONL,LCONR,COFINT)
IMPLICIT REAL *8(A-H,O-Z)

```


C
 C THIS SUBROUTINE READS HOURLY RAINFALL DATA
 C

```

SUBROUTINE IN2PUT(HORAIN)
IMPLICIT REAL *8(A-H,O-Z)
REAL *8 HORAIN(24)
READ(5,*)(HORAIN(J),J=1,24)
RETURN
END

```

C
 C THIS SUBROUTINE GIVES CONFIGURATION TO IDENTIFY
 C INTERNAL AND EXTERNAL NODES
 C

```

SUBROUTINE CONFIG(ICONF,NODTOT,IBRCON,NOROW,NOCOL,LINKNT,LINKNB,
+NOBAR,NONODE,NRCONF,NCCONF,NODE,NEXUP,INTNO,IBET,ICRNO)
IMPLICIT REAL *8(A-H,O-Z)
PARAMETER (NOD=100,NEX=50,NI=80,NO=100,NOR=20,NOC=50)
INTEGER ICONF(NOD),NEXUP(NEX),INTNO(NI),NODE(NO),NRCONF(NOD),NCCON
+F(NOD),ICRNO(NOR,NOC),IBRCON(NOD),IBET(NOD),LINKNT(NOD),LINKNB(NOD
+)
I1=1
I2=1
DO 240 I=1,NODTOT
IF((ICONF(I).EQ.1).OR.((ICONF(I).EQ.2).OR.(ICONF(I).EQ.3)))THEN
NEXUP(I1)=I
I1=I1+1
ELSE
IF(ICONF(I).EQ.5)THEN
INTNO(I2)=I
I2=I2+1
ELSE
GOTO 240
ENDIF
ENDIF
240 CONTINUE
I3=1
DO 250 I=1,NODTOT
IF((ICONF(I).EQ.1).OR.(ICONF(I).EQ.2).OR.(ICONF(I).EQ.3).OR.(ICONF
+(I).EQ.5))THEN
NODE(I3)=I
I3=I3+1
ELSE
GOTO 250
ENDIF

```

```
250 CONTINUE
    DO 275 I=1,NOROW
    DO 275 J=1,NOCOL
    ICRNO(I,J)=0
275 CONTINUE
    DO 260 I=1,NODTOT
    DO 265 J=1,NOROW
    IF(NRCONF(I).EQ.J)IR=J
265 CONTINUE
    DO 270 K=1,NOCOL
    IF(NCCONF(I).EQ.K)IC=K
270 CONTINUE
    ICRNO(IR,IC)=I
260 CONTINUE
    DO 280 I=1,NOROW
    DO 280 J=1,NOCOL
    NN=ICRNO(I,J)
    IF(NN.EQ.0)GOTO 280
    IF(IBRCON(NN).EQ.0)THEN
    NNT=NN
300 NNT=LINKNT(NNT)
    IF(NNT.EQ.0)THEN
    NNB=NN
295 NNB=LINKNB(NNB)
    IF(NNB.EQ.0)GOTO 280
    DO 285 K=1,NOBAR
    IF((IBRCON(NNB).EQ.K).OR.(ICONF(NNB).EQ.4))THEN
    IBET(NN)=2*K+1
    ELSE
    IBET(NN)=0
    ENDIF
285 CONTINUE
    IF(IBET(NN).EQ.0)GOTO 295
    ELSE
    DO 290 K=1,NOBAR
    IF(IBRCON(NNT).EQ.K)THEN
    IBET(NN)=2*K-1
    ELSE
    IBET(NN)=0
    ENDIF
290 CONTINUE
    IF(IBET(NN).EQ.0)GOTO 300
    ENDIF
    ELSE
    IBET(NN)=IBRCON(NN)
    ENDIF
280 CONTINUE
    RETURN
    END
```

C
 C THIS SUBROUTINE COMPUTES TOTAL DAILY SNOWMELT
 C

```

SUBROUTINE SNOW(RAINN,TEMPP,SOLRDD,WINDD,FORCOV,EXPOC,ALBEDO,
+SNOWP,COFM,TEMPD)
IMPLICIT REAL *8(A-H,O-Z)
COMMON /MELT/ TOTMET,SNOWF
IF(RAINN.GT.0.0)THEN
TEMA=TEMPP-32.0
RAINMT=0.029+0.0084*EXPOC*WINDD+0.007*RAINN
TOTMET=0.09+RAINMT*TEMA
IF(TOTMET.GT.SNOWP)TOTMET=SNOWP
IF(TOTMET.LT.0.0)TOTMET=0.0
ELSE
SOLMET=0.004*SOLRDD*COFM*(1.0-FORCOV)*(1.0-ALBEDO)
TEMT=0.22*(TEMPP-32.0)
TEMDM=0.78*(TEMPD-32.0)
FMT=0.029*FORCOV*(TEMPP-32.0)
TEMWMT=0.0084*EXPOC*WINDD*(TEMT+TEMDM)
TOTMET=SOLMET+FMT+TEMWMT
IF(TOTMET.GT.SNOWP)TOTMET=SNOWP
IF(TOTMET.LT.0.0)TOTMET=0.0
ENDIF
SNOWF=SNOWP-TOTMET+RAINN
IF(SNOWF.LT.0.0)SNOWF=0.0
SNOWP=SNOWF
RETURN
END

```

C
 C THIS SUBROUTINE COMPUTES DAILY VALUES OF ACTUAL
 C ACTUAL EVAPOTRANSPIRATION
 C

```

SUBROUTINE EVAT(EDIFF,TEMPER,ALBEDO,SOLRDD,PSYCRO,IGRRO,XXLAI,
+WCFE,AS,THETA,WP,FC,RAINN,TAVRUN,TOTMET,ACEVAT,J1)
IMPLICIT REAL *8(A-H,O-Z)
COMMON /AET/ IT,PT,EVAS,IN,SNOP,ESS
IN=0
TK=5.0*(TEMPER-32.0)/9.0 + 273.0
EXPO=21.255-5304.0/TK
SLOPVP=5304.0*EXP(EXPO)/(TK**2)
SOLNET=(1.0-ALBEDO)*SOLRDD/58.3
POTEV=1.28*SLOPVP*SOLNET/(25.4*(SLOPVP+PSYCRO))
TINTMX=0.05*XXLAI/3.0
IF(TINTMX.EQ.0.0)TINTMX=0.05

```

```

TINTI=0.0
TINEXP=0.0
J2=J1-1
IF((SNOP+RAINN).GE.POTEV)ESS=POTEV
IF((SNOP+RAINN).LT.POTEV)ESS=SNOP+RAINN
IF(TEMPER.LT.32.0)SNOF=SNOP+RAINN-ESS
IF((TEMPER.GT.32.0).AND.(ESS.LE.(TOTMET+RAINN)))SNOF=SNOP-ESS
IF((TEMPER.GT.32.0).AND.(ESS.GT.(TOTMET+RAINN)))SNOF=SNOP-ESS
SNOF=SNOF
IF(IGRRO.EQ.1)THEN
POTEVS=POTEV*EXP(-0.4*XXLAI)
ELSE
POTEVS=POTEV*EXP(-0.4*WCFF)
ENDIF
UPLIM=(9.0*(AS-3.0)**0.42)/25.4
EVAS=POTEV
AVINFF=RAINN+TOTMET-TAVRUN-ESS
IF(AVINFF.LT.0.0)AVINFF=0.0
EDIFF=EDIFF+EVAS-AVINFF
IF(EDIFF.LT.UPLIM)THEN
IN=1
IT=J1
EVAS1=POTEVS
ELSE
IF(J1.EQ.1)PT=0.0
YT=J1-PT
YT1=YT-1.0
IF(YT1.LE.0.0)YT1=0
IF(YT.LE.0.0)YT=0
EVAS2=AS*(SQRT(YT)-SQRT(YT1))/25.4
ENDIF
IF((EDIFF.LT.UPLIM).AND.((ESS+EVAS1).LE.POTEV))EVAS=EVAS1
IF((EDIFF.GE.UPLIM).AND.((ESS+EVAS2).LE.POTEV))EVAS=EVAS2
IF((EDIFF.LT.UPLIM).AND.((ESS+EVAS1).GT.POTEV))EVAS=POTEV-ESS
IF((EDIFF.GE.UPLIM).AND.((ESS+EVAS2).GT.POTEV))EVAS=POTEV-ESS
POTEVP=(POTEV*XXLAI)/3.0
SUM=POTEVP+ESS+EVAS
IF(SUM.LE.POTEV)THEN
EPD=POTEVP
ELSE
EPD=POTEV-ESS-EVAS
ENDIF
EVAP=EPD*(1.20-4.0*EPD+(THETA-WP)/(FC-WP))
IF(EVAP.GT.EPD)EVAP=EPD
ADDEV=EVAS+EVAP
DIFFEV=POTEV-ESS
IF(ADDEV.GT.DIFFEV)THEN
ACEVAT=DIFFEV
ELSE
ACEVAT=ADDEV

```

```

ENDIF
RETURN
END

```

```

C
C THIS SUBROUTINE COMPUTES SURFACE RUNOFF ALONG THE UPPER BOUNDARY
C

```

```

SUBROUTINERUNOFF(TOTMET,PRECIP,HIN,SLOPE,ROUGH,DTR,THETAS,HCONS,H
+ORAIN,HC,FC,II,J1,NEXTUP,NEXUP,LINKNB,ICONF,LCONL,LCONR,NOCOL,FLON
+ET,DX,TEMP,NODE,NONODE,BCOND,COFINT,HCONCOV)
IMPLICIT REAL *8(A-H,O-Z)
COMMON/RUN/HF,PINFIL,T,RUNOF,TOTRUN,THETAP,TAVINF,TAVRUN,TRR,HR
PARAMETER (NOD=100,NOC=50,NS=220,NCC=20)
REAL *8 PRECIP(2,365),HF(NOC,24),HIN(NOC),HP(NOC),SLOPE(NOC),PINFI
+L(NOC),THETAP(NOD),RUNOF(NOC),TOTRUN(NOC,24),HORAIN(24),HR(NOC),QF
+F(NOC),QFP(NOC),QFB(NOC),TEMP(2,365),DDX(NOC),FLOW(NOC),QPP(NOC),H
+FF(NOC),COFINT(NOD,4)
INTEGERNEXUP(NOC),LINKNB(NOD),ICONF(NOD),N1(NOD),LCONL(NOC),LCONR
+(NOC)
TOTLEN=0.0
DO 150 I=1,NEXTUP
XX=ATAN(SLOPE(I))
DDX(I)=DX/(COS(XX))
TOTLEN=TOTLEN+DDX(I)
150 CONTINUE
IF(TOTMET.GT.0.0)PRECIP(II,J1)=PRECIP(II,J1)+TOTMET
NO1COL=NOCOL+1
NO2COL=NOCOL+2
IF(J1.EQ.1)AVRUN=0.0
TAVRUN=0.0
TAVINF=0.0
NOR=0
TRS=0.0
HRON=0.0
OUTFLOW=0.0
TOTFLOW=0.0
TOTDEP=0.0
FLONET=0.0
HHRON=0.0
TR=0.0
140 TR=TR+DTR
J11=TR/DTR
IF(TOTMET.GT.0.0)HORAIN(J11)=HORAIN(J11)+TOTMET/24.0
DTRR=3600.0
TTR=0.0
HOFLOW=0.0

```

```
HDEP=0.0
142 TTR=TTR+DTRR
    J1R=TTR/DTRR
    DO 135 I=1,NEXTUP+2
      HIN(I)=0.0
      IF((II.NE.1).AND.(J1.EQ.1))THEN
        IF(TEMP(II-1,365).LE.32.0)THEN
          IF((J11.EQ.1).AND.(J1R.EQ.1))THEN
            QFP(I)=HIN(I)
          ELSE
            QFP(I)=QFF(I)
          ENDIF
        ELSE
          IF((J11.EQ.1).AND.(J1R.EQ.1))THEN
            QFP(I)=QFB(I)
          ELSE
            QFP(I)=QFF(I)
          ENDIF
        ENDIF
      ELSE
        IF((II.NE.1).AND.(J1.NE.1))THEN
          IF(TEMP(II,J1-1).LE.32.0)THEN
            IF((J11.EQ.1).AND.(J1R.EQ.1))THEN
              QFP(I)=HIN(I)
            ELSE
              QFP(I)=QFF(I)
            ENDIF
          ELSE
            IF((J11.EQ.1).AND.(J1R.EQ.1))THEN
              QFP(I)=QFB(I)
            ELSE
              QFP(I)=QFF(I)
            ENDIF
          ENDIF
        ELSE
          IF((II.EQ.1).AND.(J1.NE.1))THEN
            IF(TEMP(II,J1-1).LE.32.0)THEN
              IF((J11.EQ.1).AND.(J1R.EQ.1))THEN
                QFP(I)=HIN(I)
              ELSE
                QFP(I)=QFF(I)
              ENDIF
            ELSE
              IF((J11.EQ.1).AND.(J1R.EQ.1))THEN
                QFP(I)=QFB(I)
              ELSE
                QFP(I)=QFF(I)
              ENDIF
            ENDIF
          ENDIF
        ENDIF
      ENDIF
    ENDIF
```

```

ENDIF
ELSE
ENDIF
IF((II.EQ.1).AND.(J1.EQ.1))THEN
IF((J11.EQ.1).AND.(J1R.EQ.1))THEN
QFP(I)=HIN(I)
ELSE
QFP(I)=QFF(I)
ENDIF
ELSE
ENDIF
135 CONTINUE
DO 173 I=1,NEXTUP+2
IF((II.EQ.1).AND.(J1.EQ.1))QPP(I)=QFP(I)
173 CONTINUE
DO 155 I=1,NEXTUP
ALPHA=(ROUGH/(1.49*DSQRT(SLOPE(I))))**0.6
HP(I)=(QFP(I)**0.6)*ALPHA
155 CONTINUE
IF(HORAIN(J11).EQ.0.0)TRS=0.0
IF(HORAIN(J11).GT.0.0)TRS=TRS+DTRR
TRS=TRS+DTRR
IF(J11.EQ.1)THEN
TT=TTR
ELSE
IF((J11.NE.1).AND.(HORAIN(J11-1).EQ.0.0).AND.(HORAIN(J11).EQ.0.0))
+THEN
TT=TTR
ELSE
IF((J11.NE.1).AND.(HORAIN(J11-1).GT.0.0).AND.(HORAIN(J11).GT.0.0))
+THEN
TT=TRS
ELSE
IF((J11.NE.1).AND.(HORAIN(J11-1).EQ.0.0).AND.(HORAIN(J11).GT.0.0))
+THEN
TT=TRS
ELSE
IF((J11.NE.1).AND.(HORAIN(J11-1).GT.0.0).AND.(HORAIN(J11).EQ.0.0))
+THEN
TT=TTR
ELSE
ENDIF
ENDIF
ENDIF
ENDIF
ENDIF
145 DO 118 I=1,NEXTUP
136 NN=NEXUP(I)
NNL=LCONL(I)
NNR=LCONR(I)

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THETDF=THETAS-THETAP(NN)
IF(HDF.LT.0.0)HDF=0.0
XX=ATAN(SLOPE(I))
COSX=COS(XX)
HDF=(HP(I)+HC/12.0)*COSX
IF(HDF.LT.0.0)HDF=0.0
SORP=COSX*DSQRT(2.0*HCONCOV*THETDF*HDF/(24.0*3600.0))
CAPINL=0.5*SORP/(SQRT(TT))
IF((HORAIN(J11)/(12.0*3600.0)).LE.CAPINL)THEN
PINFIL(I)=HORAIN(J11)/(12.0*3600.0)
PFLOW=0.0
ELSE
PINFIL(I)=CAPINL
PFLOW=HORAIN(J11)/(12.0*3600.0)-PINFIL(I)
ENDIF
HRANN=HORAIN(J11)/(12.0)
FLOW(I)=PFLOW
118 CONTINUE
DO 112 I=1,NEXTUP
NNL=LCONL(I)
NNR=LCONR(I)
IF(I.EQ.1)THEN
ALPHA=(ROUGH/(1.49*SQRT(SLOPE(I))))**0.6
AVQ=(QFP(I)+QPP(I))/2.0
QFNL2=DTRR*(HORAIN(J11)/(12.0*3600.0))*(AVQ**0.4)/(ALPHA*0.6)
BCONDR=1.49*DSQRT(SLOPE(I))*1.67*((HORAIN(J11)/(12.0*3600.0))**0.6
+7)*BCOND/ROUGH
QFNL3=DTRR*BCONDR*(AVQ**0.4)/(ALPHA*0.6)
QFF(NNL)=QFP(NNL)+QFNL2-QFNL3
DDX(NNL)=DDX(I)
FLOW(NNL)=FLOW(I)
ELSE
IF(I.EQ.NEXTUP)THEN
ALPHA=(ROUGH/(1.49*SQRT(SLOPE(I))))**0.6
AVQ=(QFP(NNR)+QPP(NNR))/2.0
QFNL2=DTRR*(HORAIN(J11)/(12.0*3600.0))*(AVQ**0.4)/(ALPHA*0.6)
BCONDR=1.49*DSQRT(SLOPE(I))*1.67*((HORAIN(J11)/(12.0*3600.0))**0.6
+7)*BCOND/ROUGH
QFNL3=DTRR*BCONDR*(AVQ**0.4)/(ALPHA*0.6)
QFF(NNR)=QFP(NNR)+QFNL2-QFNL3
DDX(NNR)=DDX(I)
FLOW(NNR)=FLOW(I)
ELSE
ENDIF
ENDIF
112 CONTINUE
DO 116 I=1,NEXTUP
NN=NEXUP(I)
NNL=LCONL(I)
NNR=LCONR(I)

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IF(ICONF(NN).EQ.3)GOTO 116
ALPHA=(ROUGH/(1.49*SQRT(SLOPE(I))))**0.6
IF((ICONF(NNL).EQ.3).AND.(I.NE.1))THEN
ALPHA=(ROUGH/(1.49*SQRT(SLOPE(NNL))))**0.6
AVQ=(QFP(NNL)+QPP(NNL))/2.0
QFNL2=DTRR*(HORAIN(J11)/(12.0*3600.0))*(AVQ**0.4)/(ALPHA*0.6)
BCONDR=1.49*DSQRT(SLOPE(I))*1.67*((HORAIN(J11)/(12.0*3600.0))**0.6
+7)*BCOND/ROUGH
QFNL3=DTRR*BCONDR*(AVQ**0.4)/(ALPHA*0.6)
QFF(NNL)=QFP(NNL)+QFNL2-QFNL3
ELSE
ENDIF
VOL1=DTRR*QFF(NNL)/(COFINT(NN,4)*DDX(NNL))
AVG=(QFP(I)+QPP(I))/2.0
IF(AVG.LE.0.0)THEN
VOL2=0.0
ELSE
VOL2=ALPHA*0.6*QFP(I)*(AVG**(-0.4))
ENDIF
VOL3=DTRR*FLOW(I)
VVOL=VOL1+VOL2+VOL3
IF(AVG.LE.0.0)THEN
DELAX=DTRR/(COFINT(NN,4)*DDX(NNL))
ELSE
DELAX=DTRR/(COFINT(NN,4)*DDX(NNL))+ALPHA*0.6*(AVG**(-0.4))
ENDIF
QFF(I)=VVOL/DELAX
IF(QFF(I).LT.0.0)QFF(I)=0.0
IF(J11.EQ.24)QFB(I)=QFF(I)
116 CONTINUE
IK=NEXTUP+1
170 IK=IK-1
NN=NEXUP(IK)
NNR=LCONR(IK)
IF(ICONF(NN).NE.3)GOTO 171
IF((ICONF(NNR).NE.3).AND.(IK.NE.NEXTUP))THEN
ALPHA=(ROUGH/(1.49*SQRT(SLOPE(NNR))))**0.6
AVQ=(QFP(NNR)+QPP(NNR))/2.0
QFNL2=DTRR*(HORAIN(J11)/(12.0*3600.0))*(AVQ**0.4)/(ALPHA*0.6)
BCONDR=1.49*DSQRT(SLOPE(I))*1.67*((HORAIN(J11)/(12.0*3600.0))**0.6
+7)*BCOND/ROUGH
QFNL3=DTRR*BCONDR*(AVQ**0.4)/(ALPHA*0.6)
QFF(NNR)=QFP(NNR)+QFNL2-QFNL3
ELSE
ENDIF
ALPHA=(ROUGH/(1.49*SQRT(SLOPE(IK))))**0.6
VOL1=DTRR*QFF(NNR)/(COFINT(NN,2)*DDX(NNR))
AVG=(QFP(IK)+QPP(IK))/2.0
IF(AVG.LE.0.0)THEN
VOL2=0.0

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ELSE
VOL2=ALPHA*0.6*QFP(IK)*(AVG**(-0.4))
ENDIF
VOL3=DTRR*FLOW(IK)
VVOL=VOL1+VOL2+VOL3
IF(AVG.LE.0.0)THEN
DELAX=DTRR/(COFINT(NN,2)*DDX(NNR))
ELSE
DELAX=DTRR/(COFINT(NN,2)*DDX(NNR))+ALPHA*0.6*(AVG**(-0.4))
ENDIF
QFF(IK)=VVOL/DELAX
IF(QFF(IK).LT.0.0)QFF(IK)=0.0
IF(J11.EQ.24)QFB(IK)=QFF(IK)
171 IF(IK.LT.1)GOTO 170
QHF=0.0
FFL=0.0
DO 167 I=1,NEXTUP
ALPHA=(ROUGH/(1.49*SQRT(SLOPE(I))))**0.6
HFF(I)=(QFF(I)**0.6)*ALPHA
NN=NEXUP(I)
THETDF=THETAS-THETAP(NN)
XX=ATAN(SLOPE(I))
COSX=COS(XX)
HDF=(HFF(I)+HC/12.0)*COSX
SORP=COSX*DSQRT(2.0*0.2835*THETDF*HDF/(24.0*3600.0))
CAPINL=0.5*SORP/(DSQRT(TT))
IF((HORAIN(J11)/(12.0*3600.0)).LE.CAPINL)THEN
FFLOW=0.0
ELSE
FFLOW=HORAIN(J11)/(12.0*3600.0)-CAPINL
ENDIF
FFL=FFL+FFLOW
QHF=QHF+QFF(I)
167 CONTINUE
QHF=QHF/NEXTUP
FFL=FFL/NEXTUP
OUT=QHF*DTRR
IF(HORAIN(J11).GT.0.0)OUTFLOW=OUTFLOW+OUT
HOFLOW=HOFLOW+OUT
HDEP=HDEP+FFL*12.0*DTRR
FLONET=FLONET+FFL*DTRR*12.0
HRON=HRON+DTRR*12.0*HORAIN(J11)/(12.0*3600.0)
DO 172 I=1,NEXTUP
QPP(I)=QFP(I)
172 CONTINUE
IF(TTR.LT.3600.0)GOTO 142
AVFLOW=HOFLOW/3600.0
137 AVINFI=0.0
AVRUN=0.0
121 DO 156 I=1,NEXTUP

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NN=NEXUP(I)
AVINFI=AVINFI+PINFIL(I)
ALPHA=(ROUGH/(1.49*SQRT(SLOPE(I))))**0.6
HF(I,J11)=(AVFLOW**0.6)*ALPHA*12.0
XM=5.0/3.0
TOTRUN(I,J11)=1.49*((HF(I,J11)/12.0)**XM)*SQRT(SLOPE(I))/ROUGH
156 CONTINUE
AVINFI=AVINFI/NEXTUP
AVRUN=AVRUN/NEXTUP
NOR=NOR+1
TAVRUN=TAVRUN+AVRUN
TAVINF=TAVINF+AVINFI
TOTDEP=TOTDEP+HDEP
TOTFLOW=TOTFLOW+HOFLOW
HHRON=HHRON+HORAIN(J11)
IF(TR.LT.24.0)GOTO 140
DEPTH=12.0*TOTFLOW/TOTLEN
IF(DEPTH.GT.1.0)WRITE(3,*)'DEPTH=',DEPTH,'J1=',J1,'P=',PRECIP(II,J
+1)
TRUN=0.0
DO 111 I=1,NEXTUP
HR(I)=0.0
N1(I)=0
DO 117 J=1,24
HR(I)=HR(I)+HF(I,J)
IF(HR(I).NE.0.0)N1(I)=N1(I)+1
117 CONTINUE
IF(N1(I).EQ.0)N1(I)=1
HR(I)=HR(I)/N1(I)
TRUN=TRUN+HR(I)
111 CONTINUE
IF(NOR.EQ.0)NOR=1
TRUN=TRUN/NEXTUP
TAVRUN=TOTDEP
RETURN
END

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C

C THIS SUBROUTINE USES GAUSS-SEIDEL ITERATION TO
C SOLVE THE UNSATURATED MOISTURE-FLOW EQUATION

C

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SUBROUTINEFDMITER(THETIN,SLOPMC,HCONS,SAIE,THETAS,LINKNT,LINKNR,L
+INKNB,LINKNL,ICONF,THETA F,DT,DX,DZ,J1,NONODE,NEXTUP,NEXUP,II,NODE,
+PCP,HR,TEMPER,RAINN,SOLRDD,IGRRO,XXLAI,WCFE,TAVRUN,THETAP,COFINT,M
+OND)
IMPLICIT REAL*8(A-H,O-Z)

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COMMON /AET/ IT,PT,EVAS,IN,SNOP,ESS
COMMON /GAUMAIN/ ALBEDO,PSYCRO,AS,WP,FC,BC,SNOWP,ACEVAT,DIFFUS,
+HCON
PARAMETER (NOD=100,NEX=50,NO=100)
REAL *8 DIFFUS(NOD,4),THETAP(NOD),THETPP(NOD),THETAF(NOD),BC(NEX),
+HCON(NOD,4),HR(NEX),THETAB(NOD),COFINT(NOD,4)
INTEGER LINKNT(NOD),LINKNB(NOD),LINKNR(NOD),LINKNL(NOD),ICONF(NOD)
+,NEXUP(NEX),NODE(NO),MOND(NOD)
EXTERNAL EVAT
CONV=0.0
DO 30 I=1,NONODE
NN=NODE(I)
NNT=LINKNT(NN)
IF((II.EQ.1).AND.(J1.EQ.1))THEN
IF(NNT.NE.0)THEN
IF((MOND(NN).EQ.1).AND.(MOND(NNT).EQ.0))THETAP(I)=THETIN
IF((MOND(NN).EQ.1).AND.(MOND(NNT).EQ.1))THETAP(I)=THETAS
IF((MOND(NN).EQ.0).AND.(MOND(NNT).EQ.0))THETAP(I)=THETIN
ELSE
IF(NNT.EQ.0)THEN
IF(MOND(NN).EQ.1)THETAP(I)=THETAS
IF(MOND(NN).EQ.0)THETAP(I)=THETIN
ELSE
ENDIF
ENDIF
THETPP(I)=THETAP(I)
ELSE
IF(NNT.NE.0)THEN
IF((MOND(NN).EQ.1).AND.(MOND(NNT).EQ.1))THETPP(I)=THETAS
IF((MOND(NN).EQ.1).AND.(MOND(NNT).EQ.0))THETPP(I)=THETAP(I)
IF((MOND(NN).EQ.0).AND.(MOND(NNT).EQ.0))THETPP(I)=THETAP(I)
IF((MOND(NN).EQ.1).AND.(MOND(NNT).EQ.1))THETAP(I)=THETAS
IF((MOND(NN).EQ.1).AND.(MOND(NNT).EQ.0))THETAP(I)=THETAF(I)
IF((MOND(NN).EQ.0).AND.(MOND(NNT).EQ.0))THETAP(I)=THETAF(I)
ELSE
IF(NNT.EQ.0)THEN
IF(MOND(NN).EQ.1)THETPP(I)=THETAS
IF(MOND(NN).EQ.1)THETAP(I)=THETAS
IF(MOND(NN).EQ.0)THETPP(I)=THETAP(I)
IF(MOND(NN).EQ.0)THETAP(I)=THETAF(I)
ELSE
ENDIF
ENDIF
ENDIF
30 CONTINUE
EX1=2.0*SLOPMC+3.0
EX2=SLOPMC+2.0
EX3=SLOPMC*HCONS*SAIE/THETAS
DO 50 I=1,NONODE
NN=NODE(I)

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NNT=LINKNT(NN)
IF(NNT.EQ.0)THEN
TP1=(THETAP(I)+THETPP(I))/2.0
ELSE
DO 51 J=1,NONODE
IF(NODE(J).EQ.NNT)NT=J
51 CONTINUE
TP1=(THETAP(I)+THETPP(I)+THETAP(NT)+THETPP(NT))/4.0
ENDIF
HCON(I,1)=HCONS*(TP1/THETAS)**EX1
DIFFUS(I,1)=-EX3*(TP1/THETAS)**EX2
NNR=LINKNR(NN)
IF(NNR.EQ.0)THEN
TP2=(THETAP(I)+THETPP(I))/2.0
ELSE
DO 52 J=1,NONODE
IF(NODE(J).EQ.NNR)NR=J
52 CONTINUE
TP2=(THETAP(I)+THETPP(I)+THETAP(NR)+THETPP(NR))/4.0
ENDIF
HCON(I,2)=HCONS*(TP2/THETAS)**EX1
DIFFUS(I,2)=-EX3*(TP2/THETAS)**EX2
NNB=LINKNB(NN)
IF(ICONF(NNB).EQ.4)THEN
TP3=(THETAP(I)+THETPP(I))/2.0
ELSE
DO 53 J=1,NONODE
IF(NODE(J).EQ.NNB)NB=J
53 CONTINUE
TP3=(THETAP(I)+THETPP(I)+THETAP(NB)+THETPP(NB))/4.0
ENDIF
HCON(I,3)=HCONS*(TP3/THETAS)**EX1
DIFFUS(I,3)=-EX3*(TP3/THETAS)**EX2
NNL=LINKNL(NN)
IF(NNL.EQ.0)THEN
TP4=(THETAP(I)+THETPP(I))/2.0
ELSE
DO 54 J=1,NONODE
IF(NODE(J).EQ.NNL)NL=J
54 CONTINUE
TP4=(THETAP(I)+THETPP(I)+THETAP(NL)+THETPP(NL))/4.0
ENDIF
HCON(I,4)=HCONS*(TP4/THETAS)**EX1
DIFFUS(I,4)=-EX3*(TP4/THETAS)**EX2
50 CONTINUE
IF(J1.EQ.1)SNOP=SNOWP
IF(J1.EQ.1)ESS=0.0
IF(J1.EQ.1)EDIFF=0.0
THETA=0.0
DO 61 I=1,NEXTUP

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      NN=NEXUP(I)
      DO 62 J=1,NONODE
      IF(NODE(J).EQ.NN)NU=J
62  CONTINUE
      THETA=THETA+THETAP(NU)
61  CONTINUE
      THETA=THETA/NEXTUP
      CALL EVAT(EDIFF,TEMPER,ALBEDO,SOLRDD,PSYCRO,IGRRO,XXLAI,WCFF,AS,TH
+ETA,WP,FC,RAINN,TAVRUN,TOTMET,ACEVAT,J1)
      DO 56 I=1,NEXTUP
      NN=NEXUP(I)
      DO 57 J=1,NONODE
      IF(NODE(J).EQ.NN)NU1=J
57  CONTINUE
      BC(I)=(RAINN+TOTMET-TAVRUN-ACEVAT)/12.0
      IF(TEMPER.LE.32.0)BC(I)=-ACEVAT/12.0
      IF(BC(I).GT.0.5)THEN
      ELSE
      ENDIF
56  CONTINUE
      DO 21 I=1,NONODE
      THETAF(I)=THETAP(I)
21  CONTINUE
      ITER=0
20  TDIFF=0.0
      ITER=ITER+1
      DO 35 I=1,NONODE
      THETAB(I)=THETAF(I)
35  CONTINUE
      DO 45 I=1,NONODE
      VOL=(THETAF(I)-THETAP(I))/DT
      NN=NODE(I)
      NNR=LINKNR(NN)
      NNL=LINKNL(NN)
      NNT=LINKNT(NN)
      NNB=LINKNB(NN)
      IF(NNR.NE.0)THEN
      DO 46 J=1,NONODE
      IF(NODE(J).EQ.NNR)NR=J
46  CONTINUE
      ELSE
      ENDIF
      IF(NNL.NE.0)THEN
      DO 47 J=1,NONODE
      IF(NODE(J).EQ.NNL)NL=J
47  CONTINUE
      ELSE
      ENDIF
      IF(NNT.NE.0)THEN
      DO 48 J=1,NONODE

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      IF(NODE(J).EQ.NNT)NT=J
48  CONTINUE
      ELSE
      ENDIF
      IF(NNB.NE.0)THEN
      DO 49 J=1,NONODE
      IF(NODE(J).EQ.NNB)NB=J
49  CONTINUE
      ELSE
      ENDIF
      IF((ICONF(NN).EQ.1).OR.(ICONF(NN).EQ.2).OR.(ICONF(NN).EQ.3))THEN
      DO 163 J=1,NEXTUP
      IF(NEXUP(J).EQ.NN)NU=J
163 CONTINUE
      ELSE
      ENDIF
      IF(ICONF(NN).EQ.1)THEN
      IF(CONV.EQ.1.0)THETAF(NB)=THETAF(I)
      IF((ICONF(NNB).EQ.4).OR.(MOND(NNB).EQ.1))THETAF(NB)=THETAF(I)
      FLUXZB=DIFFUS(I,3)*(THETAF(NB)-THETAF(I))
      FLUXX=0.0
      FLUXZ=FLUXZB/(COFINT(NN,3)*DZ)
      GRAVIT=HCON(I,3)/(COFINT(NN,3)*DZ)
      BFLUX=BC(NU)/(COFINT(NN,3)*DZ)
      DELAXZ=DIFFUS(I,3)/((COFINT(NN,3)*DZ)**2)
      DELAX=1.0/DT+DELAXZ
      RES=-VOL+FLUXZ-GRAVIT+BFLUX
      ELSE
      IF(ICONF(NN).EQ.3)THEN
      NNR=LINKNR(NN)
      IF((LINKNB(NN).EQ.0).OR.(ICONF(NNB).EQ.4).OR.(CONV.EQ.1.0).OR.(MON
+D(NNB).EQ.1))THEN
      THNB=THETAF(I)
      THNP=THETAP(I)
      ELSE
      THNB=THETAF(NB)
      THNP=THETAP(NB)
      ENDIF
      IF(CONV.EQ.1.0)THETAF(NR)=THETAF(I)
      IF((ICONF(NNR).EQ.4).OR.(MOND(NNR).EQ.1))THETAF(NR)=THETAF(I)
      DELAXZ=DIFFUS(I,3)/((COFINT(NN,3)*DZ)**2)
      BFLUX=BC(NU)/(COFINT(NN,3)*DZ)
      FLUXZB=DIFFUS(I,3)*(THNB-THETAF(I))
      FLUXXR=DIFFUS(I,2)*(THETAF(NR)-THETAF(I))
      FLUXX=FLUXXR/((COFINT(NN,2)*DX)**2)
      FLUXZ=FLUXZB/((COFINT(NN,3)*DZ)**2)
      DELAXX=DIFFUS(I,2)/((COFINT(NN,2)*DX)**2)
      GRAVIT=HCON(I,3)/(COFINT(NN,3)*DZ)
      DELAX=1.0/DT+DELAXX+DELAXZ
      RES=-VOL+FLUXX+FLUXZ+BFLUX-GRAVIT

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ELSE
IF(ICONF(NN).EQ.2)THEN
NNL=LINKNL(NN)
IF((LINKNB(NN).EQ.0).OR.(ICONF(NNB).EQ.4).OR.(CONV.EQ.1.0).OR.(MON
+D(NNB).EQ.1))THEN
THNB=THETAF(I)
THNP=THETAP(I)
ELSE
THNB=THETAF(NB)
THNP=THETAP(NB)
ENDIF
IF(CONV.EQ.1.0)THETAF(NL)=THETAF(I)
IF((ICONF(NNL).EQ.4).OR.(MOND(NNL).EQ.1))THETAF(NL)=THETAF(I)
BFLUX=BC(NU)/(COFINT(NN,3)*DZ)
FLUXZB=DIFFUS(I,3)*(THNB-THETAF(I))
DELAXZ=DIFFUS(I,3)/((COFINT(NN,3)*DZ)**2)
FLUXXL=DIFFUS(I,4)*(THETAF(NL)-THETAF(I))
FLUXX=FLUXXL/((COFINT(NN,4)*DX)**2)
FLUXZ=FLUXZB/((COFINT(NN,3)*DZ)**2)
DELAXX=DIFFUS(I,4)/((COFINT(NN,4)*DX)**2)
DELAX=1.0/DT+DELAXZ+DELAXX
GRAVIT=HCON(I,3)/(COFINT(NN,3)*DZ)
RES=-VOL+FLUXX+FLUXZ+BFLUX-GRAVIT
ELSE
IF(ICONF(NN).EQ.5)THEN
IF((ICONF(NNB).EQ.4).OR.(MOND(NNB).EQ.1))THETAF(NB)=THETAF(I)
IF((ICONF(NNR).EQ.4).OR.(MOND(NNR).EQ.1))THETAF(NR)=THETAF(I)
IF((ICONF(NNL).EQ.4).OR.(MOND(NNL).EQ.1))THETAF(NL)=THETAF(I)
FLUXXR=DIFFUS(I,2)*(THETAF(NR)-THETAF(I))/(COFINT(NN,2)*DX)
FLUXXL=DIFFUS(I,4)*(THETAF(I)-THETAF(NL))/(COFINT(NN,4)*DX)
FLUXX=2.0*(FLUXXR-FLUXXL)/((COFINT(NN,2)+COFINT(NN,4))*DX)
FLUXZT=DIFFUS(I,1)*(THETAF(NT)-THETAF(I))/(COFINT(NN,1)*DZ)
FLUXZB=DIFFUS(I,3)*(THETAF(I)-THETAF(NB))/(COFINT(NN,3)*DZ)
FLUXZ=2.0*(FLUXZT-FLUXZB)/((COFINT(NN,1)+COFINT(NN,3))*DZ)
GRAVIT=2.0*(HCON(I,1)-HCON(I,3))/((COFINT(NN,1)+COFINT(NN,3))*DZ)
DELX1=DIFFUS(I,2)/(COFINT(NN,2)*DX)
DELX2=DIFFUS(I,4)/(COFINT(NN,4)*DX)
DELAXX=2.0*(DELX1+DELX2)/(DX*(COFINT(NN,2)+COFINT(NN,4)))
DELZ1=DIFFUS(I,1)/(COFINT(NN,1)*DZ)
DELZ2=DIFFUS(I,3)/(COFINT(NN,3)*DZ)
DELAXZ=2.0*(DELZ1+DELZ2)/(DZ*(COFINT(NN,1)+COFINT(NN,3)))
DELAX=1.0/DT+DELAXX+DELAXZ
RES=-VOL+FLUXX+FLUXZ-GRAVIT
ELSE
ENDIF
ENDIF
ENDIF
ENDIF
RELAX=1.0/DELAX
THETAF(I)=THETAF(I)+RES*RELAX

```

```

45  CONTINUE
    DO 60 I=1,NONODE
      IF(THETAF(I).GT.THETAS)THETAF(I)=THETAS
      IF(THETAF(I).LT.FC)THETAF(I)=FC
60  CONTINUE
    DO 90 I=1,NONODE
      TDIFF=TDIFF+ABS(THETAF(I)-THETAB(I))
90  CONTINUE
      IF(ITER.EQ.1)THEN
        BDIFF=TDIFF
      ELSE
        ENDIF
      CONV=0.0
      IF(ITER.GT.150)WRITE(*,*)'TOO MANY ITERATIONS IN G.S.I. METHOD'
      IF((TDIFF.GT.BDIFF).OR.(ITER.GT.20))THEN
        CONV=1.0
      ELSE
        BDIFF=TDIFF
      ENDIF
      IF(TDIFF.GT.0.04)GOTO 20
      IF(IN.EQ.1)PT=IT
      RETURN
    END

```

C

C THIS SUBROUTINE COMPUTES THE COEFFICIENT MATRICES
C TO SOLVE THE SATURATED FLOW EQUATION

C

```

SUBROUTINE SATELIM(NODTOT,NOBAR,NOROW,NOCOL,ICRNO,ICONF,IBET,
+LINKNB,NEXTUP,NEXUP,THETAS,BC,HDIN,ICOL,LCONL,LCONR,HCONS,
+CONDB,EFPORO,BCONL,BCONR,DT,DX,DZ,THICK,TAVRUN,ACEVAT,
+PRCPP,NODE,NONODE,SLOPBR,THETAF,II,J1,GRAD,THETAM,MOND,
+LINKNT,HDD,HDP,TAVSOR,HCOL,SLOPMC,SAIE,THETAP,VDISTB,HCON,
+DIFFUS,FC,LINKNR,LINKNL,DRC)
  IMPLICIT REAL *8(A-H,O-Z)
  COMMON /COFMN/ COEFM
  PARAMETER (NC=50,NT=100,NEX=50,NR=20)
  INTEGER LCONL(NC),LCONR(NC),ICOL(NC),IBET(NT),ICRNO(NR,NC),ICONF(N
+T),NEXUP(NEX),MOND(NT),NUM(NC),LINKNB(NT),NODE(NT),LINKNT(NT),NODU
+N(NC),NODCOL(NR,NC),IDN(NC),LINKNL(NT),LINKNR(NT)
  REAL *8 THICK(NC),HDIN(NC),SOURCE(NC),HDP(NC),HPP(NC),HDD(NC),HPS(
+NC,2),COEFM(NT,NT),SLOPBR(NC),HCOL(1,NEX),THETAM(NT),THETAF(NT),BC
+(NEX),GRAD(1,NC),SOURCHK(NC),THETAP(NT),VDISTB(NT,1),HCON(NT,4),DI
+FFUS(NT,4),DRAIN(NC)
  EXTERNAL ELIM,SOURCH
  SLOPL=0.0

```

```

SLOPR=0.0
DO 12 I=1,NOCOL/2
SLOPL=SLOPL+SLOPBR(I)
12 CONTINUE
SLOPL=SLOPL*2/NOCOL
DO 20 I=NOCOL/2,NOCOL
SLOPR=SLOPR+SLOPBR(I)
20 CONTINUE
SLOPR=SLOPR*2/NOCOL
DO 35 I=1,NOCOL
IF((II.EQ.1).AND.(J1.EQ.1))THEN
HDP(I)=HDIN(I)
HPP(I)=HDP(I)
ELSE
HPP(I)=HDP(I)
HDP(I)=HDD(I)
ENDIF
35 CONTINUE
TAVSOR=0.0
DO 426 I=1,NOCOL
SOURCE(I)=0.0
SOURCHK(I)=0.0
NUM(I)=0
426 CONTINUE
HHDP=0.0
HHDPP=0.0
DO 427 I=1,NOCOL
HHDP=HHDP+HDP(I)
HHDPP=HHDPP+HPP(I)
427 CONTINUE
CALL SOURCH(NOROW,NOCOL,MOND,THETA,THETAP,VDISTB,SOURCHK,ICRNO,
+NONODE,NODE,FC,DRAIN,NODUN,NODCOL,BC,LINKNB,DIFFUS,DX,LINKNR,
+LINKNL,DZ,THETAS,SAIE,HCONS,DRC)
DO 575 I=1,NOROW
DO 575 J=1,NOCOL
NN=ICRNO(I,J)
IF(NN.EQ.0)GOTO 575
DO 430 K=1,NONODE
IF(NODE(K).EQ.NN)N1=K
430 CONTINUE
IF(MOND(NN).EQ.1)GOTO 575
NNT=LINKNT(NN)
NNB=LINKNB(NN)
IF(NNB.EQ.0)THEN
THNB=THETA(N1)
ELSE
IF(ICONF(NNB).NE.4)THEN
DO 431 K=1,NONODE
IF(NODE(K).EQ.NNB)NB1=K
431 CONTINUE

```

```

      THNB=THETA(NB1)
      ELSE
      THNB=THETA(N1)
      ENDIF
      ENDIF
      SOURCE(J)=SOURCE(J)+HCON(N1,3)-DIFFUS(N1,3)*(THNB-THETA(N1))/DZ
      NUM(J)=NUM(J)+1
575  CONTINUE
      DO 580 I=1,NOROW
      DO 580 J=1,NOCOL
      NN=ICRNO(I,J)
      IF(NN.EQ.0)GOTO 580
      IF(ICONF(NN).EQ.4)GOTO 580
      IF(ICONF(NN).EQ.5)GOTO 580
      IF(NUM(J).EQ.0)THEN
      DO 434 K=1,NEXTUP
      IF(NEXUP(K).EQ.NN)NU=K
434  CONTINUE
      IF(BC(NU).GT.HCON(NU,1))THEN
      SOURCE(J)=SOURCE(J)+HCON(NU,1)
      NUM(J)=NUM(J)+1
      ELSE
      SOURCE(J)=SOURCE(J)+BC(NU)
      NUM(J)=NUM(J)+1
      ENDIF
      ELSE
      GOTO 580
      ENDIF
580  CONTINUE
      TAVIN=(PRCPP-TAVRUN-ACEVAT)/12.0
433  DO 432 I=1,NOCOL
      SOURCE(I)=SOURCE(I)/NUM(I)
      IF(SOURCE(I).LT.0.0)SOURCE(I)=0.0
      IF(SOURCE(I).GT.(0.1*PRCPP/12.0))SOURCE(I)=SOURCHK(I)
      IF(SOURCE(I).LT.(0.05*PRCPP/12.0))SOURCE(I)=DRAIN(I)
      TAVSOR=TAVSOR+SOURCE(I)
432  CONTINUE
      TAVSOR=12.0*TAVSOR/NOCOL
      DO 10 I=1,NOCOL
     >NNL=LCONL(I)
     >NNR=LCONR(I)
      IF(ICOL(NNL).EQ.2)THEN
     >HPS(I,2)=(HDP(I)+HPP(I))/2.0
     >HPS(I,1)=(HDP(I)+HPP(I)+HDP(NNR)+HPP(NNR))/4.0
      ELSE
      IF(ICOL(NNR).EQ.2)THEN
     >HPS(I,1)=(HDP(I)+HPP(I))/2.0
     >HPS(I,2)=(HDP(I)+HPP(I)+HDP(NNL)+HPP(NNL))/4.0
      ELSE
     >HPS(I,1)=(HDP(I)+HPP(I)+HDP(NNR)+HPP(NNR))/4.0

```

```

HPS(I,2)=(HDP(I)+HPP(I)+HDP(NNL)+HPP(NNL))/4.0
ENDIF
ENDIF
10 CONTINUE
DO 45 I=1,NOCOL
DO 45 J=1,NOCOL
COEFM(I,J)=0.0
45 CONTINUE
DO 15 I=1,NOCOL
IF(I.LT.(NOCOL/2))SLOPB=SLOPL
IF(I.GE.(NOCOL/2))SLOPB=SLOPR
NNL=LCONL(I)
NNR=LCONR(I)
IF(ICOL(NNL).EQ.2)THEN
A=(-HCONS*HPS(I,1))/DX**2-CONDB/THICK(I)-EFPORO/DT
B=0.0
C=HCONS*HPS(I,1)/DX**2
P=-SOURCE(I)+CONDB-EFPORO*HDP(I)/DT+HCONS*HPS(I,2)*BCONL/DX-HCONS*
+SLOPB*BCONL
COEFM(I,I)=A
COEFM(I,NNR)=C
HDD(I)=P
ELSE
IF(ICOL(NNR).EQ.2)THEN
A=(-HCONS*HPS(I,2))/DX**2-CONDB/THICK(I)-EFPORO/DT
B=HCONS*HPS(I,2)/DX**2
C=0.0
P=-SOURCE(I)+CONDB-EFPORO*HDP(I)/DT-HCONS*HPS(I,1)*BCONR/DX-HCONS*
+SLOPB*BCONR
COEFM(I,I)=A
COEFM(I,NNL)=B
HDD(I)=P
ELSE
IF((ICOL(NNL).EQ.1).AND.(ICOL(NNR).EQ.1))THEN
A=HCONS*(-HPS(I,1)-HPS(I,2))/DX**2-CONDB/THICK(I)-EFPORO/DT
B=HCONS*HPS(I,2)/DX**2-HCONS*SLOPB/(2.0*DX)
C=HCONS*HPS(I,1)/DX**2+HCONS*SLOPB/(2.0*DX)
P=-SOURCE(I)+CONDB-EFPORO*HDP(I)/DT
COEFM(I,I)=A
COEFM(I,NNL)=B
COEFM(I,NNR)=C
HDD(I)=P
ELSE
ENDIF
ENDIF
ENDIF
15 CONTINUE
CALL ELIM(HDD,NOCOL)
DO 991 I=1,NOBAR
DO 991 J=1,NOCOL

```

```

NNL=LCONL(J)
NNR=LCONR(J)
IF(HDD(J).LT.0.0)HDD(J)=0.0
IF(HDD(J).GT.HCOL(I,J))HDD(J)=HCOL(I,J)
GRAD(I,J)=ABS(HDD(NNL)-HDD(NNR))/(2.0*DX)
991 CONTINUE
RETURN
END

```

C

C THIS SUBROUTINE COMPUTES LEACHATE ACCRETION IN THE

C UNSATURATED ZONE

C

```

SUBROUTINE SOURCH(NOROW,NOCOL,MOND,THETA F,THETA P,VDISTB,
+SOURCHK,ICRNO,NONODE,NODE,FC,DRAIN,NODUN,NODCOL,BC,LINKNB,
+DIFFUS,DX,LINKNR,LINKNL,DZ,THETAS,SAIE,HCONS,DRC)
IMPLICIT REAL *8(A-H,O-Z)
PARAMETER(NT=100,NC=50,NR=20)
REAL *8 VDISTB(NT,1),THETA F(NT),THETA P(NT),SOURCHK(NC),SOLRIGHT(NR
+),CCF(NR,NR),CCP(NR,NR),DRAIN(NC),BC(NC),STORE(NC,NR),DIFFUS(NT,4)
INTEGER MOND(NT),NODUN(NC),NODCOL(NC,NR),ICRNO(NR,NC),NODE(NT),LIN
+KNB(NT),LINKNR(NT),LINKNL(NT)
DO 30 I=1,NOCOL
NODUN(I)=0
30 CONTINUE
DO 10 I=1,NOROW
DO 10 J=1,NOCOL
NN=ICRNO(I,J)
IF(NN.EQ.0)GOTO 10
IF(MOND(NN).NE.1)THEN
NODUN(J)=NODUN(J)+1
NODCOL(J,NODUN(J))=NN
ELSE
ENDIF
10 CONTINUE
DO 40 I=1,NOCOL
NOUN=NODUN(I)
SOURCHK(I)=0.0
DO 20 J=1,NOUN
NN1=NODCOL(I,J)
DO 432 IJ=1,NONODE
IF(NODE(IJ).EQ.NN1)NI1=IJ
432 CONTINUE
STORPOT=THETAS-THETA P(NI1)
IF(THETA P(NI1).GT.FC)THEN
GRAV=THETAS-FC
SOURCHK(I)=SOURCHK(I)+DRC*(1.0-STORPOT/GRAV)**3

```

```

ELSE
DCHK=HCONS*DABS(SAIE)*(THETAP(NI1)**4.50)
SOURCHK(I)=SOURCHK(I)+DCHK
ENDIF
20 CONTINUE
SOURCHK(I)=SOURCHK(I)/NOUN
40 CONTINUE
DO 80 I=1,NOCOL
NOUN=NODUN(I)
DRAN=0.0
DO 80 J=1,NOUN
IF(J.EQ.NOUN)GOTO 80
NN1=NODCOL(I,J)
NN2=NODCOL(I,J+1)
VVDIST=VDISTB(NN2,1)-VDISTB(NN1,1)
DO 430 IJ=1,NONODE
IF(NODE(IJ).EQ.NN1)NI1=IJ
IF(NODE(IJ).EQ.NN2)NI2=IJ
430 CONTINUE
AREAF=0.5*(THETAF(NI1)+THETAF(NI2))*VVDIST
AREAP=0.5*(THETAP(NI1)+THETAP(NI2))*VVDIST
STORE(I,J)=AREAF-AREAP
80 CONTINUE
DO 50 I=1,NOCOL
NOUN=NODUN(I)-1
IF(NOUN.LT.0)THEN
DRAIN(I)=0.0
GOTO 50
ELSE
ENDIF
IF(NOUN.EQ.0)THEN
NN1=NODCOL(I,NODUN(I))
NN2=LINKNB(NN1)
NNR=LINKNR(NN1)
NNL=LINKNL(NN1)
DO 61 K=1,NONODE
IF(NODE(K).EQ.NN1)N1=K
IF(NODE(K).EQ.NNL)NL1=K
IF(NODE(K).EQ.NNR)NR1=K
61 CONTINUE
DR1=(THETAF(N1)-THETAP(N1))*(VDISTB(NN1,1)-VDISTB(NN2,1))/2.
IF((NNL.EQ.0).OR.(NNR.EQ.0))THEN
DR2=0.0
DR3=0.0
ELSE
DR2=DIFFUS(N1,2)*(THETAF(NR1)-THETAF(N1))*DZ/DX
DR3=DIFFUS(N1,4)*(THETAF(NL1)-THETAF(N1))*DZ/DX
ENDIF
DR= BC(I)-DR1-DR2-DR3
DRAIN(I)=DR

```

```

GOTO 50
ELSE
ENDIF
J=NOUN
60 IF(J.EQ.NOUN)THEN
  NN=NODCOL(I,J)
  NNR=LINKNR(NN)
  NNL=LINKNL(NN)
  DO 63 K=1,NONODE
  IF(NODE(K).EQ.NN)N1=K
  IF(NODE(K).EQ.NNL)NL1=K
  IF(NODE(K).EQ.NNR)NR1=K
63 CONTINUE
  IF((NNL.EQ.0).OR.(NNR.EQ.0))THEN
    DR2=0.0
    DR3=0.0
  ELSE
    DR2=DIFFUS(N1,2)*(THETAF(NR1)-THETAF(N1))*DZ/DX
    DR3=DIFFUS(N1,4)*(THETAF(NL1)-THETAF(N1))*DZ/DX
  ENDIF
  DR=BC(I)-DR2-DR3-STORE(I,J)
  ELSE
    NN=NODCOL(I,J)
    NNL=LINKNL(NN)
    NNR=LINKNR(NN)
    DO 64 K=1,NONODE
    IF(NODE(K).EQ.NN)N1=K
    IF(NODE(K).EQ.NNL)NL1=K
    IF(NODE(K).EQ.NNR)NR1=K
64 CONTINUE
    IF((NNL.EQ.0).OR.(NNR.EQ.0))THEN
      DR2=0.0
      DR3=0.0
    ELSE
      DR2=DIFFUS(N1,2)*(THETAF(NR1)-THETAF(N1))*DZ/DX
      DR3=DIFFUS(N1,4)*(THETAF(NL1)-THETAF(N1))*DZ/DX
    ENDIF
    DR=DR-DR2-DR3-STORE(I,J)
  ENDIF
  J=J-1
  IF(J.GE.1)GOTO 60
  DRAIN(I)=DR
50 CONTINUE
  DO 70 I=1,NOCOL
  IF((SOURCHK(I)).GT.(DRAIN(I)))SOURCHK(I)=DRAIN(I)
70 CONTINUE
  RETURN
END

```

C
 C THIS SUBROUTINE SOLVES SIMULTANEOUS EQUATIONS
 C USING GAUSS-JORDAN ELIMINATION METHOD
 C

```

SUBROUTINE ELIM(COLMAT,NOBEL)
  IMPLICIT REAL *8(A-H,O-Z)
  PARAMETER (NT=100)
  COMMON /COFMN/ COEFF
  REAL*8 COEFF(NT,NT),PP(NT),COLMAT(NT)
  DO 947 I=1,NOBEL
    IF(I.EQ.NOBEL)GOTO 956
    BIG=DABS(COEFF(I,I))
    I1=I+1
    DO 951 L=I1,NOBEL
      IF(DABS(COEFF(L,I)).GT.BIG)BIG=DABS(COEFF(L,I))
951  CONTINUE
    DO 952 L=I,NOBEL
      IF(DABS(COEFF(L,I)).EQ.BIG)NR=L
952  CONTINUE
    DO 953 M=1,NOBEL
      PP(M)=COEFF(I,M)
953  CONTINUE
    NON1=NOBEL+1
    PP(NON1)=COLMAT(I)
    DO 954 M=1,NOBEL
      COEFF(I,M)=COEFF(NR,M)
954  CONTINUE
    COLMAT(I)=COLMAT(NR)
    DO 955 M=1,NOBEL
      COEFF(NR,M)=PP(M)
955  CONTINUE
    COLMAT(NR)=PP(NON1)
956  DEN=COEFF(I,I)
    DO 948 J=1,NOBEL
      COEFF(I,J)=COEFF(I,J)/DEN
948  CONTINUE
    COLMAT(I)=COLMAT(I)/DEN
    DO 949 K=1,NOBEL
      IF(K.EQ.I)GOTO 949
      IF(COEFF(K,I).EQ.0.0)GOTO 949
      CON=COEFF(K,I)
      DO 950 J=1,NOBEL
        IF(DABS(COEFF(I,J)).LT.1.0E-30)THEN
          COEFF(I,J)=1.0E-30
        ELSE
          COEFF(K,J)=COEFF(K,J)-COEFF(I,J)*CON
        ENDIF
950  CONTINUE
    COLMAT(K)=COLMAT(K)-COLMAT(I)*CON
  
```

```

949 CONTINUE
947 CONTINUE
  RETURN
  END

```

```

C
C THIS SUBROUTINE COMPUTES LATERAL FLOW UPON THE BARRIER
C AND VERTICAL FLOW THROUGH THE BARRIER
C
  SUBROUTINE FLOW(NOCOL,NOBAR,HDD,SLOPB,HCONS,J1,II,RAINN,THICK,CON
+DB)
  IMPLICIT REAL *8(A-H,O-Z)
  COMMON /MFLO/ AQV,FLOWL,FLOWR,FLOWVV,FLOWLV,FLOWLI,BEFFICY,
+SMEFICY,EFFICY
  COMMON /MONFO/ GRAD,BIG
  PARAMETER(NT=100,NC=50,NR=50)
  REAL *8 HDD(NC),BIG(1),AVSLPL(1),AVSLPR(1),SLOPB
+R(NC),FLOWL(1),FLOWR(1),GRAD(1,NC),AVGDL(1),AVGDR(1),FLOWVV(1),FLO
+WLV(1),FLOWLI(1),BEFFICY(2),SMEFICY(2),AVHDL(1),AVHDR(1),THICK(NC)
  INTEGER MX(1)
  DO 350 I=1,NOBAR
    AVGDR(I)=0.0
    AVGDL(I)=0.0
350 CONTINUE
    AQV=0.0
    DO 336 I=1,NOCOL
      QV=CONDB*(HDD(I)+THICK(I))/THICK(I)
      AQV=AQV+QV
336 CONTINUE
      AQV=AQV/NOCOL
      IF(RAINN.NE.0.0)EFFICY=100.0-100.0*AQV/RAINN
      IF(RAINN.EQ.0.0)EFFICY=100.0
      IF(RAINN.NE.0.0)THEN
        IF(J1.EQ.1)THEN
          BEFFICY(II)=EFFICY
          SMEFICY(II)=EFFICY
        ELSE
          IF(EFFICY.GT.BEFFICY(II))BEFFICY(II)=EFFICY
          IF(EFFICY.LT.SMEFICY(II))SMEFICY(II)=EFFICY
        ENDIF
      ELSE
        ENDIF
      DO 505 I=1,NOBAR
        BIG(I)=HDD(1)
        DO 505 J=2,NOCOL
          IF(HDD(J).GT.BIG(I))BIG(I)=HDD(J)
505 CONTINUE

```

```

DO 510 I=1,NOBAR
DO 510 J=1,NOCOL
IF(HDD(J).EQ.BIG(I))THEN
MX(I)=J
ELSE
GOTO 510
ENDIF
510 CONTINUE
DO 515 I=1,NOBAR
KN=MX(I)
DO 520 J=1,KN
AVHDL(I)=AVHDL(I)+HDD(J)
AVSLPL(I)=AVSLPL(I)+SLOPBR(J)
AVGDL(I)=AVGDL(I)+GRAD(I,J)
520 CONTINUE
AVHDL(I)=AVHDL(I)/KN
AVSLPL(I)=AVSLPL(I)/KN
AVGDL(I)=AVGDL(I)/KN
DO 525 J=KN,NOCOL
AVHDR(I)=AVHDR(I)+HDD(J)
AVSLPR(I)=AVSLPR(I)+SLOPBR(J)
AVGDR(I)=AVGDR(I)+GRAD(I,J)
525 CONTINUE
KN1=NOCOL-KN+1
AVHDR(I)=AVHDR(I)/KN1
AVSLPR(I)=AVSLPR(I)/KN1
AVGDR(I)=AVGDR(I)/KN1
515 CONTINUE
DO 540 I=1,NOBAR
541 FLOWL(I)=HCONS*BIG(I)*(AVGDL(I)+AVSLPL(I))/2.0
FLOWR(I)=HCONS*BIG(I)*(AVGDR(I)+AVSLPR(I))/2.0
540 CONTINUE
RETURN
END

```

C

C THIS SUBROUTINE COMPUTES MONTHLY VALUES OF THE FLOW PARAMETERS

C

```

SUBROUTINE SELECT(NOBAR,FLOWL,FLOWR,AQV,J1,II,AVTHET,RAIN,AVIN,TAV
+RUN,ACEVAT,TAVSOR,BIG,BBIG,TAVINF,FLONET,CAPINF,FLORUN)
IMPLICIT REAL *8(A-H,O-Z)
COMMON/SILA/FLOMLL,FLOMLR,FLOWMV,BIGLAT,SMALAT,AVMOST,RAINM,
+PERCM,RUNM,EVAM,RECHM,BIGVER,SMAVER
REAL *8 FLOMLL(1,2,12),FLOMLR(1,2,12),FLOWMV(1,2,12),BIGLAT(1),SMA
+LAT(1),FLOWL(1),FLOWR(1),AVMOST(1,2,12),RAINM(2,12),PERCM(2
+,12),RUNM(2,12),EVAM(2,12),RAIN(2,365),RECHM(1,2,12),BIGVER(1),SMA
+VER(1),BIG(1),BBIG(1,12),CAPINF(2,12),FLORUN(2,12)

```

```

DO 585 I=1,NOBAR
IF(J1.LE.31)THEN
IF(J1.EQ.1)BBIG(I,1)=0.0
FLOMLL(I,II,1)=FLOMLL(I,II,1)+FLOWL(I)
FLOMLR(I,II,1)=FLOMLR(I,II,1)+FLOWR(I)
FLOWMV(I,II,1)=FLOWMV(I,II,1)+AQV
AVMOST(I,II,1)=AVMOST(I,II,1)+AVTHET
RECHM(I,II,1)=RECHM(I,II,1)+TAVSOR
IF(BIG(I).GT.BBIG(I,1))BBIG(I,1)=BIG(I)
ELSE
IF((J1.GT.31).AND.(J1.LE.59))THEN
IF(J1.EQ.32)BBIG(I,2)=0.0
IF(J1.EQ.32)AVMOST(I,II,1)=AVMOST(I,II,1)/31.0
FLOMLL(I,II,2)=FLOMLL(I,II,2)+FLOWL(I)
FLOMLR(I,II,2)=FLOMLR(I,II,2)+FLOWR(I)
FLOWMV(I,II,2)=FLOWMV(I,II,2)+AQV
AVMOST(I,II,2)=AVMOST(I,II,2)+AVTHET
RECHM(I,II,2)=RECHM(I,II,2)+TAVSOR
IF(BIG(I).GT.BBIG(I,2))BBIG(I,2)=BIG(I)
ELSE
IF((J1.GT.59).AND.(J1.LE.90))THEN
IF(J1.EQ.60)BBIG(I,3)=0.0
IF(J1.EQ.60)AVMOST(I,II,2)=AVMOST(I,II,2)/28.0
FLOMLL(I,II,3)=FLOMLL(I,II,3)+FLOWL(I)
FLOMLR(I,II,3)=FLOMLR(I,II,3)+FLOWR(I)
FLOWMV(I,II,3)=FLOWMV(I,II,3)+AQV
AVMOST(I,II,3)=AVMOST(I,II,3)+AVTHET
RECHM(I,II,3)=RECHM(I,II,3)+TAVSOR
IF(BIG(I).GT.BBIG(I,3))BBIG(I,3)=BIG(I)
ELSE
IF((J1.GT.90).AND.(J1.LE.120))THEN
IF(J1.EQ.91)BBIG(I,4)=0.0
IF(J1.EQ.91)AVMOST(I,II,3)=AVMOST(I,II,3)/31.0
FLOMLL(I,II,4)=FLOMLL(I,II,4)+FLOWL(I)
FLOMLR(I,II,4)=FLOMLR(I,II,4)+FLOWR(I)
FLOWMV(I,II,4)=FLOWMV(I,II,4)+AQV
AVMOST(I,II,4)=AVMOST(I,II,4)+AVTHET
RECHM(I,II,4)=RECHM(I,II,4)+TAVSOR
IF(BIG(I).GT.BBIG(I,4))BBIG(I,4)=BIG(I)
ELSE
IF((J1.GT.120).AND.(J1.LE.151))THEN
IF(J1.EQ.121)BBIG(I,5)=0.0
IF(J1.EQ.121)AVMOST(I,II,4)=AVMOST(I,II,4)/30.0
FLOMLL(I,II,5)=FLOMLL(I,II,5)+FLOWL(I)
FLOMLR(I,II,5)=FLOMLR(I,II,5)+FLOWR(I)
FLOWMV(I,II,5)=FLOWMV(I,II,5)+AQV
AVMOST(I,II,5)=AVMOST(I,II,5)+AVTHET
RECHM(I,II,5)=RECHM(I,II,5)+TAVSOR
IF(BIG(I).GT.BBIG(I,5))BBIG(I,5)=BIG(I)
ELSE

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IF((J1.GT.151).AND.(J1.LE.181))THEN
IF(J1.EQ.152)BBIG(I,6)=0.0
IF(J1.EQ.152)AVMOST(I,II,5)=AVMOST(I,II,5)/31.0
FLOMLL(I,II,6)=FLOMLL(I,II,6)+FLOWL(I)
FLOMLR(I,II,6)=FLOMLR(I,II,6)+FLOWR(I)
FLOWMV(I,II,6)=FLOWMV(I,II,6)+AQV
AVMOST(I,II,6)=AVMOST(I,II,6)+AVTHET
RECHM(I,II,6)=RECHM(I,II,6)+TAVSOR
IF(BIG(I).GT.BBIG(I,6))BBIG(I,6)=BIG(I)
ELSE
IF((J1.GT.181).AND.(J1.LE.212))THEN
IF(J1.EQ.182)BBIG(I,7)=0.0
IF(J1.EQ.182)AVMOST(I,II,6)=AVMOST(I,II,6)/30.0
FLOMLL(I,II,7)=FLOMLL(I,II,7)+FLOWL(I)
FLOMLR(I,II,7)=FLOMLR(I,II,7)+FLOWR(I)
FLOWMV(I,II,7)=FLOWMV(I,II,7)+AQV
AVMOST(I,II,7)=AVMOST(I,II,7)+AVTHET
RECHM(I,II,7)=RECHM(I,II,7)+TAVSOR
IF(BIG(I).GT.BBIG(I,7))BBIG(I,7)=BIG(I)
ELSE
IF((J1.GT.212).AND.(J1.LE.243))THEN
IF(J1.EQ.213)BBIG(I,8)=0.0
IF(J1.EQ.213)AVMOST(I,II,7)=AVMOST(I,II,7)/31.0
FLOMLL(I,II,8)=FLOMLL(I,II,8)+FLOWL(I)
FLOMLR(I,II,8)=FLOMLR(I,II,8)+FLOWR(I)
FLOWMV(I,II,8)=FLOWMV(I,II,8)+AQV
AVMOST(I,II,8)=AVMOST(I,II,8)+AVTHET
RECHM(I,II,8)=RECHM(I,II,8)+TAVSOR
IF(BIG(I).GT.BBIG(I,8))BBIG(I,8)=BIG(I)
ELSE
IF((J1.GT.243).AND.(J1.LE.273))THEN
IF(J1.EQ.244)BBIG(I,9)=0.0
IF(J1.EQ.244)AVMOST(I,II,8)=AVMOST(I,II,8)/31.0
FLOMLL(I,II,9)=FLOMLL(I,II,9)+FLOWL(I)
FLOMLR(I,II,9)=FLOMLR(I,II,9)+FLOWR(I)
FLOWMV(I,II,9)=FLOWMV(I,II,9)+AQV
AVMOST(I,II,9)=AVMOST(I,II,9)+AVTHET
RECHM(I,II,9)=RECHM(I,II,9)+TAVSOR
IF(BIG(I).GT.BBIG(I,9))BBIG(I,9)=BIG(I)
ELSE
IF((J1.GT.273).AND.(J1.LE.304))THEN
IF(J1.EQ.274)BBIG(I,10)=0.0
IF(J1.EQ.274)AVMOST(I,II,9)=AVMOST(I,II,9)/30.0
FLOMLL(I,II,10)=FLOMLL(I,II,10)+FLOWL(I)
FLOMLR(I,II,10)=FLOMLR(I,II,10)+FLOWR(I)
FLOWMV(I,II,10)=FLOWMV(I,II,10)+AQV
AVMOST(I,II,10)=AVMOST(I,II,10)+AVTHET
RECHM(I,II,10)=RECHM(I,II,10)+TAVSOR
IF(BIG(I).GT.BBIG(I,10))BBIG(I,10)=BIG(I)
ELSE

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IF((J1.GT.304).AND.(J1.LE.334))THEN
IF(J1.EQ.305)BBIG(I,11)=0.0
IF(J1.EQ.305)AVMOST(I,II,10)=AVMOST(I,II,10)/31.0
FLOMLL(I,II,11)=FLOMLL(I,II,11)+FLOWL(I)
FLOMLR(I,II,11)=FLOMLR(I,II,11)+FLOWR(I)
FLOWMV(I,II,11)=FLOWMV(I,II,11)+AQV
AVMOST(I,II,11)=AVMOST(I,II,11)+AVTHET
RECHM(I,II,11)=RECHM(I,II,11)+TAVSOR
IF(BIG(I).GT.BBIG(I,11))BBIG(I,11)=BIG(I)
ELSE
IF(J1.EQ.335)BBIG(I,12)=0.0
IF(J1.EQ.335)AVMOST(I,II,11)=AVMOST(I,II,11)/30.0
FLOMLL(I,II,12)=FLOMLL(I,II,12)+FLOWL(I)
FLOMLR(I,II,12)=FLOMLR(I,II,12)+FLOWR(I)
FLOWMV(I,II,12)=FLOWMV(I,II,12)+AQV
AVMOST(I,II,12)=AVMOST(I,II,12)+AVTHET
RECHM(I,II,12)=RECHM(I,II,12)+TAVSOR
IF(BIG(I).GT.BBIG(I,12))BBIG(I,12)=BIG(I)
IF(J1.EQ.365)AVMOST(I,II,12)=AVMOST(I,II,12)/31.0
ENDIF
ENDIF
ENDIF
ENDIF
ENDIF
ENDIF
ENDIF
ENDIF
ENDIF
ENDIF
ENDIF
ENDIF
585 CONTINUE
DO 595 I=1,NOBAR
IF((II.EQ.1).AND.(J1.EQ.1))THEN
BIGLAT(I)=FLOWL(I)
SMALAT(I)=FLOWL(I)
BIGVER(I)=AQV
SMAVER(I)=AQV
ELSE
IF(FLOWL(I).GT.BIGLAT(I))BIGLAT(I)=FLOWL(I)
IF(FLOWL(I).LT.SMALAT(I))SMALAT(I)=FLOWL(I)
IF(AQV.GT.BIGVER(I))BIGVER(I)=AQV
IF(AQV.LT.SMAVER(I))SMAVER(I)=AQV
ENDIF
595 CONTINUE
IF(J1.LE.31)THEN
RAINM(II,1)=RAINM(II,1)+RAIN(II,J1)
PERCM(II,1)=PERCM(II,1)+AVIN
RUNM(II,1)=RUNM(II,1)+TAVRUN
EVAM(II,1)=EVAM(II,1)+ACEVAT
CAPINF(II,1)=CAPINF(II,1)+TAVINF

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FLORUN(II,1)=FLORUN(II,1)+FLONET
ELSE
IF((J1.GT.31).AND.(J1.LE.59))THEN
RAINM(II,2)=RAINM(II,2)+RAIN(II,J1)
PERCM(II,2)=PERCM(II,2)+AVIN
RUNM(II,2)=RUNM(II,2)+TAVRUN
EVAM(II,2)=EVAM(II,2)+ACEVAT
CAPINF(II,2)=CAPINF(II,2)+TAVINF
FLORUN(II,2)=FLORUN(II,2)+FLONET
ELSE
IF((J1.GT.59).AND.(J1.LE.90))THEN
RAINM(II,3)=RAINM(II,3)+RAIN(II,J1)
PERCM(II,3)=PERCM(II,3)+AVIN
RUNM(II,3)=RUNM(II,3)+TAVRUN
EVAM(II,3)=EVAM(II,3)+ACEVAT
CAPINF(II,3)=CAPINF(II,3)+TAVINF
FLORUN(II,3)=FLORUN(II,3)+FLONET
ELSE
IF((J1.GT.90).AND.(J1.LE.120))THEN
RAINM(II,4)=RAINM(II,4)+RAIN(II,J1)
PERCM(II,4)=PERCM(II,4)+AVIN
RUNM(II,4)=RUNM(II,4)+TAVRUN
EVAM(II,4)=EVAM(II,4)+ACEVAT
CAPINF(II,4)=CAPINF(II,4)+TAVINF
FLORUN(II,4)=FLORUN(II,4)+FLONET
ELSE
IF((J1.GT.120).AND.(J1.LE.151))THEN
RAINM(II,5)=RAINM(II,5)+RAIN(II,J1)
PERCM(II,5)=PERCM(II,5)+AVIN
RUNM(II,5)=RUNM(II,5)+TAVRUN
EVAM(II,5)=EVAM(II,5)+ACEVAT
CAPINF(II,5)=CAPINF(II,5)+TAVINF
FLORUN(II,5)=FLORUN(II,5)+FLONET
ELSE
IF((J1.GT.151).AND.(J1.LE.181))THEN
RAINM(II,6)=RAINM(II,6)+RAIN(II,J1)
PERCM(II,6)=PERCM(II,6)+AVIN
RUNM(II,6)=RUNM(II,6)+TAVRUN
EVAM(II,6)=EVAM(II,6)+ACEVAT
CAPINF(II,6)=CAPINF(II,6)+TAVINF
FLORUN(II,6)=FLORUN(II,6)+FLONET
ELSE
IF((J1.GT.181).AND.(J1.LE.212))THEN
RAINM(II,7)=RAINM(II,7)+RAIN(II,J1)
PERCM(II,7)=PERCM(II,7)+AVIN
RUNM(II,7)=RUNM(II,7)+TAVRUN
EVAM(II,7)=EVAM(II,7)+ACEVAT
CAPINF(II,7)=CAPINF(II,7)+TAVINF
FLORUN(II,7)=FLORUN(II,7)+FLONET
ELSE

```


END

```

C
C THIS SUBROUTINE COMPUTES AVERAGE DAILY VALUES OF THE
C FLOW PARAMETERS
C
  SUBROUTINE FIND(NYEARS,NMONTH,NOBAR,FLOMLL,FLOMLR,FLOWMV,
+BIGLAT,SMALAT,TDAYS,RAINM,RUNM,EVAM,RECHM,FLORUN,CAPINF,
+ANFLOW,BALFLO)
  IMPLICIT REAL *8(A-H,O-Z)
  COMMON /BIND/ FLOML,TOTML,TOMFLO,BIGEST,SMALST,TOTFL,TOTFV,
+AVDALF,AVMOLF,AVMOVF,AVDAVF,AVDL,AVML,DAIFTO,TOTLL,TOTLR,
+AVDLL,AVDLR,AVDALL,AVDALR,AVMLL,AVMLR,ANRAIN,ANPERC,
+ANRUN,ANEVAT,AVDAIN,AVDAIC,AVDAUN,AVDAET,AVEFCY,ANVPE,ANVERT
  REAL *8 FLOML(1,2,12),TOTML(2,12),TOMFLO(2,12),TOTFL(1),TOTFV(1),A
+VDALF(1),AVMOLF(1),AVMOVF(1),AVDAVF(1),FLOMLL(1,2,12),FLOMLR(1,2,1
+2),FLOWMV(1,2,12),BIGLAT(1),SMALAT(1),TOTLL(1),TOTLR(1),AVDALL(1
+),AVDALR(1),RAINM(2,12),RUNM(2,12),EVAM(2,12),ANRUN(2),ANPERC(2),A
+NEVAT(2),ANRAIN(2),AVDAIN(2),AVDAIC(2),AVDAUN(2),AVDAET(2),VERFLO(
+1,2),AVEFCY(2),PRENT(2),FLTENT(2),RECHM(2,12)
  REAL *8 ANFLOW(2),ANINF(2),FLORUN(2,12),CAPINF(2,12),ANVPE(2),ANVE
+RT(2),BALFLO(2)
  TOTAL=0.0
  TOTALL=0.0
  TOTALR=0.0
  TOTAV=0.0
  DO 600 I=1,NOBAR
  DO 600 J=1,NYEARS
  DO 600 K=1,NMONTH
  FLOML(I,J,K)=(FLOMLL(I,J,K)+FLOMLR(I,J,K))/2.0
600 CONTINUE
  DO 601 I=1,NYEARS
  DO 601 J=1,NMONTH
  IF(RECHM(I,J).LT.0.0)RECHM(I,J)=0.0
601 CONTINUE
  DO 605 I=1,NOBAR
  DO 605 J=1,NYEARS
  DO 605 K=1,NMONTH
  TOTML(J,K)=TOTML(J,K)+FLOML(I,J,K)
605 CONTINUE
  DO 610 I=1,NYEARS
  ANRAIN(I)=0.0
  ANPERC(I)=0.0
  ANRUN(I)=0.0
  ANEVAT(I)=0.0
  ANINF(I)=0.0
  ANFLOW(I)=0.0
  DO 610 J=1,NMONTH
  TOMFLO(I,J)=TOTML(I,J)+FLOWMV(NOBAR,I,J)

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ANRAIN(I)=ANRAIN(I)+RAINM(I,J)
ANRUN(I)=ANRUN(I)+RUNM(I,J)
ANEVAT(I)=ANEVAT(I)+EVAM(I,J)
ANINF(I)=ANINF(I)+CAPINF(I,J)
ANFLOW(I)=ANFLOW(I)+FLORUN(I,J)
610 CONTINUE
DO 611 I=1,NYEARS
AVDAIN(I)=ANRAIN(I)/TDAYS
AVDAUN(I)=ANRUN(I)/TDAYS
AVDAET(I)=ANEVAT(I)/TDAYS
ANPERC(I)=ANRAIN(I)-ANRUN(I)-ANEVAT(I)
AVDAIC(I)=AVDAIN(I)-AVDAUN(I)-AVDAET(I)
611 CONTINUE
DO 615 I=1,NOBAR
BIGEST=BIGEST+BIGLAT(I)
SMALST=SMALST+SMALAT(I)
615 CONTINUE
DO 621 I=1,NOBAR
DO 621 J=1,NYEARS
VERFLO(I,J)=0.0
DO 621 K=1,NMONTH
VERFLO(I,J)=VERFLO(I,J)+FLOWMV(I,J,K)
621 CONTINUE
DO 620 I=1,NOBAR
DO 620 J=1,NYEARS
DO 620 K=1,NMONTH
TOTFL(I)=TOTFL(I)+FLOML(I,J,K)
TOTLL(I)=TOTLL(I)+FLOMLL(I,J,K)
TOTLR(I)=TOTLR(I)+FLOMLR(I,J,K)
TOTFV(I)=TOTFV(I)+FLOWMV(I,J,K)
620 CONTINUE
TOTDAY=NYEARS*TDAYS
TOTMO=NYEARS*NMONTH
DO 625 I=1,NOBAR
AVDALF(I)=TOTFL(I)/TOTDAY
AVDALL(I)=TOTLL(I)/TOTDAY
AVDALR(I)=TOTLR(I)/TOTDAY
AVMOLF(I)=TOTFL(I)/TOTMO
AVMOVF(I)=TOTFV(I)/TOTMO
AVDAVF(I)=TOTFV(I)/TOTDAY
625 CONTINUE
DO 626 I=1,NYEARS
IF(ANRAIN(I).EQ.0.0)THEN
AVEFCY(I)=100.0
PRENT(I)=0.0
ELSE
AVEFCY(I)=100.0-100.0*VERFLO(I,I)/ANRAIN(I)
ENDIF
BALFLO(I)=ANVPE(I)-TOTLL(I)-ANVERT(I)
626 CONTINUE

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```
DO 630 I=1,NOBAR
TOTAL=TOTAL+TOTFL(I)
TOTALL=TOTALL+TOTLL(I)
TOTALR=TOTALR+TOTLR(I)
TOTAV=TOTAV+TOTFV(I)
630 CONTINUE
AVDL=TOTAL/TOTDAY
AVML=TOTAL/TOTMO
AVDLL=TOTALL/TOTDAY
AVDLR=TOTALR/TOTDAY
AVMLL=TOTALL/TOTMO
AVMLR=TOTALR/TOTMO
DAIFTO=AVDL+AVDAVF(NOBAR)
RETURN
END
```

APPENDIX D

**FORTRAN PROGRAM FOR DIRECT GREEN'S FUNCTION
BOUNDARY INTEGRAL SOLUTION**

```

C
C DIRECT GREEN'S FUNCTION BOUNDARY INTEGRAL SOLUTION FOR
C TWO-DIMENSIONAL UNSATURATED MOISTURE-FLOW EQUATION
C
C MAIN PROGRAM
C
  IMPLICIT REAL *8(A-H,O-Z)
  COMMON /COFMN/ COEFF
  PARAMETER(NB=20,NI=5,NT=28,NV=21,ND=25)
  REAL *8 X(NV),Y(NV),BC(NB),COEFF(NB,NB),COLMAT(NB),XIN(NI),YIN(NI),
+ SOLIN(NI),XM(ND),YM(ND),THETAP(ND),DIFFUS(ND),GRAVIT(ND),AREA(NT),
+ DTDX(ND),DTDZ(ND),THETPP(ND)
  INTEGER KODE(ND),IBOUND(NB),NINT(5),NOSING(ND),ISING(ND,9),LINKNO(
+ NT,3)
  EXTERNAL DATA,MATRIX,ELIM,OUTPUT,SOLINT,FLUXIN,CAMPINT,
+ NONLIN,GRAVIN
  OPEN(2,FILE='BEMLAP.DAT',STATUS='OLD')
  OPEN(3,FILE='BEMLAP.OUT',STATUS='UNKNOWN')
  OPEN(4,FILE='INTBOU.OUT',STATUS='UNKNOWN')
  OPEN(5,FILE='INFLUX.OUT',STATUS='UNKNOWN')
  OPEN(6,FILE='INFLUZ.OUT',STATUS='UNKNOWN')
  CALL DATA(NOBEL,NINTER,KODE,X,Y,BC,XIN,YIN,THETAS,SLOPMC,SAIE,HCON
+ S,THETIN,NOEL,DT,TIME,NOSING,ISING,LINKNO)
  NOTEL=NINTER+NOBEL
  I1=0
  I2=0
  DO 10 I=1,NOTEL
  IF((KODE(I).EQ.0).OR.(KODE(I).EQ.1))THEN
  I1=I1+1
  IBOUND(I1)=I
  ELSE
  I2=I2+1
  NINT(I2)=I
  ENDIF
10 CONTINUE
  X(NOBEL+1)=X(1)
  Y(NOBEL+1)=Y(1)
  DO 25 I=1,NOBEL
  N=IBOUND(I)
  XM(N)=(X(I)+X(I+1))/2.0
  YM(N)=(Y(I)+Y(I+1))/2.0
25 CONTINUE
  DO 80 I=1,NOEL
  K=1

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```

K1=LINKNO(I,K)
IF((KODE(K1).EQ.0).OR.(KODE(K1).EQ.1))THEN
XK1=XM(K1)
YK1=YM(K1)
ELSE
XK1=XIN(K1)
YK1=YIN(K1)
ENDIF
K2=LINKNO(I,K+1)
IF((KODE(K2).EQ.0).OR.(KODE(K2).EQ.1))THEN
XK2=XM(K2)
YK2=YM(K2)
ELSE
XK2=XIN(K2)
YK2=YIN(K2)
ENDIF
K3=LINKNO(I,K+2)
IF((KODE(K3).EQ.0).OR.(KODE(K3).EQ.1))THEN
XK3=XM(K3)
YK3=YM(K3)
ELSE
XK3=XIN(K3)
YK3=YIN(K3)
ENDIF
ARM1=XK2*YK3-YK2*XK3
ARM2=-XK1*(YK3-YK2)
ARM3=YK1*(XK3-XK2)
AREA(I)=DABS(ARM1+ARM2+ARM3)/2.0
80 CONTINUE
WRITE(4,47)
47 FORMAT(5X,'TIME',5X,'NODE #1',5X,'NODE #2',5X,'NODE #3',5X,'NODE #
+4',5X,'NODE #5',/)
WRITE(5,47)
WRITE(6,47)
TT=0.0
12 T=0.0
11 T=T+DT
TT=TT+DT
J1=TT/DT
DO 45 I=1,NOBEL
N=IBOUND(I)
IF(KODE(N).EQ.1)BC(I)=0.001
IF(KODE(N).EQ.0)BC(I)=0.3+0.0982*T-0.0265*T**2+0.00107*T**3+0.0000
+595*T**4
IF(J1.EQ.1)THEN
IF(KODE(N).EQ.1)THEN
THETAP(N)=THETIN
THETPP(N)=THETAP(N)
ELSE
THETAP(N)=BC(I)

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```

    THETPP(N)=THETAP(N)
  ENDIF
  ELSE
    IF(KODE(N).EQ.1)THEN
      THETPP(N)=THETAP(N)
      THETAP(N)=COLMAT(I)
    ELSE
      THETPP(N)=THETAP(N)
      THETAP(N)=BC(I)
    ENDIF
  ENDIF
45  CONTINUE
    DO 46 I=1,NINTER
      N=NINT(I)
      IF(J1.EQ.1)THEN
        THETAP(N)=THETIN
        THETPP(N)=THETAP(N)
      ELSE
        THETPP(N)=THETAP(N)
        THETAP(N)=SOLIN(I)
      ENDIF
46  CONTINUE
    DO 40 I=1,NOBEL
      N=IBOUND(I)
      DIFF1=-SLOPMC*HCONS*SAIE/THETAS
      TPP=THETAP(N)
      DIFF2=(TPP/THETAS)**(SLOPMC+2.0)
      DIFFUS(N)=DIFF1*DIFF2
40  CONTINUE
    DO 41 I=1,NINTER
      N=NINT(I)
      DIFF1=-SLOPMC*HCONS*SAIE/THETAS
      TPP=THETAP(N)
      DIFF2=(TPP/THETAS)**(SLOPMC+2.0)
      DIFFUS(N)=DIFF1*DIFF2
41  CONTINUE
    DO 20 I=1,NOBEL
      N=IBOUND(I)
      N=IBOUND(I)
      IF(J1.EQ.1)THEN
        GRAVIT(N)=0.0
      ELSE
        CALL NONLIN(X(I),X(I+1),Y(I),Y(I+1),BC(I),KODE(I),COLMAT(I),SLOPMC
+,HCONS,THETAS,GRAVIT(N))
      ENDIF
20  CONTINUE
    DO 23 I=1,NINTER
      N=NINT(I)
      IF(J1.EQ.1)THEN
        GRAVIT(N)=0.0

```

```

ELSE
CALL GRAVIN(DTDZ(N),SLOPMC,HCONS,THETAS,SOLIN(I),GRAVIT(N))
ENDIF
23 CONTINUE
CALL MATRIX(NOBEL,KODE,X,Y,BC,COEFF,COLMAT,THETAP,XM,YM,IBOUND,DIF
+FUS,DT,AREA,XIN,YIN,LINKNO,NOEL)
IIN=0
DO 21 I=1,NOBEL
N=IBOUND(I)
CALL CAMPINT(I,N,LINKNO,XIN,YIN,BBB,NOEL,XM,YM,AREA,IIN,KODE,DIFFU
+S,DT,GRAVIT,NOSING,ISING)
COLMAT(I)=COLMAT(I)-BBB
21 CONTINUE
CALL ELIM(COLMAT,NOBEL)
CALL SOLINT(NOBEL,NINTER,X,Y,XIN,YIN,BC,COLMAT,SOLIN,KODE,IBOUND,N
+INT,DT,THETAP,DIFFUS,GRAVIT,AREA,NOSING,ISING,NOEL,LINKNO,XM,YM)
CALL FLUXIN(NOBEL,NINTER,X,Y,XIN,YIN,BC,COLMAT,DTDZ,KODE,LINK
+NO,NOEL,AREA,XM,YM,NINT,IBOUND,THETAP,DIFFUS,GRAVIT,DT)
CALL OUTPUT(NOBEL,NINTER,BC,COLMAT,SOLIN,KODE,IBOUND,NINT,TT,DTDZ,
+DTDZ)
IF(T.LT.TIME)GOTO 11
IF(TT.LT.30.0)GOTO 12
STOP
END

```

C

C THIS SUBROUTINE READS DATA FROM THE INPUT FILE

C

```

SUBROUTINE DATA(NOBEL,NINTER,KODE,X,Y,BC,XIN,YIN,THETAS,SLOPMC,SAI
+E,HCONS,THETIN,NOEL,DT,TIME,NOSING,ISING,LINKNO)
IMPLICIT REAL *8(A-H,O-Z)
PARAMETER(NB=20,NI=5,NT=28,NV=21,ND=25)
REAL*8 X(NV),Y(NV),BC(NB),XIN(NI),YIN(NI)
INTEGER KODE(ND),ISING(ND,9),NOSING(ND),LINKNO(NT,3)
READ(2,*)NOBEL,NINTER,NEXEL,NOEL
READ(2,*)THETAS,SLOPMC,SAIE,HCONS,THETIN,DT,TIME
DO 10 I=1,NOBEL
READ(2,*)X(I),Y(I)
10 CONTINUE
DO 20 I=1,NOEL
READ(2,*)(LINKNO(I,K),K=1,3)
20 CONTINUE
NOTEL=NINTER+NEXEL
DO 30 I=1,NINTER
READ(2,*)XIN(I),YIN(I)
30 CONTINUE
DO 40 I=1,NOTEL

```

```

      READ(2,*)KODE(I),NOSING(I)
40  CONTINUE
      DO 45 I=1,NOTEL
      MS=NOSING(I)
      READ(2,*)(ISING(I,K),K=1,MS)
45  CONTINUE
      RETURN
      END

```

```

C
C THIS SUBROUTINE COMPUTES THE GRAVITY TERM AT THE BOUNDARY
C
      SUBROUTINE NONLIN(XX1,XX2,YY1,YY2,BC1,K1,COLMAZ,SLOPMC,HCONS,THETA
+S,GVIT)
      IMPLICIT REAL *8(A-H,O-Z)
      R2=DSQRT((XX2-XX1)**2+(YY2-YY1)**2)
      CO=(XX1-XX2)/R2
      EX1=2.0*SLOPMC+3.0
      EX2=2.0*SLOPMC+2.0
      GRA1=EX1*HCONS/(THETAS**EX1)
      IF(K1.EQ.1)THEN
      GVIT=GRA1*(COLMAZ**EX2)*BC1*CO
      ELSE
      GVIT=GRA1*(BC1**EX2)*COLMAZ*CO
      ENDIF
      RETURN
      END

```

```

C
C THIS SUBROUTINE EVALUATES THE GRAVITY TERM IN THE DOMAIN
C
      SUBROUTINE GRAVIN(DTZZ,SLOPMC,HCONS,THETAS,SOLNN,GVIT)
      IMPLICIT REAL *8(A-H,O-Z)
      EX1=2.0*SLOPMC+3.0
      EX2=2.0*SLOPMC+2.0
      GRA1=EX1*HCONS/(THETAS**EX1)
      GVIT=GRA1*(SOLNN**EX2)*DTZZ
      RETURN
      END

```

```

C
C THIS SUBROUTINE COMPUTES THE ELEMENTS OF THE MATRICES FOR BIM
C
  SUBROUTINE MATRIX(NOBEL,KODE,X,Y,BC,COEFF,COLMAT,THETAP,XM,YM,IB
+OUND,DIFFUS,DT,AREA,XIN,YIN,LINKNO,NOEL)
  IMPLICIT REAL *8(A-H,O-Z)
  PARAMETER(NB=20,NI=5,NT=28,NV=21,ND=25)
  REAL*8 X(NV),Y(NV),BC(NB),COEFF(NB,NB),COLMAT(NB),XM(ND),YM(ND),TH
+ETAP(ND),DIFFUS(ND),AREA(NT),XIN(NI),YIN(NI)
  INTEGER KODE(ND),IBOUND(NB),LINKNO(NT,3)
  EXTERNAL INTEG,SOURCE
  IN=0
  DO 20 I=1,NOBEL
    N=IBOUND(I)
    COLMAT(I)=0.0
    DO 20 J=1,NOBEL
      M=IBOUND(J)
      CALL INTEG(I,J,X(J),X(J+1),Y(J),Y(J+1),XM(N),YM(N),HH,GG,IN,DIFFUS
+(M),DT)
      IF(KODE(N).EQ.1)THEN
        IF(I.EQ.J)THEN
          COEFF(I,J)=0.50
        ELSE
          IF(KODE(M).EQ.1)COEFF(I,J)=HH
          IF(KODE(M).EQ.0)COEFF(I,J)=-GG
        ENDIF
      ELSE
        IF(KODE(M).EQ.1)COEFF(I,J)=HH
        IF(KODE(M).EQ.0)COEFF(I,J)=-GG
      ENDIF
      IF(KODE(M).EQ.1)COLMAT(I)=COLMAT(I)+BC(J)*GG
      IF((KODE(M).EQ.0).AND.(I.NE.J))COLMAT(I)=COLMAT(I)-BC(J)*HH
      IF((KODE(M).EQ.0).AND.(I.EQ.J))COLMAT(I)=COLMAT(I)-BC(J)*0.5
20  CONTINUE
  IIN=0
  DO 30 I=1,NOBEL
    N=IBOUND(I)
    CALL SOURCE(I,N,LINKNO,XIN,YIN,BBB,NOEL,XM,YM,AREA,IIN,KODE,THETAP
+,DIFFUS,DT)
    COLMAT(I)=COLMAT(I)+BBB
30  CONTINUE
  RETURN
  END

```

C
 C THIS SUBROUTINE COMPUTES THE DOMAIN INTEGRAL
 C ASSOCIATED WITH THE GRAVITY TERM
 C

```

SUBROUTINE CAMPINT(I,N,LINKNO,XIN,YIN,BCM,NOEL,XM,YM,AREA,IIN,KODE
+,DIFFUS,DT,GRAVIT,NOSING,ISING)
IMPLICIT REAL *8(A-H,O-Z)
PARAMETER(NB=20,NI=5,NT=28,NV=21,ND=25)
REAL *8 XIN(NI),YIN(NI),AREA(NT),XM(ND),YM(ND),DIFFUS(ND),GRAVIT(ND
+),TC1(3),TC2(3),TC3(3),OMEG(3),XCO(3),YCO(3)
INTEGER LINKNO(NT,3),KODE(ND),NOSING(ND),ISING(ND,9)
EXTERNAL ERROR1
CC=0.5771
OMEG(1)=1.0/3.0
OMEG(2)=1.0/3.0
OMEG(3)=1.0/3.0
TC1(1)=0.0
TC1(2)=0.0
TC1(3)=1.0
TC2(1)=0.0
TC2(2)=1.0
TC2(3)=0.0
TC3(1)=1.0
TC3(2)=0.0
TC3(3)=0.0
BCM=0.0
MS=NOSING(N)
DO 20 J=1,NOEL
K1=LINKNO(J,1)
IF((KODE(K1).EQ.0).OR.(KODE(K1).EQ.1))THEN
XC1=XM(K1)
YC1=YM(K1)
ELSE
XC1=XIN(K1)
YC1=YIN(K1)
ENDIF
K2=LINKNO(J,2)
IF((KODE(K2).EQ.0).OR.(KODE(K2).EQ.1))THEN
XC2=XM(K2)
YC2=YM(K2)
ELSE
XC2=XIN(K2)
YC2=YIN(K2)
ENDIF
K3=LINKNO(J,3)
IF((KODE(K3).EQ.0).OR.(KODE(K3).EQ.1))THEN
XC3=XM(K3)
YC3=YM(K3)
ELSE
XC3=XIN(K3)

```

```

YC3=YIN(K3)
ENDIF
AVDIFF=(DIFFUS(K1)+DIFFUS(K2)+DIFFUS(K3))/3.0
AVGRAV=(GRAVIT(K1)+GRAVIT(K2)+GRAVIT(K3))/3.0
NSING=0
DO 15 KS=1,MS
IF(ISING(N,KS).EQ.J)NSING=J
15  CONTINUE
IF(NSING.EQ.J)THEN
R2=DSQRT((XC2-XC3)**2+(YC2-YC3)**2)
CO=(XC3-XC2)/R2
SI=(YC3-YC2)/R2
SINGI1=0.0
SINGI21=0.0
DO 16 K=1,3
K1=LINKNO(J,K)
XCO(K)=TC1(K)*XC1+TC2(K)*XC2+TC3(K)*XC3
YCO(K)=TC1(K)*YC1+TC2(K)*YC2+TC3(K)*YC3
IF(IIN.EQ.0)THEN
RAD=DSQRT((XM(N)-XCO(K))**2+(YM(N)-YCO(K))**2)
ELSE
RAD=DSQRT((XIN(I)-XCO(K))**2+(YIN(I)-YCO(K))**2)
ENDIF
DENLO=4.0*DIFFUS(K1)*DT
DENO=4.0*3.14*DIFFUS(K1)
SINGI1=SINGI1-AREA(J)*OMEG(K)*GRAVIT(K1)*CC/DENO
SINGI21=SINGI21-AREA(J)*OMEG(K)*GRAVIT(K1)*DLOG(1.0/DENLO)/DENO
16  CONTINUE
AX=(XC2-XC3)/2.0
AY=(YC2-YC3)/2.0
IF(AX.EQ.0.0)THEN
DIST=DABS(XC1-XC3)
ELSE
TA=AY/AX
DIST=DABS(TA*XC1-YC1+YC3-TA*XC3)/DSQRT(TA**2+1.0)
ENDIF
SIG=(XC3-XC1)*(YC2-YC1)-(XC2-XC1)*(YC3-YC1)
IF(SIG.LT.0.0)THEN
DIST=-DIST
ELSE
ENDIF
SIC1=YC1*SI+XC1*CO
ETA1=XC1*SI-YC1*CO
SIC2=YC2*SI+XC2*CO
ETA2=XC2*SI-YC2*CO
SIC3=YC3*SI+XC3*CO
ETA3=XC3*SI-YC3*CO
ALPHA=(SIC3-SIC1)/(ETA3-ETA1)
BETA=(SIC2-SIC1)/(ETA2-ETA1)
BRAC1=(ALPHA-BETA)*(DLOG(DIST)-1.5)

```

```

BRAC2=ALPHA*(DLOG(1.0+ALPHA**2))/2.0
BRAC3=BETA*(DLOG(1.0+BETA**2))/2.0
BRAC4=ATAN(ALPHA)-ATAN(BETA)
PRO=2.0*AVGRAV/(4.0*AVDIFF*3.14)
SINGI22=-PRO*(BRAC1+BRAC2-BRAC3+BRAC4)*(DIST**2)/2.0
BCM=BCM+SINGIN1+SINGI21+SINGI22
ELSE
BCM1=0.0
DO 21 K=1,3
K1=LINKNO(J,K)
XCO(K)=TC1(K)*XC1+TC2(K)*XC2+TC3(K)*XC3
YCO(K)=TC1(K)*YC1+TC2(K)*YC2+TC3(K)*YC3
IF(IIN.EQ.0)THEN
RAD=DSQRT((XM(N)-XCO(K))**2+(YM(N)-YCO(K))**2)
ELSE
RAD=DSQRT((XIN(I)-XCO(K))**2+(YIN(I)-YCO(K))**2)
ENDIF
A=RAD**2/(4.0*DIFFUS(K1)*DT)
DENO=4.0*3.14*DIFFUS(K1)
BCM1=BCM1+OMEG(K)*GRAVIT(K1)*ERROR1(A)/DENO
21 CONTINUE
BCM=BCM+AREA(J)*BCM1
ENDIF
20 CONTINUE
RETURN
END

```

```

C
C THIS SUBROUTINE SOLVES SIMULTANEOUS EQUATIONS BY
C GAUSS-JORDAN ELIMINATION METHOD
C

```

```

SUBROUTINE ELIM(COLMAT,NOBEL)
IMPLICIT REAL *8(A-H,O-Z)
COMMON /COFMN/ COEFF
PARAMETER (NB=20,NT=28)
REAL*8 COEFF(NB,NB),PP(NT),COLMAT(NB)
DO 947 I=1,NOBEL
IF(I.EQ.NOBEL)GOTO 956
BIG=DABS(COEFF(I,I))
II=I+1
DO 951 L=II,NOBEL
IF(DABS(COEFF(L,I)).GT.BIG)BIG=DABS(COEFF(L,I))
951 CONTINUE
DO 952 L=I,NOBEL
IF(DABS(COEFF(L,I)).EQ.BIG)NR=L
952 CONTINUE
DO 953 M=1,NOBEL

```

```

      PP(M)=COEFF(I,M)
953  CONTINUE
      NON1=NOBEL+1
      PP(NON1)=COLMAT(I)
      DO 954 M=1,NOBEL
      COEFF(I,M)=COEFF(NR,M)
954  CONTINUE
      COLMAT(I)=COLMAT(NR)
      DO 955 M=1,NOBEL
      COEFF(NR,M)=PP(M)
955  CONTINUE
      COLMAT(NR)=PP(NON1)
956  DEN=COEFF(I,I)
      DO 948 J=1,NOBEL
      COEFF(I,J)=COEFF(I,J)/DEN
948  CONTINUE
      COLMAT(I)=COLMAT(I)/DEN
      DO 949 K=1,NOBEL
      IF(K.EQ.I)GOTO 949
      IF(COEFF(K,I).EQ.0.0)GOTO 949
      CON=COEFF(K,I)
      DO 950 J=1,NOBEL
      IF(DABS(COEFF(I,J)).LT.1.0E-30)THEN
      COEFF(I,J)=1.0E-30
      ELSE
      COEFF(K,J)=COEFF(K,J)-COEFF(I,J)*CON
      ENDIF
950  CONTINUE
      COLMAT(K)=COLMAT(K)-COLMAT(I)*CON
949  CONTINUE
947  CONTINUE
      RETURN
      END

```

C

C THIS SUBROUTINE COMPUTES THE MOISTURE CONTENTS IN THE DOMAIN

C

```

      SUBROUTINE SOLINT(NOBEL,NINTER,X,Y,XIN,YIN,BC,COLMAT,SOLIN,KODE,IB
+OUND,NINT,DT,THETAP,DIFFUS,GRAVIT,AREA,NOSING,ISING,NOEL,LINKNO,XM
+,YM)
      IMPLICIT REAL *8(A-H,O-Z)
      PARAMETER(NB=20,NI=5,NT=28,NV=21,ND=25)
      REAL *8 X(NV),Y(NV),XIN(NI),YIN(NI),SOLIN(NI),BC(NB),COLMAT(NB),THE
+TAP(ND),DIFFUS(ND),GRAVIT(ND),AREA(NT),XM(NB),YM(NB)
      INTEGER KODE(ND),IBOUND(NB),NINT(5),NOSING(ND),ISING(ND,9),LINKNO(
+NT,3)
      EXTERNAL INTEG,CAMPINT,SOURCE

```

```

IN=1
IIN=1
DO 10 I=1,NINTER
N=NINT(I)
CALL SOURCE(I,N,LINKNO,XIN,YIN,BCB,NOEL,XM,YM,AREA,IIN,KODE,THETAP
+,DIFFUS,DT)
SOLIN(I)=0.0
DO 11 J=1,NOBEL
M=IBOUND(J)
CALL INTEG(I,J,X(J),X(J+1),Y(J),Y(J+1),XIN(I),YIN(I),AA,BB,IN,DIFF
+US(M),DT)
IF(KODE(M).EQ.1)SOLIN(I)=SOLIN(I)+BC(J)*BB-COLMAT(J)*AA
IF(KODE(M).EQ.0)SOLIN(I)=SOLIN(I)+COLMAT(J)*BB-BC(J)*AA
11 CONTINUE
SOLIN(I)=SOLIN(I)+BCB
10 CONTINUE
DO 20 I=1,NINTER
N=NINT(I)
IIN=1
CALL CAMPINT(I,N,LINKNO,XIN,YIN,BBB,NOEL,XM,YM,AREA,IIN,KODE,DIFFU
+S,DT,GRAVIT,NOSING,ISING)
SOLIN(I)=SOLIN(I)-BBB
20 CONTINUE
RETURN
END

```

C

C THIS SUBROUTINE COMPUTES THE MOISTURE-FLUX IN THE DOMAIN

C

```

SUBROUTINEFLUXIN(NOBEL,NINTER,X,Y,XIN,YIN,BC,COLMAT,DTDZ,DTDZ,KOD
+E,LINKNO,NOEL,AREA,XM,YM,NINT,IBOUND,THETAP,DIFFUS,GRAVIT,DT)
IMPLICIT REAL *8(A-H,O-Z)
PARAMETER(NB=20,NI=5,NT=28,NV=21,ND=25)
REAL*8 X(NV),Y(NV),XIN(NI),YIN(NI),DTDZ(ND),DTDZ(ND),BC(NB),COLMAT
+(NB),AREA(NT),XM(NB),YM(NB),THETAP(ND),DIFFUS(ND),GRAVIT(ND)
INTEGER KODE(ND),LINKNO(NT,3),NINT(5),IBOUND(NB)
EXTERNAL FINTEG,FSOURC
IN=1
IIN=1
DO 10 I=1,NINTER
N=NINT(I)
ID=0
CALL FSOURC(I,LINKNO,XM,YM,XIN,YIN,BCCX,BCCZ,NOEL,AREA,KODE,THETAP
+,DIFFUS,ID,DT)
ID=1
CALL FSOURC(I,LINKNO,XM,YM,XIN,YIN,BCGX,BCGZ,NOEL,AREA,KODE,GRAVIT
+,DIFFUS,ID,DT)

```

```

DTDX(N)=0.0
DTDZ(N)=0.0
DO 11 J=1,NOBEL
M=IBOUND(J)
CALL FINTEG(X(J),X(J+1),Y(J),Y(J+1),XIN(I),YIN(I),AF,BF,AFZ,BFZ,DI
+FFUS(M),DT)
IF(KODE(M).EQ.1) DTDX(N)=DTDX(N)+BC(J)*BF-COLMAT(J)*AF
IF(KODE(M).EQ.1) DTDZ(N)=DTDZ(N)+BC(J)*BFZ-COLMAT(J)*AFZ
IF(KODE(M).EQ.0) DTDX(N)= DTDX(N)+COLMAT(J)*BF-BC(J)*AF
IF(KODE(M).EQ.0) DTDZ(N)=DTDZ(N)+COLMAT(J)*BFZ-BC(J)*AFZ
11 CONTINUE
DTDX(N)=DTDX(N)+BCCX-BCGX
DTDZ(N)=DTDZ(N)+BCCZ-BCGZ
10 CONTINUE
RETURN
END

```

```

C
C THIS FUNCTION COMPUTES THE EXPONENTIAL INTEGRAL
C

```

```

REAL*8 FUNCTION ERROR1(Z)
IMPLICIT REAL *8(A-H,O-Z)
REAL*8 Z
IF(Z.LE.1.0)THEN
A0=-0.57721566
A1=0.99999193
A2=-0.24991055
A3=0.05519968
A4=-0.00976004
A5=0.00107857
ERROR1=-DLOG(Z)+A0+A1*Z+A2*Z**2+A3*Z**3+A4*Z**4+A5*Z**5
ELSE
A1=8.5733287401
A2=18.059016973
A3=8.6347608925
A4=0.2677737343
B1=9.5733223454
B2=25.6329561486
B3=21.0996530827
B4=3.9584969228
XNUM=Z**4+A1*Z**3+A2*Z**2+A3*Z+A4
XEN1=DEXP(-Z)
DEN2=Z**4+B1*Z**3+B2*Z**2+B3*Z+B4
DENO=Z*DEN2
ERROR1=XEN1*XNUM/DENO
ENDIF
RETURN

```

END

C
 C THIS FUNCTION COMPUTES THE SERIES TO EVALUATE THE
 C SINGULAR INTEGRAL AT THE BOUNDARY

C
 REAL*8 FUNCTION ERROR2(Z)
 IMPLICIT REAL *8(A-H,O-Z)
 REAL*8 Z
 C=0.5771
 D=DLOG(Z)
 N=0
 CC=0.0
 SUM=0.0
 DD=1.0
 10 N=N+1
 CC=CC+1.0
 DD=DD*CC
 DENO=N*(2.0*N+1.0)*DD
 PROD=((-1)**(N-1))*(Z**N)/DENO
 SUM=SUM+PROD
 IF(DABS(PROD).LT.1.0E-9)THEN
 GOTO 11
 ELSE
 GOTO 10
 ENDIF
 11 ERROR2=-C-D+2.0+SUM
 RETURN
 END

C
 C THIS SUBROUTINE EVALUATES THE BOUNDARY INTEGRAL
 C

SUBROUTINE INTEG(I,J,XX1,XX2,YY1,YY2,XP,YP,HH,GG,IN,DIFFS1,DT)
 IMPLICIT REAL *8(A-H,O-Z)
 REAL*8 SCIE(4),OME(4),XCO(4),YCO(4)
 EXTERNAL ERROR1,ERROR2
 SCIE(1)=0.86113631
 SCIE(2)=-SCIE(1)
 SCIE(3)=0.33998104
 SCIE(4)=-SCIE(3)
 OME(1)=0.34785485
 OME(2)=OME(1)
 OME(3)=0.65214515

```

OME(4)=OME(3)
AX=(XX2-XX1)/2.0
BX=(XX2+XX1)/2.0
AY=(YY2-YY1)/2.0
BY=(YY2+YY1)/2.0
IF(((I-J).EQ.0).AND.(IN.EQ.0))THEN
DELS=2.0*DSQRT(AX**2+AY**2)
ALPHA=DELS**2/(16.0*DIFFS1*DT)
GG=DELS*ERROR2(ALPHA)/(4.0*3.14)
HH=0.0
ELSE
IF(AX.EQ.0.0)THEN
DIST=DABS(XP-XX1)
ELSE
TA=AY/AX
DIST=DABS((TA*XP-YP+YY1-TA*XX1)/DSQRT(TA**2+1))
ENDIF
SIG=(XX1-XP)*(YY2-YP)-(XX2-XP)*(YY1-YP)
IF(SIG.LT.0.0)THEN
DIST=-DIST
ELSE
GOTO 32
ENDIF
32  GG=0.0
    HH=0.0
    AXY=DSQRT(AX**2+AY**2)
    DO 15 K=1,4
    XCO(K)=AX*SCIE(K)+BX
    YCO(K)=AY*SCIE(K)+BY
    RAD=DSQRT((XP-XCO(K))**2+(YP-YCO(K))**2)
    A=(RAD**2)/(4.0*DIFFS1*DT)
    BRAC=44.0*RAD**2
    GG=GG+7.0*ERROR1(A)*OME(K)*AXY/88.0
    HH=HH-7.0*(DEXP(-A)*DIST*OME(K)*AXY)/BRAC
15  CONTINUE
    ENDIF
    RETURN
    END

```

C
C THIS SUBROUTINE EVALUATES THE BOUNDARY INTEGRAL
C FOR MOISTURE-FLUX

C
SUBROUTINE FINTEG(XX1,XX2,YY1,YY2,XP,YP,HHX,GGX,HHZ,GGZ,DIFFS1,DT)
IMPLICIT REAL *8(A-H,O-Z)
REAL *8 SCIE(4),OME(4),XCO(4),YCO(4)
SCIE(1)=0.86113631

```

SCIE(2)=-SCIE(1)
SCIE(3)=0.33998104
SCIE(4)=-SCIE(3)
OME(1)=0.34785485
OME(2)=OME(1)
OME(3)=0.65214515
OME(4)=OME(3)
AX=(XX2-XX1)/2.0
BX=(XX2+XX1)/2.0
AY=(YY2-YY1)/2.0
BY=(YY2+YY1)/2.0
R2=DSQRT((XX2-XX1)**2+(YY2-YY1)**2)
CO=(XX1-XX2)/R2
SI=(YY1-YY2)/R2
IF(AX.EQ.0.0)THEN
DIST=DABS(XP-XX1)
ELSE
TA=AY/AX
DIST=DABS((TA*XP-YP+YY1-TA*XX1)/SQRT(TA**2+1))
ENDIF
SIG=(XX1-XP)*(YY2-YP)-(XX2-XP)*(YY1-YP)
IF(SIG.LT.0.0)THEN
DIST=-DIST
ELSE
GOTO 32
ENDIF
32  GGX=0.0
    GGZ=0.0
    HHX=0.0
    HHZ=0.0
    DO 15 K=1,4
    XCO(K)=AX*SCIE(K)+BX
    YCO(K)=AY*SCIE(K)+BY
    XXX=-XCO(K)+XP
    ZZZ=-YCO(K)+YP
    RAD=DSQRT((XP-XCO(K))**2+(YP-YCO(K))**2)
    SEE=-SI*(YCO(K)-YP)-CO*(XCO(K)-XP)
34  DRDX=-(-SEE*CO-DIST*SI)/RAD
    DRDZ=-(DIST*CO-SEE*SI)/RAD
    A=RAD**2/(4.0*DIFFS1*DT)
    GGX=GGX-7.0*(DRDX/RAD)*DEXP(-A)*OME(K)*DSQRT(AX**2+AY**2)/44.0
    GGZ=GGZ-7.0*(DRDZ/RAD)*DEXP(-A)*OME(K)*DSQRT(AX**2+AY**2)/44.0
    BHX1=DIST*DRDX/((RAD**3)*3.14)
    BHX2=DEXP(-A)*(1.0+A)
    BHX3=BHX1*BHX2
    BHX4=SI*DEXP(-A)/((RAD**2)*6.28)
    BHX=BHX3-BHX4
    HHX=HHX+BHX*OME(K)*DSQRT(AX**2+AY**2)
    BHZ1=DIST*DRDZ/((RAD**3)*3.14)
    BHZ2=DEXP(-A)*(1.0+A)

```

```

    BHZ3=BHZ1*BHZ2
    BHZ4=-CO*DEXP(-A)/((RAD**2)*6.28)
    BHZ=BHZ3-BHZ4
    HHZ=HHZ+BHZ*OME(K)*DSQRT(AX**2+AY**2)
15  CONTINUE
    RETURN
    END

```

C

C THIS SUBROUTINE EVALUATES THE DOMAIN INTEGRAL

C ASSOCIATED WITH THE INITIAL CONDITION

C

```

    SUBROUTINE SOURCE(I,N,LINKNO,XIN,YIN,BCM,NOEL,XM,YM,AREA,IIN,KODE,
+THETAP,DIFFUS,DT)
    IMPLICIT REAL *8(A-H,O-Z)
    PARAMETER(NB=20,NI=5,NT=28,ND=25)
    REAL*8 XIN(NI),YIN(NI),AREA(NT),OMEG(3),XCO(4),YCO(4),TC1(3),TC2(3
+),TC3(3),XM(ND),YM(ND),DIFFUS(ND),THETAP(ND)
    INTEGER LINKNO(NT,3),KODE(ND)
    OMEG(1)=1.0/3.0
    OMEG(2)=1.0/3.0
    OMEG(3)=1.0/3.0
    TC1(1)=0.5
    TC1(2)=0.5
    TC1(3)=0.0
    TC2(1)=0.0
    TC2(2)=0.5
    TC2(3)=0.5
    TC3(1)=0.5
    TC3(2)=0.0
    TC3(3)=0.5
    BCM=0.0
    DO 20 J=1,NOEL
    K1=LINKNO(J,1)
    IF((KODE(K1).EQ.0).OR.(KODE(K1).EQ.1))THEN
    XC1=XM(K1)
    YC1=YM(K1)
    ELSE
    XC1=XIN(K1)
    YC1=YIN(K1)
    ENDIF
    K2=LINKNO(J,2)
    IF((KODE(K2).EQ.0).OR.(KODE(K2).EQ.1))THEN
    XC2=XM(K2)
    YC2=YM(K2)
    ELSE
    XC2=XIN(K2)

```

```

YC2=YIN(K2)
ENDIF
K3=LINKNO(J,3)
IF((KODE(K3).EQ.0).OR.(KODE(K3).EQ.1))THEN
XC3=XM(K3)
YC3=YM(K3)
ELSE
XC3=XIN(K3)
YC3=YIN(K3)
ENDIF
BCM1=0.0
DO 21 K=1,3
K1=LINKNO(J,K)
IF(K.NE.3)THEN
K2=LINKNO(J,K+1)
ELSE
K2=LINKNO(J,1)
ENDIF
XCO(K)=TC1(K)*XC1+TC2(K)*XC2+TC3(K)*XC3
YCO(K)=TC1(K)*YC1+TC2(K)*YC2+TC3(K)*YC3
IF(IIN.EQ.0)THEN
RAD=DSQRT((XM(N)-XCO(K))**2+(YM(N)-YCO(K))**2)
ELSE
RAD=DSQRT((XIN(I)-XCO(K))**2+(YIN(I)-YCO(K))**2)
ENDIF
AVDIFF=(DIFFUS(K1)+DIFFUS(K2))/2.0
AVTHET=(THETAP(K1)+THETAP(K2))/2.0
A=RAD**2/(4.0*AVDIFF*DT)
DENO=4.0*3.14*AVDIFF*DT
BCM1=BCM1+OMEG(K)*AVTHET*DEXP(-A)/DENO
21 CONTINUE
BCM=BCM+AREA(J)*BCM1
20 CONTINUE
RETURN
END

```

C

C THIS SUBROUTINE COMPUTES THE DOMAIN INTEGRAL FOR MOISTURE-FLUX

C

```

SUBROUTINE FSOURC(I,LINKNO,XM,YM,XIN,YIN,BCMX,BCMZ,NOEL,AREA,KODE,
+THEGRA,DIFFUS,ID,DT)
IMPLICIT REAL*8(A-H,O-Z)
PARAMETER(NB=20,NI=5,ND=25,NT=28)
REAL*8 XIN(NI),YIN(NI),AREA(NT),XM(ND),YM(ND),THEGRA(ND),DIFFUS(ND
+)
INTEGER LINKNO(NT,3),KODE(ND)
TC1=1.0/3.0

```

```

TC2=1.0/3.0
TC3=1.0/3.0
BCMX=0.0
BCMZ=0.0
DO 20 J=1,NOEL
K1=LINKNO(J,1)
IF((KODE(K1).EQ.0).OR.(KODE(K1).EQ.1))THEN
XC1=XM(K1)
YC1=YM(K1)
ELSE
XC1=XIN(K1)
YC1=YIN(K1)
ENDIF
K2=LINKNO(J,2)
IF((KODE(K2).EQ.0).OR.(KODE(K2).EQ.1))THEN
XC2=XM(K2)
YC2=YM(K2)
ELSE
XC2=XIN(K2)
YC2=YIN(K2)
ENDIF
K3=LINKNO(J,3)
IF((KODE(K3).EQ.0).OR.(KODE(K3).EQ.1))THEN
XC3=XM(K3)
YC3=YM(K3)
ELSE
XC3=XIN(K3)
YC3=YIN(K3)
ENDIF
AVDIFG=(THEGRA(K1)+THEGRA(K2)+THEGRA(K3))/3.0
AVDIFF=(DIFFUS(K1)+DIFFUS(K2)+DIFFUS(K3))/3.0
XCO=TC1*XC1+TC2*XC2+TC3*XC3
YCO=TC1*YC1+TC2*YC2+TC3*YC3
XXX=-(XCO-XIN(I))
ZZZ=-(YCO-YIN(I))
RAD=DSQRT((XIN(I)-XCO)**2+(YIN(I)-YCO)**2)
DRDX=XXX/RAD
DRDZ=ZZZ/RAD
A=RAD**2/(4.0*AVDIFF*DT)
IF(ID.EQ.0)THEN
BCMX1=-AVDIFG*DEXP(-A)*XXX/(8.0*3.14*(AVDIFF**2)*(DT**2))
BCMZ1=-AVDIFG*DEXP(-A)*ZZZ/(8.0*3.14*(AVDIFF**2)*(DT**2))
ELSE
BCMX1=-7.0*AVDIFG*DEXP(-A)*XXX/((RAD**2)*AVDIFF*44.0)
BCMZ1=-7.0*AVDIFG*DEXP(-A)*ZZZ/((RAD**2)*AVDIFF*44.0)
ENDIF
BCMX=BCMX+AREA(J)*BCMX1
BCMZ=BCMZ+AREA(J)*BCMZ1
20 CONTINUE
RETURN

```

END

C

C THIS SUBROUTINE WRITES THE VALUES IN THE OUTPUT FILES

C

```

SUBROUTINEOUTPUT(NOBEL,NINTER,BC,COLMAT,SOLIN,KODE,IBOUND,NINT,TT
+,DTDX,DTDZ)
IMPLICIT REAL *8(A-H,O-Z)
PARAMETER(NB=20,NI=5,ND=25)
REAL*8 BC(NB),COLMAT(NB),SOLIN(ND),DTDX(ND),DTDZ(ND)
INTEGER KODE(ND),IBOUND(NB),NINT(NI)
WRITE(3,37)TT
37 FORMAT(1X,'TIME = ',F5.2)
WRITE(3,36)
36 FORMAT(5X,'BOUNDARY ELEMENT SOLUTION OF POISSON'S EQUATION
+ USING',
+/,12X,'BOUNDARY INTEGRATION FOR THE SOURCE TERM',/,5X,' _____
+ _____',/,,,/)
WRITE(3,35)
35 FORMAT(5X,'TABLE : SOLUTION AT THE BOUNDARY ',/,5X,' _____
+ _____',/)
WRITE(3,15)
15 FORMAT(5X,'NODE NO.',5X,'POTENTIAL',10X,'FLUX',/,5X,' _____',5X
+,' _____',5X,' _____',/)
DO 20 I=1,NOBEL
N=IBOUND(I)
IF(KODE(N).EQ.0)WRITE(3,25)N,BC(I),COLMAT(I)
25 FORMAT(7X,I2,7X,F10.5,6X,F10.5)
IF(KODE(N).EQ.1)WRITE(3,25)N,COLMAT(I),BC(I)
20 CONTINUE
WRITE(3,21)
21 FORMAT(5X,' _____',/,,,/)
WRITE(3,45)
45 FORMAT(5X,'TABLE : SOLUTION IN THE DOMAIN',/,5X,' _____
+ _____')
WRITE(3,40)
40 FORMAT(/,5X,'NODE NO.',5X,'POTENTIAL',5X,'FLUX IN X-DIR.',4X,'FLUX
+ IN Z-DIR.',/,5X,' _____',5X,' _____',5X,' _____',4X,
+ ' _____',/)
DO 30 I=1,NINTER
N=NINT(I)
WRITE(3,50)N,SOLIN(I)
50 FORMAT(8X,I2,6X,F10.5)
30 CONTINUE
WRITE(3,31)
31 FORMAT(5X,' _____
+ _____',/,,,/)

```

```
WRITE(4,70)TT,SOLIN(1),SOLIN(2),SOLIN(3),SOLIN(4),SOLIN(5)
70  FORMAT(6(3X,F10.5))
    IF(TT.EQ.1.0)THEN
      WRITE(5,73)
73  FORMAT(4X,'NODE NO.',4X,'FLUX IN X-DIR',4X,'FLUX IN Z-DIR')
      DO 71 I=1,NINTER
        WRITE(5,72)I,DTDZ(I),DTDZ(I)
72  FORMAT(5X,I2,5X,F10.5,5X,F10.5)
71  CONTINUE
    ELSE
      ENDIF
      RETURN
    END
```

APPENDIX E

**FORTRAN PROGRAM FOR PERTURBATION GREEN'S
FUNCTION BOUNDARY INTEGRAL SOLUTION**

```

C
C PERTURBATION GREEN'S FUNCTION BOUNDARY INTEGRAL SOLUTION FOR
C TWO-DIMENSIONAL UNSATURATED MOISTURE-FLOW EQUATION
C
C           MAIN PROGRAM
C
  IMPLICIT REAL *8(A-H,O-Z)
  COMMON /COFMN/ COEFF
  PARAMETER(NB=20,NI=5,NT=28,NV=21,ND=25)
  REAL*8 X(NV),Y(NV),BC(NB),COEFF(NB,NB),COLMAT(NB),XIN(NI),YIN(NI),
+ SOLIN(NI),XM(ND),YM(ND),THETAP(ND),DIFFUS(ND),THETPP(ND),RHMAT(NB)
+ ,GRAVIT(ND),AREA(NT),DTDZ(ND),THETAF(ND),BCC(NB)
  INTEGER KODE(ND),IBOUND(NB),NINT(5),NOSING(ND),ISING(ND,9),LINKNO(
+ NT,3)
  EXTERNAL DATA,MATRIX,ELIM,OUTPUT,SOLINT,FLUXIN,CAMPINT,NONLIN,
+ GRAVIN
  OPEN(2,FILE='BEMLAP.DAT',STATUS='OLD')
  OPEN(3,FILE='BEMLAP.OUT',STATUS='UNKNOWN')
  OPEN(4,FILE='INTBOU.OUT',STATUS='UNKNOWN')
  CALL DATA(NOBEL,NINTER,KODE,X,Y,BCC,XIN,YIN,THETAS,SLOPMC,SAIE,HCO
+ NS,THETIN,NOEL,DT,TIME,NOSING,ISING,LINKNO)
  NOTEL=NINTER+NOBEL
  I1=0
  I2=0
  DO 10 I=1,NOTEL
  IF((KODE(I).EQ.0).OR.(KODE(I).EQ.1))THEN
  I1=I1+1
  IBOUND(I1)=I
  ELSE
  I2=I2+1
  NINT(I2)=I
  ENDIF
10 CONTINUE
  WRITE(4,47)
47  FORMAT(5X,'TIME',5X,'NODE #1',5X,'NODE #2',5X,'NODE #3',5X,'NODE #
+ 4',5X,'NODE #5',/)
  TT=0.0
12  T=0.0
11  T=T+DT
  TT=TT+DT
  WRITE(*,*)'TT=',TT
  J1=TT/DT
  DO 45 I=1,NOBEL

```

```

N=IBOUND(I)
IF(KODE(N).EQ.1)BCC(I)=0.001
IF(KODE(N).EQ.0)BCC(I)=0.3+0.0982*T-0.0265*T**2+0.00107*T**3+0.000
+0595*T**4
IF(J1.EQ.1)THEN
IF(KODE(N).EQ.1)THEN
THETAP(N)=THETIN
THETPP(N)=THETAP(N)
ELSE
THETAP(N)=BCC(I)
THETPP(N)=THETAP(N)
ENDIF
ELSE
IF(KODE(N).EQ.1)THEN
THETPP(N)=THETAP(N)
THETAP(N)=THETAF(N)
ELSE
THETPP(N)=THETAP(N)
THETAP(N)=BCC(I)
ENDIF
ENDIF
45 CONTINUE
DO 49 I=1,NINTER
N=NINT(I)
IF(J1.EQ.1)THEN
THETAP(N)=THETIN
THETPP(N)=THETAP(N)
ELSE
THETPP(N)=THETAP(N)
THETAP(N)=THETAF(N)
ENDIF
49 CONTINUE
DO 40 I=1,NOBEL
N=IBOUND(I)
DIFF1=-SLOPMC*HCONS*SAIE/THETAS
TPP=THETAP(N)
DIFF2=(TPP/THETAS)**(SLOPMC+2.0)
DIFFUS(N)=DIFF1*DIFF2
40 CONTINUE
DO 41 I=1,NINTER
N=NINT(I)
DIFF1=-SLOPMC*HCONS*SAIE/THETAS
TPP=THETAP(N)
DIFF2=(TPP/THETAS)**(SLOPMC+2.0)
DIFFUS(N)=DIFF1*DIFF2
41 CONTINUE
DO 46 I=1,NOBEL
N=IBOUND(I)
THETAF(N)=0.0
46 CONTINUE

```

```

DO 48 I=1,NINTER
N=NINT(I)
THETAF(N)=0.0
48 CONTINUE
II=0
22 II=II+1
DO 31 I=1,NOBEL
N=IBOUND(I)
IF(II.EQ.1)THEN
BC(I)=BCC(I)
ELSE
BC(I)=0.0
THETAP(N)=0.0
ENDIF
31 CONTINUE
DO 27 I=1,NINTER
N=NINT(I)
IF(II.GT.1)THETAP(N)=0.0
27 CONTINUE
CALL MATRIX(NOBEL,KODE,X,Y,BC,COEFF,RHMAT,THETAP,XM,YM,IBOUND,DIFF
+US,DT,AREA,XIN,YIN,LINKNO,NOEL)
DO 20 I=1,NOBEL
N=IBOUND(I)
N=IBOUND(I)
IF((J1.EQ.1).AND.(II.EQ.1))THEN
GRAVIT(N)=0.0
ELSE
CALL NONLIN(X(I),X(I+1),Y(I),Y(I+1),BC(I),KODE(I),COLMAT(I),SLOPMC
+,HCONS,THETAS,GRAVIT(N))
ENDIF
20 CONTINUE
DO 23 I=1,NINTER
N=NINT(I)
IF((J1.EQ.1).AND.(II.EQ.1))THEN
GRAVIT(N)=0.0
ELSE
CALL GRAVIN(DTDZ(N),SLOPMC,HCONS,THETAS,SOLIN(I),GRAVIT(N))
ENDIF
23 CONTINUE
IIN=0
DO 21 I=1,NOBEL
N=IBOUND(I)
CALL CAMPINT(I,N,LINKNO,XIN,YIN,BBB,NOEL,XM,YM,AREA,IIN,KODE,DIFFU
+S,DT,GRAVIT,NOSING,ISING)
COLMAT(I)=RHMAT(I)-BBB
21 CONTINUE
CALL ELIM(COLMAT,NOBEL)
CALL SOLINT(NOBEL,NINTER,X,Y,XIN,YIN,BC,COLMAT,SOLIN,KODE,IBOUND,N
+INT,DT,THETAP,DIFFUS,GRAVIT,AREA,NOSING,ISING,NOEL,LINKNO,XM,YM)
CALL FLUXIN(NOBEL,NINTER,X,Y,XIN,YIN,BC,COLMAT,DTDZ,DTDX,DTDZ,KODE,LINK

```

```

+NO,NOEL,AREA,XM,YM,NINT,IBOUND,THETAP,DIFFUS,GRAVIT,DT)
DO 24 I=1,NINTER
N=NINT(I)
THETAF(N)=THETAF(N)+SOLIN(I)
24 CONTINUE
DO 26 I=1,NOBEL
N=IBOUND(I)
IF(KODE(N).EQ.1)THEN
THETAF(N)=THETAF(N)+COLMAT(I)
ELSE
THETAF(N)=THETAF(N)+BC(I)
ENDIF
26 CONTINUE
TSLN=0.0
DO 28 I=1,NINTER
TSLN=TSLN+SOLIN(I)
28 CONTINUE
DO 29 I=1,NOBEL
N=IBOUND(I)
IF(KODE(N).EQ.1)TSLN=TSLN+COLMAT(I)
IF(KODE(N).EQ.0)TSLN=TSLN+BC(I)
29 CONTINUE
IF(DABS(TSLN).GT.1.0E-5)GOTO 22
CALL OUTPUT(NOBEL,NINTER,BC,COLMAT,THETAF,KODE,IBOUND,NINT,TT)
IF(T.LT.TIME)GOTO 11
IF(TT.LT.30.0)GOTO 12
STOP
END

```

C

C THIS SUBROUTINE COMPUTES THE MATRICES IN THE PGF SOLUTION

C

```

SUBROUTINE MATRIX(NOBEL,KODE,X,Y,BC,COEFF,RHMAT,THETAP,XM,YM,
+IBOUND,DIFFUS,DT,AREA,XIN,YIN,LINKNO,NOEL)
IMPLICIT REAL *8(A-H,O-Z)
PARAMETER(NB=20,NI=5,NT=28,NV=21,ND=25)
REAL *8 X(NV),Y(NV),BC(NB),COEFF(NB,NB),RHMAT(NB),XM(ND),YM(ND),THE
+TAP(ND),DIFFUS(ND),AREA(NT),XIN(NI),YIN(NI)
INTEGER KODE(ND),IBOUND(NB),LINKNO(NT,3)
EXTERNAL INTEG,SOURCE
X(NOBEL+1)=X(1)
Y(NOBEL+1)=Y(1)
DO 25 I=1,NOBEL
N=IBOUND(I)
XM(N)=(X(I)+X(I+1))/2.0
YM(N)=(Y(I)+Y(I+1))/2.0

```

```

25 CONTINUE
  IN=0
  DO 20 I=1,NOBEL
    N=IBOUND(I)
    RHMAT(I)=0.0
    DO 20 J=1,NOBEL
      M=IBOUND(J)
      CALL INTEG(I,J,X(J),X(J+1),Y(J),Y(J+1),XM(N),YM(N),HH,GG,IN,DIFFUS
+(M),DT)
      IF(KODE(N).EQ.1)THEN
        IF(I.EQ.J)THEN
          COEFF(I,J)=0.50
        ELSE
          IF(KODE(M).EQ.1)COEFF(I,J)=HH
          IF(KODE(M).EQ.0)COEFF(I,J)=-GG
        ENDIF
      ELSE
        IF(KODE(M).EQ.1)COEFF(I,J)=HH
        IF(KODE(M).EQ.0)COEFF(I,J)=-GG
      ENDIF
      IF(KODE(M).EQ.1)RHMAT(I)=RHMAT(I)+BC(J)*GG
      IF((KODE(M).EQ.0).AND.(I.NE.J))RHMAT(I)=RHMAT(I)-BC(J)*HH
      IF((KODE(M).EQ.0).AND.(I.EQ.J))RHMAT(I)=RHMAT(I)-BC(J)*0.5
20 CONTINUE
  DO 10 I=1,NOEL
    K=1
    K1=LINKNO(I,K)
    IF((KODE(K1).EQ.0).OR.(KODE(K1).EQ.1))THEN
      XK1=XM(K1)
      YK1=YM(K1)
    ELSE
      XK1=XIN(K1)
      YK1=YIN(K1)
    ENDIF
    K2=LINKNO(I,K+1)
    IF((KODE(K2).EQ.0).OR.(KODE(K2).EQ.1))THEN
      XK2=XM(K2)
      YK2=YM(K2)
    ELSE
      XK2=XIN(K2)
      YK2=YIN(K2)
    ENDIF
    K3=LINKNO(I,K+2)
    IF((KODE(K3).EQ.0).OR.(KODE(K3).EQ.1))THEN
      XK3=XM(K3)
      YK3=YM(K3)
    ELSE
      XK3=XIN(K3)
      YK3=YIN(K3)
    ENDIF

```

```

ARM1=XK2*YK3-YK2*XK3
ARM2=-XK1*(YK3-YK2)
ARM3=YK1*(XK3-XK2)
AREA(I)=DABS(ARM1+ARM2+ARM3)/2.0
10 CONTINUE
IIN=0
DO 30 I=1,NOBEL
N=IBOUND(I)
CALL SOURCE(I,N,LINKNO,XIN,YIN,BBB,NOEL,XM,YM,AREA,IIN,KODE,THETAP
+,DIFFUS,DT)
RHMAT(I)=RHMAT(I)+BBB
30 CONTINUE
RETURN
END

```

C

C THIS SUBROUTINE WRITES THE VALUES IN THE OUTPUT FILES

C

```

SUBROUTINE OUTPUT(NOBEL,NINTER,BC,COLMAT,THETA,F,KODE,IBOUND,
+NINT,TT)
IMPLICIT REAL *8(A-H,O-Z)
PARAMETER(NB=20,NV=21,NI=5,NT=28,ND=25)
REAL*8 BC(NB),COLMAT(NB),THETA(F,ND)
INTEGER KODE(ND),IBOUND(NB),NINT(NI)
WRITE(3,37)TT
37 FORMAT(1X,'TIME = ',F5.2)
WRITE(3,36)
36 FORMAT(5X,'BOUNDARY ELEMENT SOLUTION OF POISSON'S
+EQUATION USING',/,12X,'BOUNDARY INTEGRATION FOR THE
+SOURCE TERM',/,5X,'
+_____
+',,////)
WRITE(3,35)
35 FORMAT(5X,'TABLE : SOLUTION AT THE BOUNDARY',/,5X,'
+_____
+',/)
WRITE(3,15)
15 FORMAT(5X,'NODE NO.',5X,'POTENTIAL',10X,'FLUX',/,5X,'
+_____
+',5X,'
+_____
',/)
DO 20 I=1,NOBEL
N=IBOUND(I)
IF(KODE(N).EQ.0)WRITE(3,25)N,BC(I),COLMAT(I)
25 FORMAT(7X,I2,7X,F10.5,6X,F10.5)
IF(KODE(N).EQ.1)WRITE(3,25)N,COLMAT(I),BC(I)
20 CONTINUE
WRITE(3,21)
21 FORMAT(5X,'
+_____
+',,////)
WRITE(3,45)

```

```

45  FORMAT(5X,'TABLE : SOLUTION IN THE DOMAIN',5X,'_____
+_____')
    WRITE(3,40)
40  FORMAT(/,5X,'NODE NO.',5X,'POTENTIAL',5X,'FLUX IN X-DIR.',4X,'FLUX
+ IN Z-DIR.',5X,'_____',5X,'_____',5X,'_____',4X,
+ '_____',/)
    DO 30 I=1,NINTER
    N=NINT(I)
    WRITE(3,50)N,THETAF(N)
50  FORMAT(8X,I2,6X,F10.5)
30  CONTINUE
    WRITE(3,31)
31  FORMAT(5X,'_____
+_____,////)
    WRITE(4,70)TT,THETAF(1),THETAF(2),THETAF(3),THETAF(4),THETAF(5)
70  FORMAT(6(3X,F10.5))
    RETURN
    END

```

C The program also includes the SUBROUTINES CAMPINT, NONLIN, GRAVIN,
C INTEG, ELIM, DATA, SOLINT, FLUXIN, FINTEG, SOURCE, FSOURC and
C FUNCTIONS ERROR1 AND ERROR2.

C The rest of the program is the same as in the DGF solution.

APPENDIX F

FORTRAN PROGRAMS FOR ANALYTICAL SOLUTIONS

C
 C ANALYTICAL SOLUTION FOR TWO-DIMENSIONAL MOISTURE-FLOW PROBLEM
 C

```

    IMPLICIT REAL *8(A-H,O-Z)
    REAL *8 THETA(10),X(10),Z(10),N,M,THETAP(10),DIFFUS(10),DTDZ(10),PP
    +(10),DTDX(10)
    EXTERNAL ANIN
    OPEN(2,FILE='ANAL1.DAT',STATUS='UNKNOWN')
    OPEN(3,FILE='ANAL1.OUT',STATUS='UNKNOWN')
    OPEN(4,FILE='COMP1.OUT',STATUS='UNKNOWN')
    OPEN(5,FILE='COMPX.OUT',STATUS='UNKNOWN')
    OPEN(6,FILE='COMPZ.OUT',STATUS='UNKNOWN')
    CALL ANIN(XX,ZZ,AA,BB,CC,DD,NINTER,X,Z,THETIN,SLOPMC,HCONS,DT,SAIE
    +,THETAS)
    WRITE(4,47)
47  FORMAT(5X,'TIME',5X,'NODE #1',5X,'NODE #2',5X,'NODE #3',5X,'NODE #
    +4',5X,'NODE #5',/)
    WRITE(6,47)
    WRITE(5,47)
    TT=0.0
17  T=0.0
15  T=T+DT
    TT=TT+DT
    J1=TT/DT
    BB=0.3+0.0982*T-0.0265*T**2+0.00107*T**3+0.0000595*T**4
    AA=0.3+0.0982*T-0.0265*T**2+0.00107*T**3+0.0000595*T**4
    DO 25 I=1,NINTER
    IF(J1.EQ.1)THEN
    THETAP(I)=THETIN
    ELSE
    THETAP(I)=THETA(I)
    ENDIF
    IF(J1.EQ.1)THEN
    PP(I)=0.0
    ELSE
    GRAV1=(2.0*SLOPMC+3.0)*HCONS*DTDZ(I)
    GRAV2=THETAP(I)**(2.0*SLOPMC+2.0)
    GRAV3=THETAS**(2.0*SLOPMC+3.0)
    PP(I)=GRAV1*GRAV2/GRAV3
    ENDIF
25  CONTINUE
    DO 10 I=1,NINTER
  
```

```

DIFF1=-SLOPMC*HCONS*SAIE/THETAS
DIFF2=(THETAP(I)/THETAS)**(SLOPMC+2.0)
DIFFUS(I)=DIFF1*DIFF2
N=-1
COSMN=0.0
SINMN=0.0
SINMXX=0.0
20 N=N+1
M=-1
SINM=0.0
SINMX=0.0
21 M=M+1
OMEGAN=(2.0*N+1.0)*22.0/(14.0*ZZ)
OMEGAM=(2.0*M+1.0)*22.0/(14.0*XX)
CHAT=4.0*CC/((2.0*N+1.0)*3.14)
SINNT=DSIN((2.0*N+1.0)*22.0/14.0)
BETAN1=2.0*(ZZ**2)*SINNT/((2.0*N+1.0)*3.14)
BETAN2=(2.0*ZZ/((2.0*N+1.0)*3.14))**2
BETAN3=BB-AA-DD*(BETAN1-BETAN2)/ZZ
BETAN=4.0*(BETAN3)/((2.0*N+1.0)*3.14)
GAMAN1=THETAP(I)-AA-DD*(BETAN1-BETAN2)/ZZ
GAMAN=4.0*GAMAN1/((2.0*N+1.0)*3.14)
SINMT=DSIN((2.0*M+1.0)*22.0/14.0)
GAMHA1=2.0*(XX**2)*SINMT/((2.0*M+1.0)*3.14)
GAMHA2=(2.0*XX/((2.0*M+1.0)*3.14))**2
GAMHA3=GAMAN-BETAN-CHAT*(GAMHA1-GAMHA2)/XX
GAMHAT=4.0*GAMHA3/((2.0*M+1.0)*3.14)
C2HAT=DIFFUS(I)*(OMEGAN**2)*CHAT
ALPHAC=DIFFUS(I)*CHAT/XX
ALPHAB=DIFFUS(I)*(OMEGAN**2)*BETAN
ALPHAD=DIFFUS(I)*DD/ZZ
TNBAR=4.0*(ALPHAD-PP(I))/((2.0*N+1.0)*3.14)
TNHAT=ALPHAC-ALPHAB-C2HAT*(X(I)**2)/(2.0*XX)+TNBAR
TMXT1=TNBAR+ALPHAC-ALPHAB-C2HAT*(GAMHA1-GAMHA2)/XX
TMN=4.0*(TMXT1)/((2.0*M+1.0)*3.14)
ALPHAT=DIFFUS(I)*(OMEGAN**2+OMEGAM**2)
CMT=TMN/ALPHAT+GAMHAT*DEXP(-ALPHAT*DT)
SINM1=CMT*DSIN(OMEGAM*X(I))
SINMX1=CMT*OMEGAM*DCOS(OMEGAM*X(I))
SINM=SINM+SINM1
SINMX=SINMX+SINMX1
IF(DABS(SINM1).GT.1.0E-10)GOTO 21
SINMM=SINM+BETAN+CHAT*(X(I)**2)/(2.0*XX)
SINMN1=SINMM*DSIN(OMEGAN*Z(I))
COSMN1=SINMM*OMEGAN*DCOS(OMEGAN*Z(I))
COSMN=COSMN+COSMN1
SINMN=SINMN+SINMN1
SINMNX=SINMX+CHAT*X(I)/XX
SINMNX1=SINMNX*DSIN(OMEGAN*Z(I))
SINMXX=SINMXX+SINMNX1

```

```

IF(DABS(SINMN1).GT.1.0E-8)GOTO 20
IF(DABS(COSMN1).GT.1.0E-8)GOTO 20
IF(DABS(SINMNX1).GT.1.0E-8)GOTO 20
THETA(I)=SINMN+AA+(Z(I)**2)*DD/(2.0*ZZ)
DTDZ(I)=COSMN+Z(I)*DD/ZZ
DTDZ(I)=SINMXX
10 CONTINUE
WRITE(*,*)'TT=',TT
WRITE(4,16)TT,THETA(1),THETA(2),THETA(3),THETA(4),THETA(5)
WRITE(6,16)TT,DTDZ(1),DTDZ(2),DTDZ(3),DTDZ(4),DTDZ(5)
WRITE(5,16)TT,DTDZ(1),DTDZ(2),DTDZ(3),DTDZ(4),DTDZ(5)
16 FORMAT(6(3X,F10.5))
WRITE(3,*)'TT=',TT
WRITE(3,12)
12 FORMAT(8X,'ANALYTICAL SOLUTION FOR RECHARD'S EQUATION',/11X,'WITH
+ FLUX AT THE UPPER BOUNDARY',/8X,'
+ _____
+ _____',//)
WRITE(3,11)
11 FORMAT(5X,'TABLE : SOLUTION FOR INTERNAL NODES',/5X,'
+ _____',/
WRITE(3,2)
2 FORMAT(5X,'NODE NO.',3X,'POTENTIAL',3X,'FLUX IN X-DIR',3X,'FLUX IN
+ Z-DIR',/5X,' _____',3X,' _____',3X,' _____',3X,'
+ _____',/
DO 30 I=1,NINTER
WRITE(3,1)I,THETA(I)
1 FORMAT(5X,I3,5X,F10.5)
30 CONTINUE
WRITE(3,31)
31 FORMAT(5X,' _____',//)
IF(T.LT.10.)GOTO 15
IF(TT.LT.30.0)GOTO 17
35 STOP
END

```

C

C THIS SUBROUTINE READS DATA FOR TWO-DIMENSIONAL ANALYTICAL
C SOLUTION OF THE UNSATURATED MOISTURE-FLOW EQUATION

C

```

SUBROUTINE ANIN(XX,ZZ,AA,BB,CC,DD,NINTER,X,Z,THETIN,SLOPMC,HCONS,D
+T,SAIE,THETAS)
IMPLICIT REAL *8(A-H,O-Z)
REAL*8 X(10),Z(10)
READ(2,*)NINTER
READ(2,*)XX,ZZ,THETIN,SLOPMC,HCONS,DT,SAIE,THETAS
READ(2,*)AA,BB,CC,DD
DO 10 I=1,NINTER

```

```

      READ(2,*)X(I),Z(I)
10  CONTINUE
      RETURN
      END

```

```

C
C ANALYTICAL SOLUTION FOR ONE-DIMENSIONAL MOISTURE-FLOW PROBLEM
C

```

```

      IMPLICIT REAL *8(A-H,O-Z)
      REAL *8 THETAF(50),ZZ(50),THETAP(50),DIFFUS(50)
      EXTERNAL DATA,ERROR
      OPEN(2,FILE='DATA.DAT',STATUS='UNKNOWN')
      OPEN(4,FILE='ANAL.OUT',STATUS='UNKNOWN')
      CALL DATA(BC,THETIN,TIME,DT,SLOPMC,SAIE,THETAS,NODE,HCONS,ZZ)
      WRITE(4,47)
47  FORMAT(5X,'TIME',5X,'NODE 4',5X,'NODE 7',5X,'NODE 10',5X,'NODE 1
      +3',5X,'NODE 16',/)
      TT=0.0
12  T=0.0
11  T=T+DT
      TT=TT+DT
      J1=TT/DT
      BC=0.3+0.0982*T-0.0265*(T**2)+0.00107*(T**3)+0.0000595*(T**4)
      DO 45 I=1,NODE
      IF(J1.EQ.1)THEN
      THETAP(I)=THETIN
      ELSE
      THETAP(I)=THETAF(I)
      ENDIF
45  CONTINUE
      DO 40 I=1,NODE
      DIFF1=-SLOPMC*HCONS*SAIE/THETAS
      TPP=THETAP(I)
      DIFF2=(TPP/THETAS)**(SLOPMC+2.0)
      DIFFUS(I)=DIFF1*DIFF2
40  CONTINUE
      DO 10 I=1,NODE
      ARG=ZZ(I)/(2.0*DSQRT(DIFFUS(I)*T))
      THETAF(I)=BC+(THETAP(I)-BC)*ERROR(ARG)
10  CONTINUE
      WRITE(4,70)TT,THETAF(4),THETAF(7),THETAF(10),THETAF(13),THETAF(16)
70  FORMAT(6(3X,F10.5))
      IF(T.LT.TIME)GOTO 11
      IF(TT.LT.30.0)GOTO 12
      STOP
      END

```

C
 C THIS SUBROUTINE COMPUTES ERROR FUNCTION FOR ONE-DIMENSIONAL
 C ANALYTICAL SOLUTION OF THE MOISTURE-FLOW EQUATION
 C

```

REAL*8 FUNCTION ERROR(Z)
  IMPLICIT REAL *8(A-H,O-Z)
  REAL*8 Z
  PP=0.3275911
  A1=0.254829592
  A2=-0.284496736
  A3=1.421413741
  A4=-1.453152027
  A5=1.061405429
  TT=1.0/(1.0+PP*Z)
  ERF1Z=A1*TT + A2*TT**2 + A3*TT**3 + A4*TT**4 + A5*TT**5
  ERFCZ=ERF1Z*DEXP(-Z**2)
  ERROR=1.0-ERFCZ
  RETURN
  END

```

C
 C THIS SUBROUTINE READS DATA FOR ONE-DIMENSIONAL ANALYTICAL SOLUTION
 C OF THE UNSATURATED MOISTURE-FLOW EQUATION
 C

```

SUBROUTINE DATA(THETIN,TIME,DT,SLOPMC,SAIE,THETAS,NODE,HCONS,ZZ)
  IMPLICIT REAL *8(A-H,O-Z)
  REAL *8 ZZ(50)
  READ(2,*)NODE,HCONS,THETAS,SAIE,SLOPMC,DT,TIME,THETIN
  DO 10 I=1,NODE
  READ(2,*)ZZ(I)
  ZZ(I)=0.6*ZZ(I)
10 CONTINUE
  RETURN
  END

```

C
 C ANALYTICAL SOLUTION FOR SIMPLIFIED LEACHATE FLOW
 C EQUATION IN THE SATURATED ZONE
 C

```

  IMPLICIT REAL *8(A-H,O-Z)
  PARAMETER (NT=25)

```

```

INTEGER LCONL(NT),LCONR(NT),ICOL(NT)
REAL *8 THICK(NT),HDIN(NT),SOURCE(NT),HDP(NT),HPP(NT),HPS(NT),HDD(
+NT),X(NT)
EXTERNAL ELIM,SATDAT
OPEN(2,FILE='SAT.DAT',STATUS='OLD')
OPEN(3,FILE='SAT.OUT',STATUS='UNKNOWN')
OPEN(4,FILE='COMPAN.OUT',STATUS='UNKNOWN')
CALL SATDAT(THICK,CONDB,DX,HCONS,NOCOL,LCONL,LCONR,ICOL,EFPORO,DT,
+SOURCE,HDIN,SLOPE,BCONL,BCONR,TIME,EPS)
AVTHICK=0.0
XLENG=DX*(NOCOL+1.0)
DO 55 I=1,NOCOL
X(I)=I*DX
55 CONTINUE
T=0.0
11 T=T+1.0
J1=T/1.0
DO 45 I=1,NOCOL
IF(J1.EQ.1)THEN
HDP(I)=HDIN(I)
HPP(I)=HDP(I)
ELSE
HPP(I)=HDP(I)
HDP(I)=HDD(I)
ENDIF
45 CONTINUE
DO 10 I=1,NOCOL
HPS(I)=(HDP(I)+HPP(I))/2.0
10 CONTINUE
DO 35 I=1,NOCOL
HDD(I)=0.0
35 CONTINUE
DO 15 I=1,NOCOL
N=-1
31 N=N+1
OMEGA=(2.0*N+1.0)*22.0/(14.0*XLENG)
XSINX=((1.0/OMEGA)**2)*((-1.0)**N)
COSX=14.0*XLENG*((-1.0)**N)/((2.0*N+1.0)*22.0)
XCOSX=14.0*(XLENG**2)*((-1.0)**N)/((2.0*N+1.0)*22.0)-(1.0/OMEGA)**
+2
X2COSX=14.0*(XLENG**3)*((-1.0)**N)/((2.0*N+1.0)*22.0)-28.0*XLENG*X
+SINX/((2.0*N+1.0)*22.0)
CNO=2.0*HDP(I)*COSX/XLENG-2.0*BCONL*XCOSX/XLENG+2.0*BCONL*X2COSX/(
+XLENG**2)-2.0*BCONR*X2COSX/(XLENG**3)
ALPHA=HCONS*HPS(I)/EFPORO
ALP2HAT=ALPHA*(OMEGA**2)
PB1=CONDB
PHAT=(-SOURCE(I)+PB1)/EFPORO
BETA=-CONDB/(THICK(I)*EFPORO)
TN1=4.0*ALPHA*BCONR*COSX/(XLENG**3)

```

```

TN2=4.0*ALPHA*BCONL*COSX/(XLENG**2)
TN3=2.0*BETA*BCONL*XCOSX/XLENG
TN4=2.0*BETA*BCONL*X2COSX/XLENG**2
TN5=2.0*BETA*BCONR*X2COSX/XLENG**3
TN6=2.0*PHAT*COSX/XLENG
TNHAT=TN1-TN2+TN3-TN4+TN5-TN6
ALPHAT=ALP2HAT-BETA
GG=CNO-TNHAT/ALPHAT
CNT=TNHAT/ALPHAT+GG*DEXP(-ALPHAT*DT)
HDST1=CNT*DCOS(OMEGA*X(I))
HDD(I)=HDD(I)+HDST1
IF(DABS(HDST1).GT.1.0E-10)GOTO 31
AX=BCONL*X(I)*(1.0-X(I)/(XLENG))+BCONR*(X(I)**2)/(XLENG**2)
HDD(I)=HDD(I)+AX
15 CONTINUE
IF(J1.EQ.1)THEN
WRITE(4,42)
42 FORMAT(1X,'TIME',2X,'# 8',2X,'# 9',2X,'# 10',2X,'# 11',2X,'# 12',2
+X,'# 13',/)
ELSE
ENDIF
WRITE(4,41)T,HDD(8),HDD(9),HDD(10),HDD(11),HDD(12),HDD(13)
41 FORMAT(1X,F5.2,6(F10.2))
WRITE(3,50)T
50 FORMAT(5X,'TIME T = ',F5.2,1X,'DAYS',/,5X,'NODE NO.',10X,'MOUND H
+EAD (FT)')
DO 52 I=1,NOCOL
WRITE(3,51)I,HDD(I)
51 FORMAT(5X,I2,8X,F10.2)
52 CONTINUE
IF(T.LT.TIME)GOTO 11
STOP
END

```

C

C THIS SUBROUTINE READS DATA FOR THE ANALYTICAL SOLUTION
C OF THE SATURATED LEACHATE FLOW EQUATION

C

```

SUBROUTINE SATDAT(THICK,CONDB,DX,HCONS,NOCOL,LCONL,LCONR,ICOL,
+EFPORO,DT,SOURCE,HDIN,SLOPE,BCONL,BCONR,TIME,EPS)
IMPLICIT REAL*8(A-H,O-Z)
PARAMETER(NT=25)
REAL *8 THICK(NT),SOURCE(NT),HDIN(NT)
INTEGER LCONL(NT),LCONR(NT),ICOL(NT)
READ(2,*)CONDB,DX,DT,HCONS,EFPORO,SLOPE
READ(2,*)BCONL,BCONR,TIME,EPS

```

```
READ(2,*)NOCOL
DO 10 I=1,NOCOL
ICOL(I)=1
READ(2,*)LCONL(I),LCONR(I),SOURCE(I),THICK(I),HDIN(I)
10 CONTINUE
ICOL(NOCOL+1)=2
ICOL(NOCOL+2)=2
RETURN
END
```

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